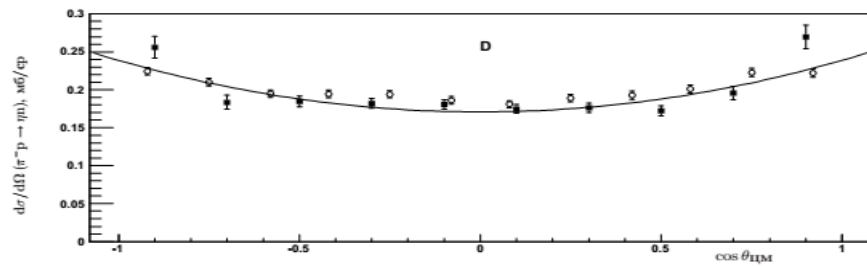
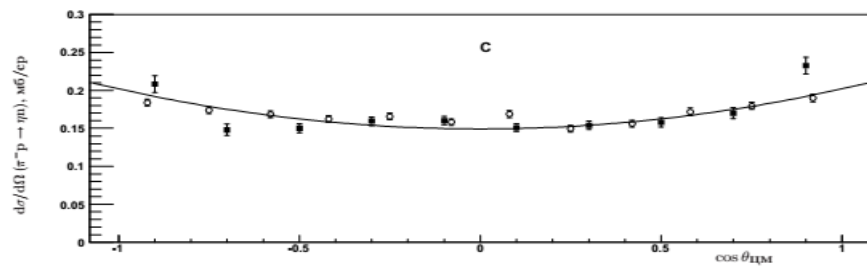
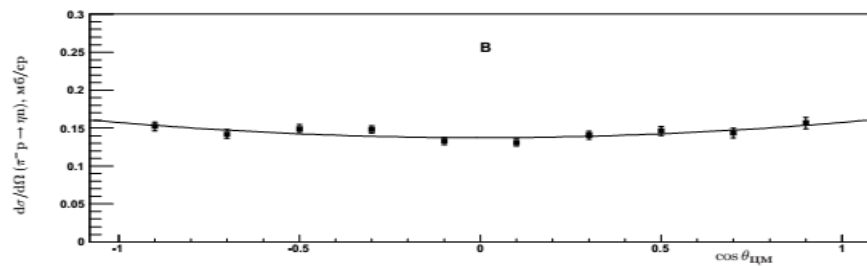
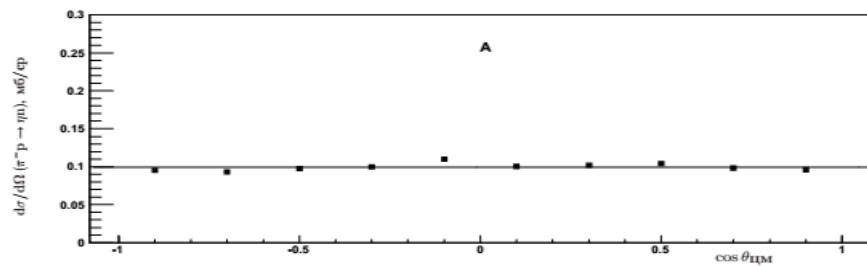
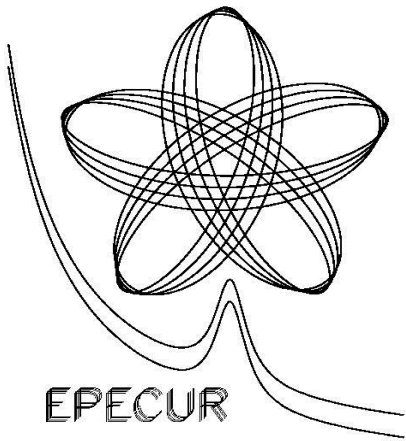




PNPI & EPECUR.





New results on narrow structure in the pion nucleon elastic scattering from the EPECUR experimentn.

I.G. Alekseev, I.G. Bordyuzhin, D.A. Fedin, V.P. Kanavets , L.I. Koroleva,
B.V. Morozov, V.M. Nesterov, V.V. Ryltsov, D.N. Svirida, A.D. Sulimov
ITEP, Moscow

V.A. Andreev, Ye.A. Filimonov, A.B. Gridnev, V.V. Golubev, A.I. Kovalev,
N.G. Kozlenko, V.S. Kozlov, A.G. Krivshich, D.V. Novinsky, V.V. Sumachev, V.I. Tarakanov,
V.Yu. Trautman
PNPI, Gatchina

M. Sadler
ACU, Abilene

Pentaquark antidecuplet

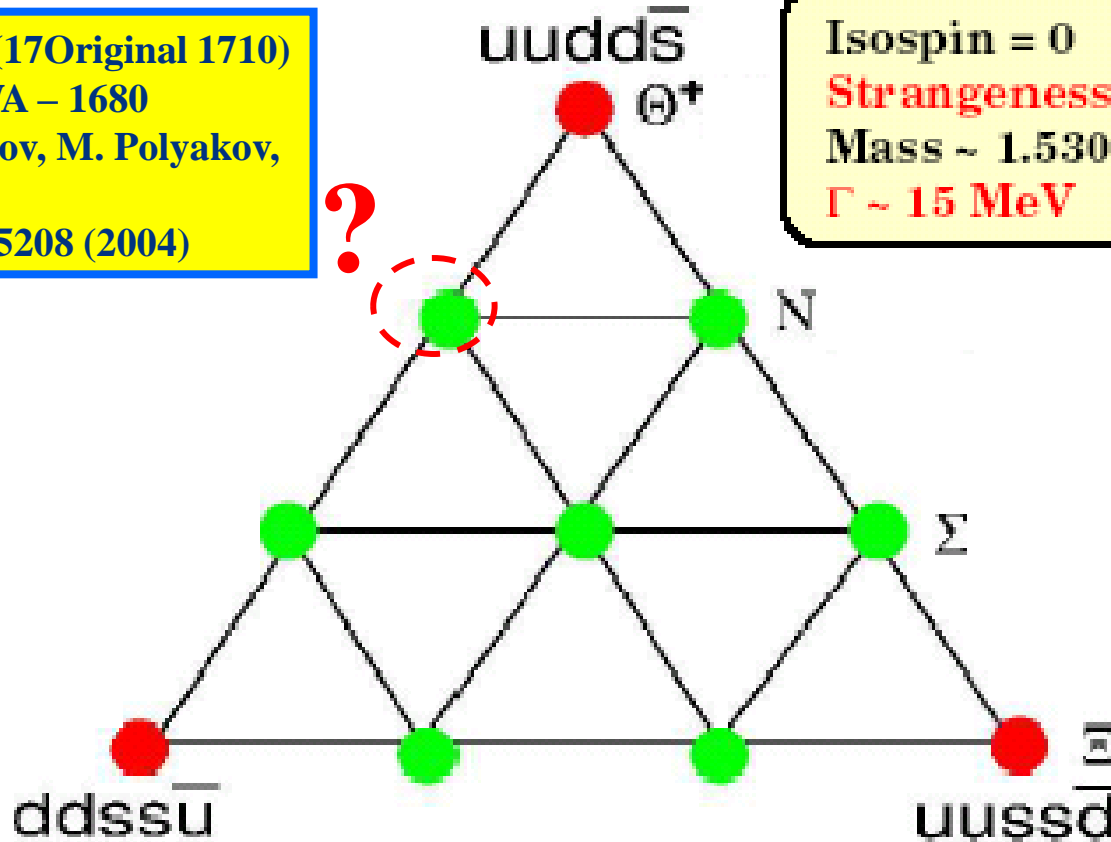


$[\bar{10}]$ Spin = $\frac{1}{2}$ NEW MULTIPLET

D.Diakonov et al. Z. Phys A359, 1997, 305

prediction – N***(17Original 1710)
From modified PWA – 1680
R. Arndt, Ya. Azimov, M. Polyakov,
IS, R. Workman,
Phys Rev C 69, 035208 (2004)

Isospin = 0
Strangeness = +1
Mass ~ 1.530 MeV
 $\Gamma \sim 15$ MeV

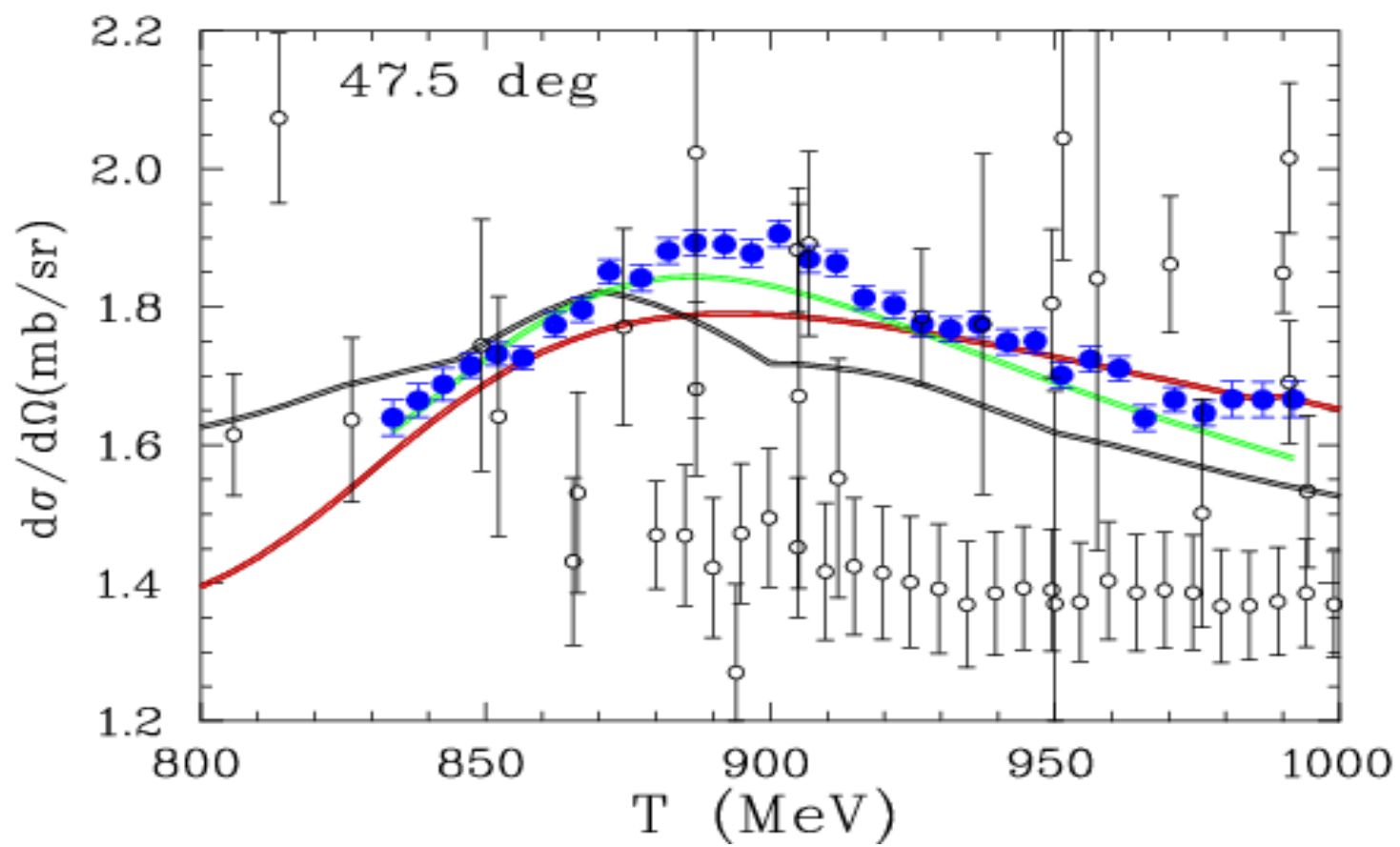


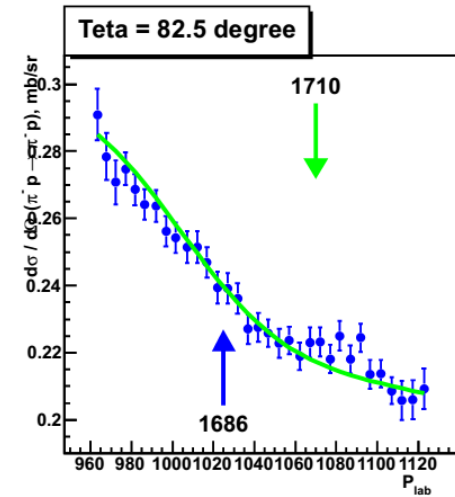
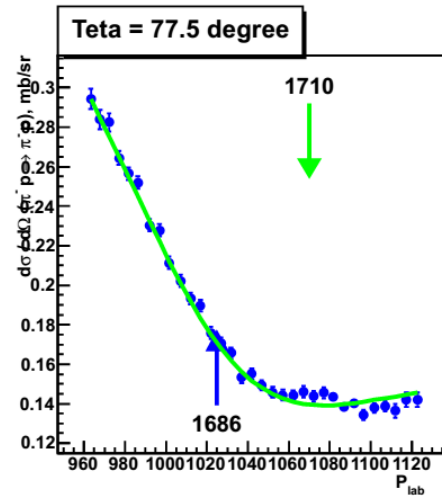
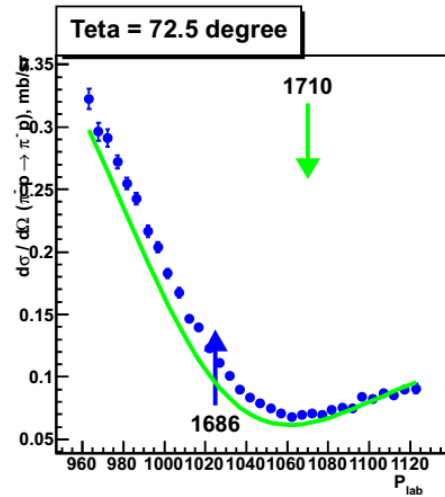
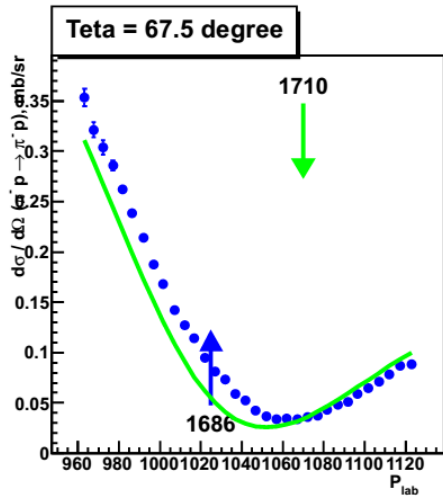
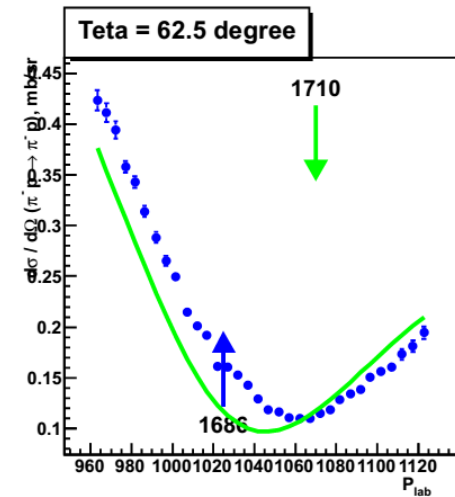
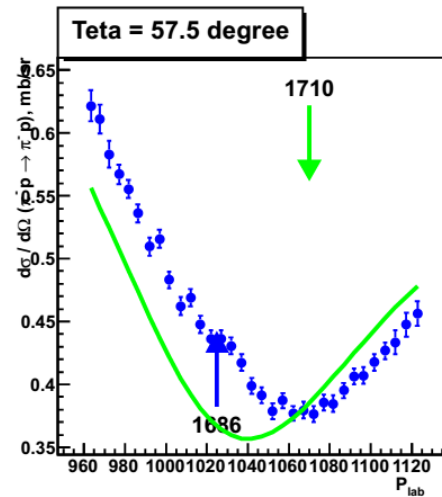
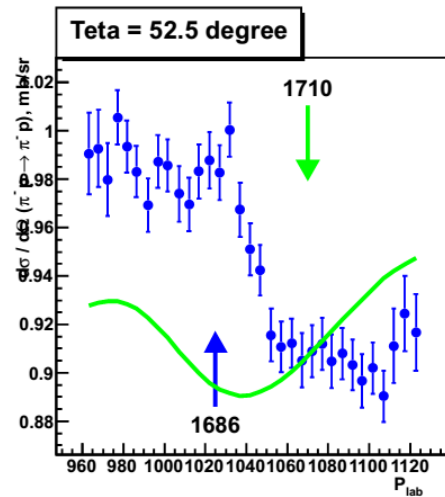
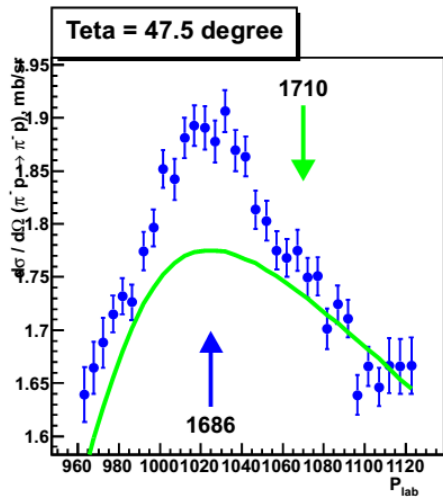


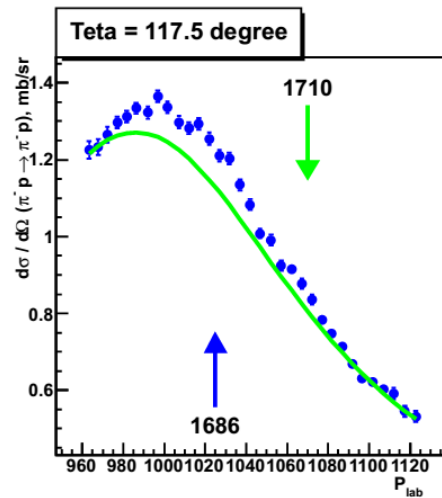
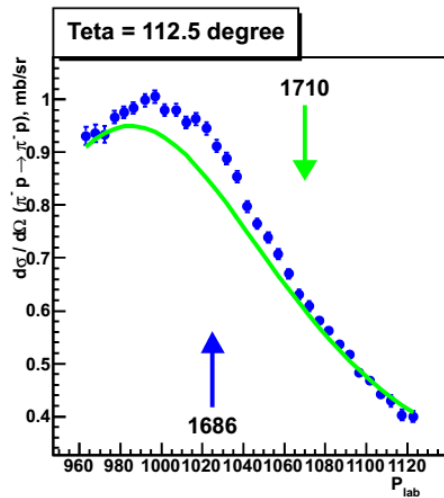
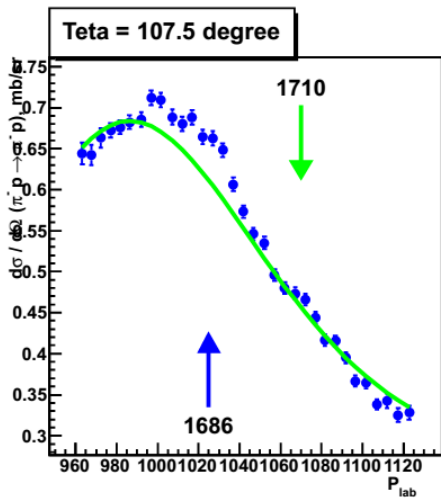
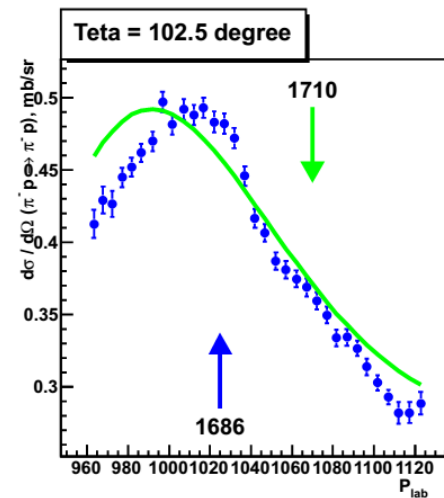
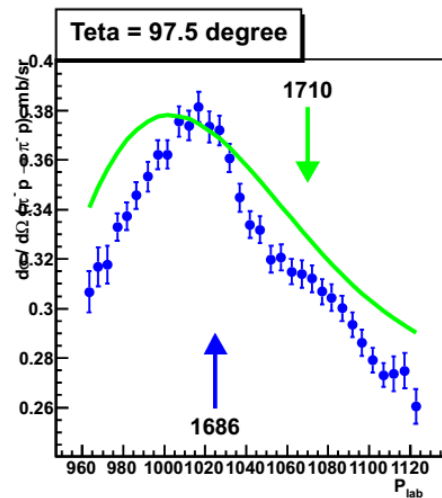
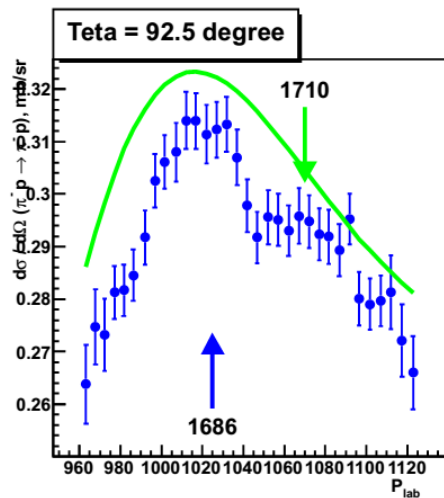
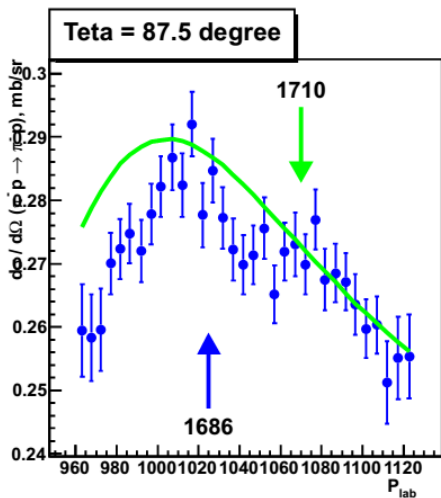
Why pions? Theory gives weak coupling to πN sector.

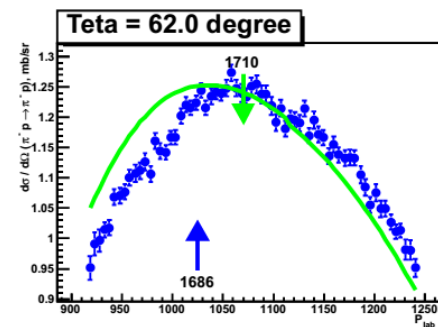
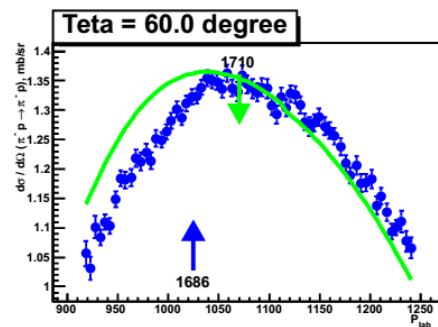
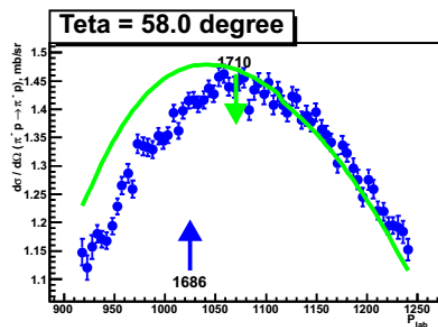
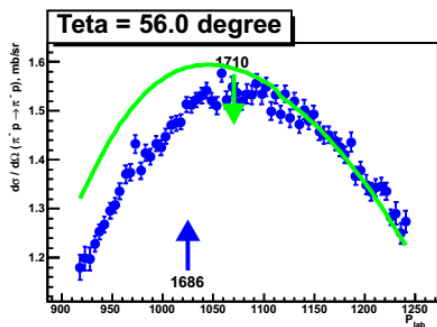
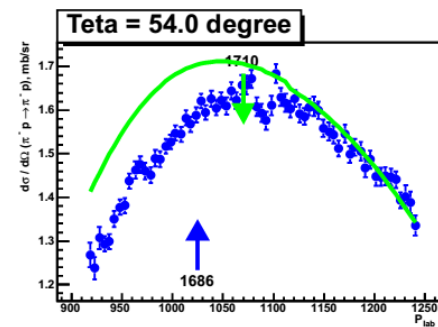
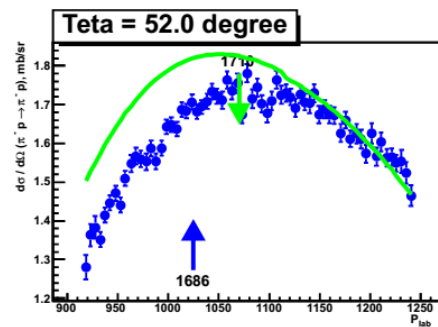
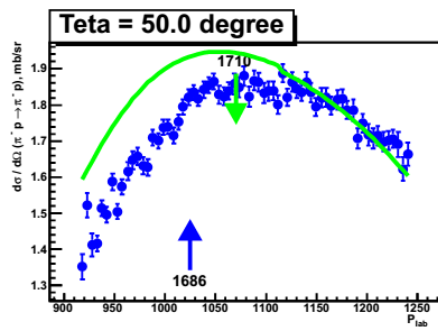
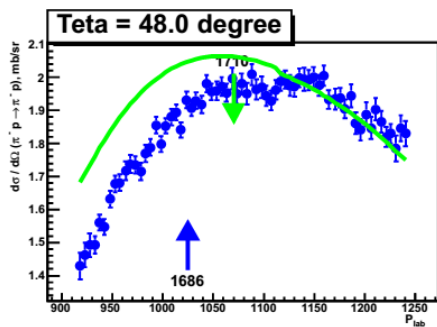
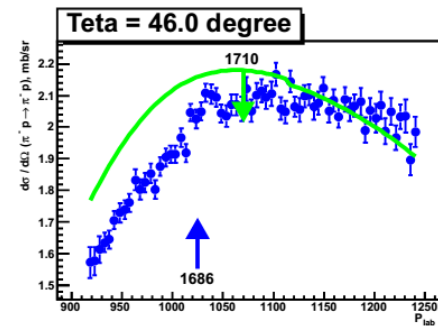
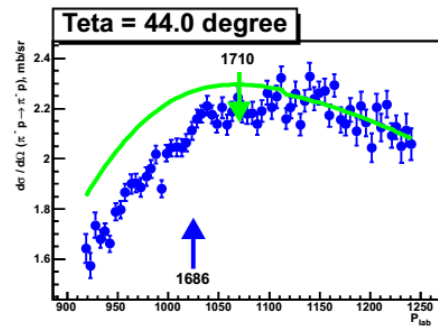
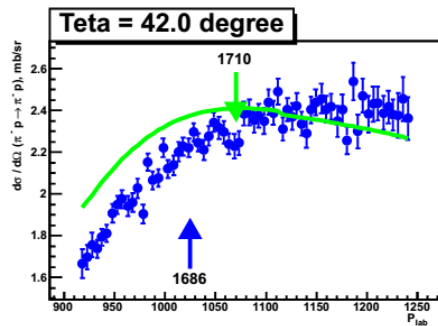
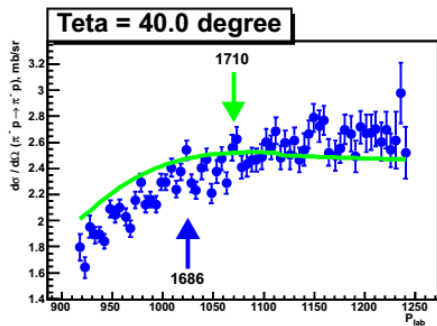
Advantages:

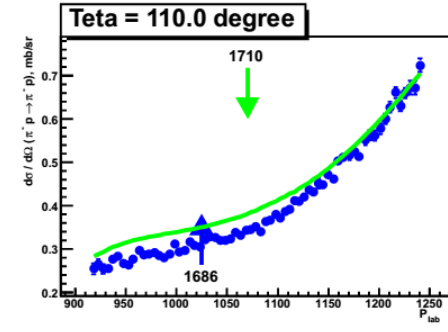
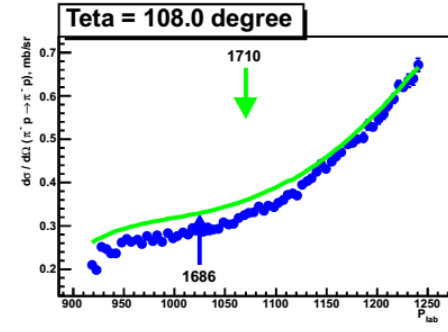
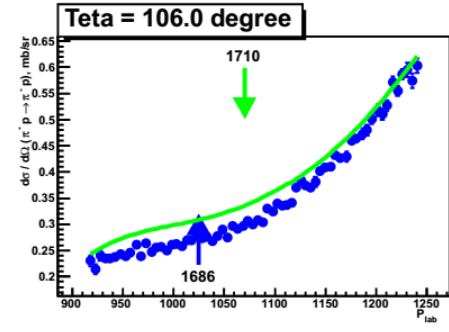
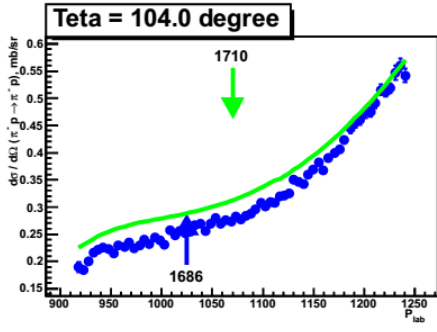
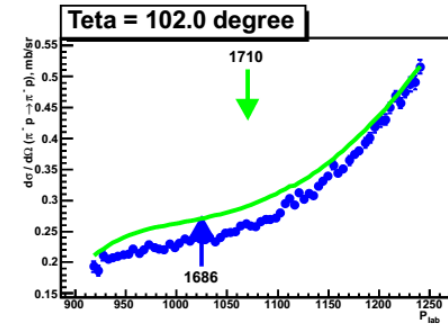
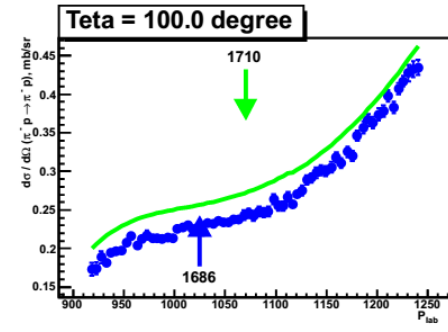
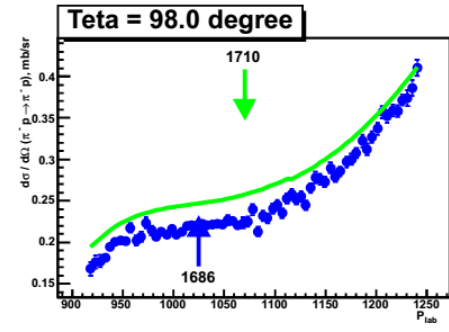
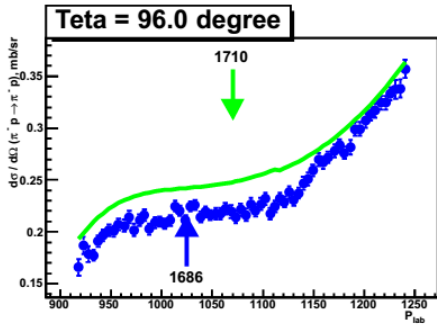
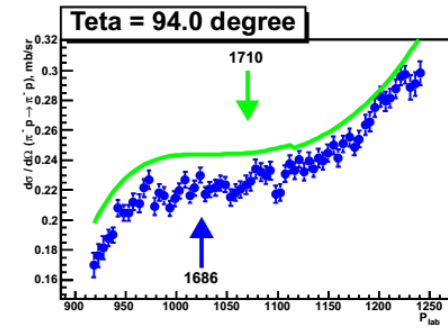
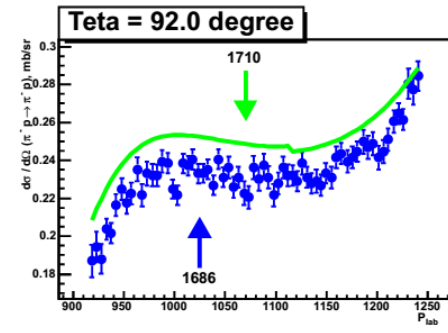
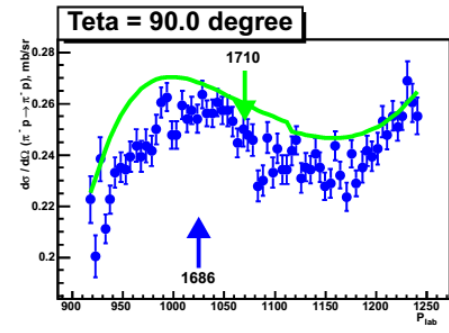
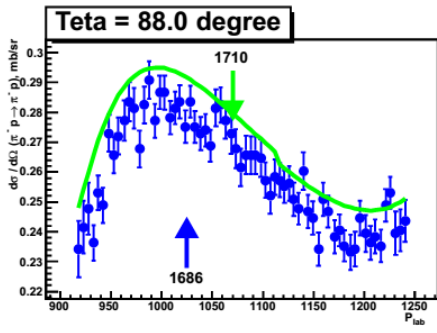
- 1. The structure of the πN amplitude is much more simpler than in photoproduction*
- 2. The πN partial waves are known very well from phase shift analysis.*
- 3. There is isospin symmetry.*
- 4. The number of free parameters smaller because there is no photocouplings*













How we can add resonance and background?

- 1. Sum of amplitudes \rightarrow can violate unitarity condition.*
- 2. Sum of phase shifts \rightarrow in classical physics it is true for INDEPENDENT sources. Not unique procedure.*

K-matrix approach with effective Lagrangians.

P.F.A. Goudsmit et al Nucl.Phys A575 (1994)673

A.B. Gridnev, N.G. Kozlenko. Eur.Phys.J.A4:187-194, (1999).

T. Feuster and U. Mosel Phys. Rec. C 58 457 (1998).



It is assumed that the K -matrix, being a solution of the equation for scattering amplitude, can be considered as a sum of the tree-level Feynman diagrams with the effective Lagrangians in the vertices.

4 resonances in s and u channels and sigma, rho like exchange in t channel.*

Multichannel:

- 1. elastic scattering*
- 2. two pion production (effective)*
- 3. η n production*
- 4. K Λ production*
- 5. K Σ production*



Free parameters \rightarrow coupling constants.

We concentrate on elastic scattering and treat inelastic channels approximately to save the number of free parameters.

Database:

EPECURE results.

SAID single energy solutions up to 1 GeV.

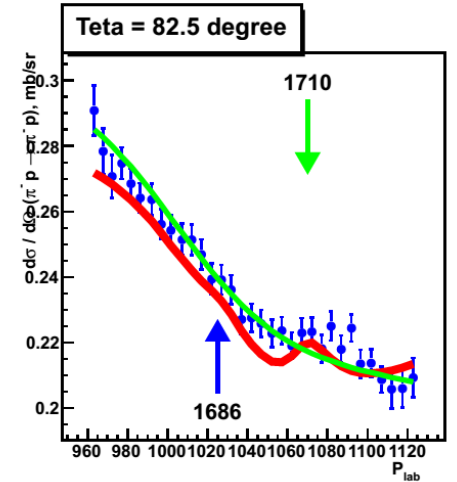
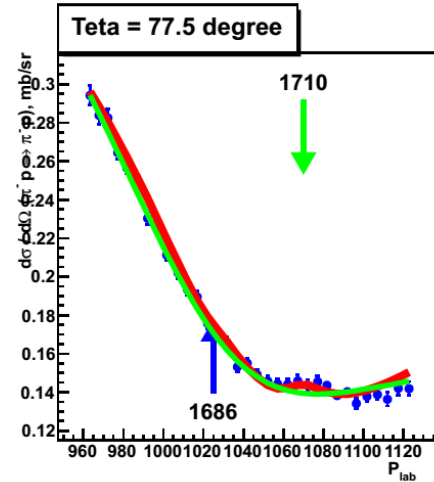
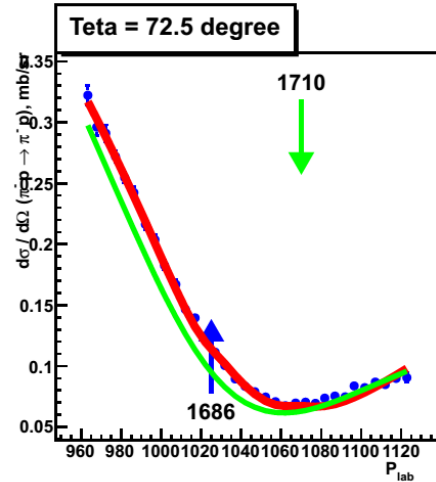
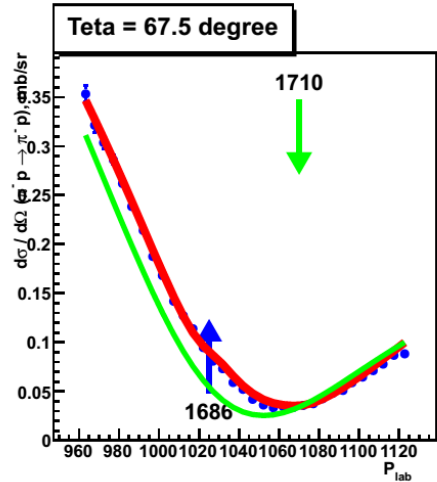
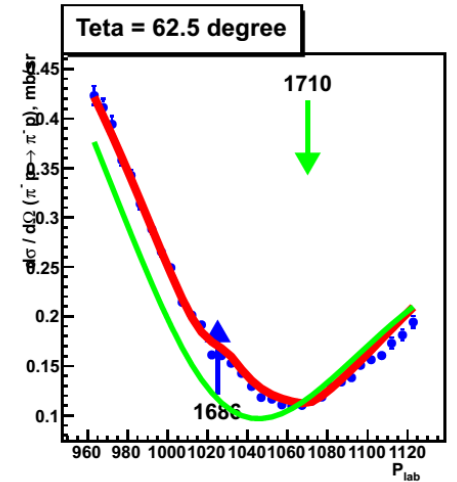
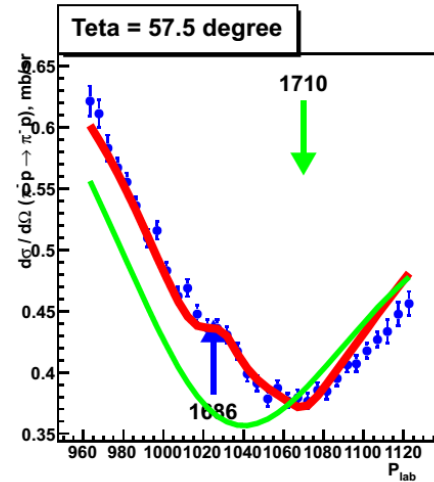
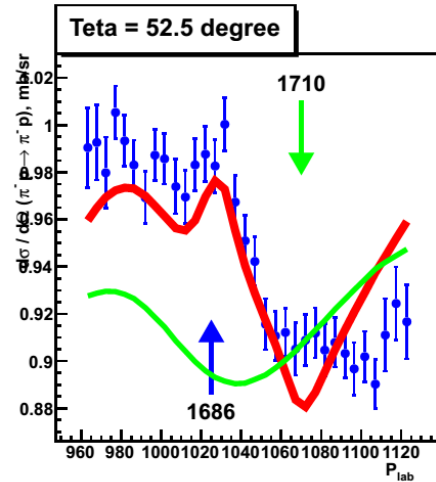
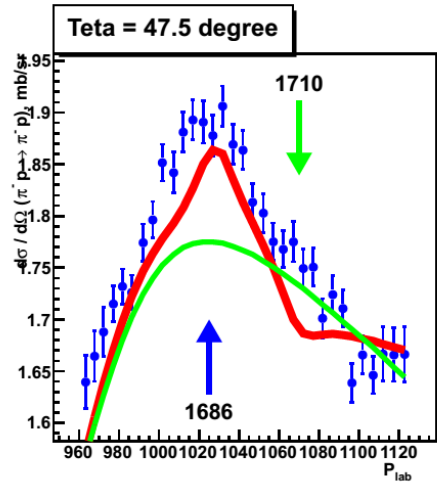
η n total cross section

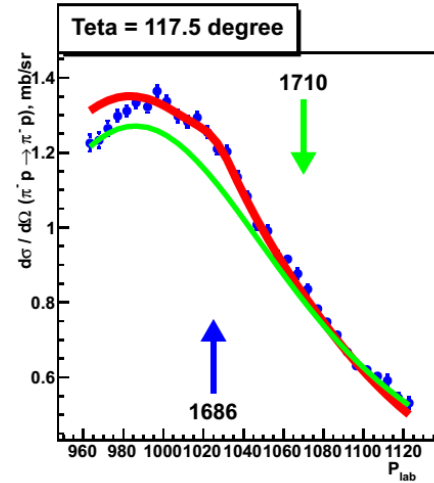
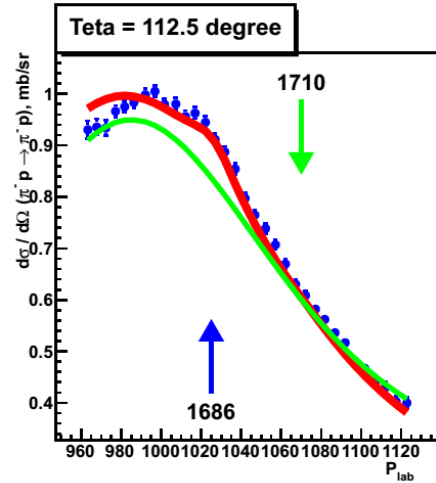
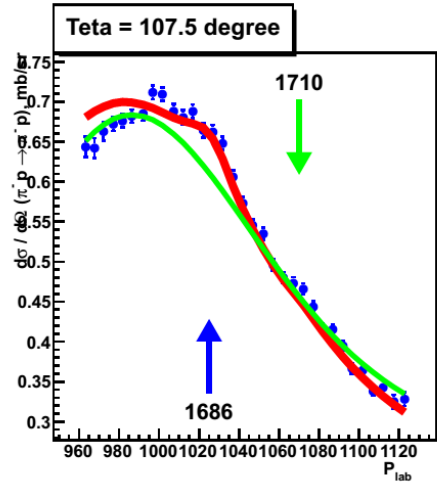
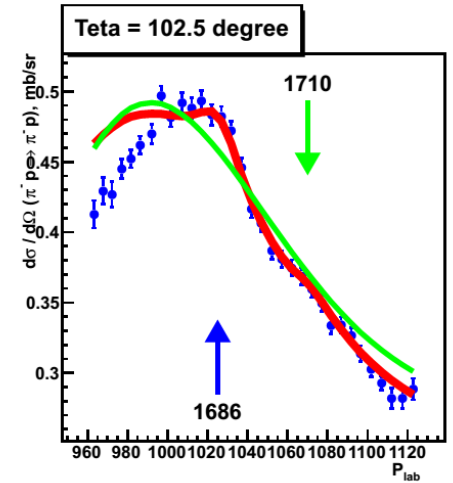
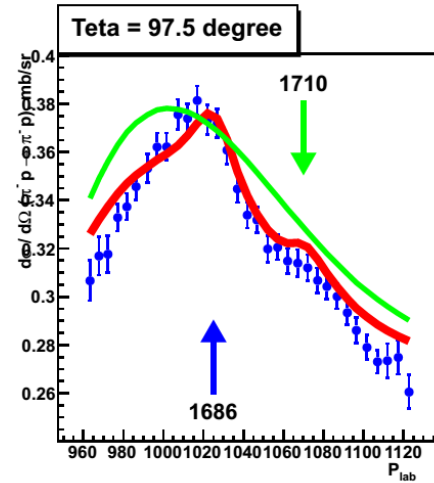
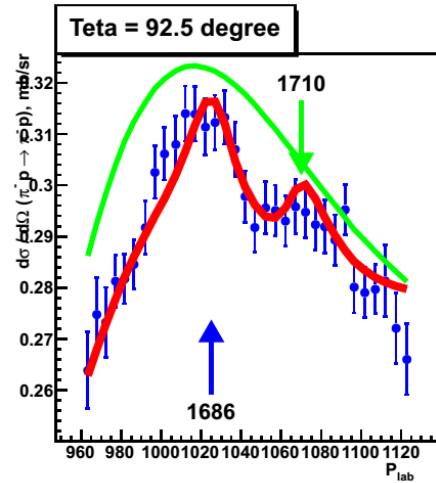
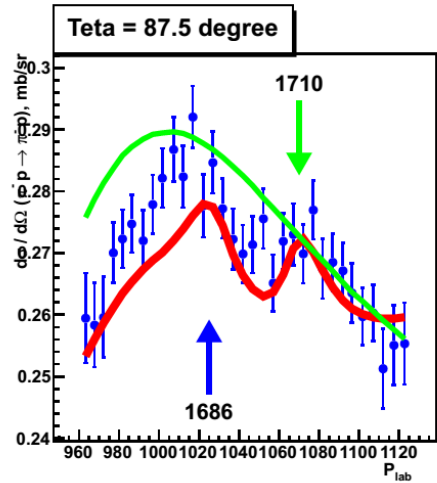
$K \Lambda$ differential cross section

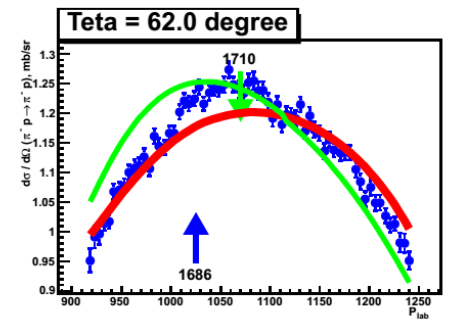
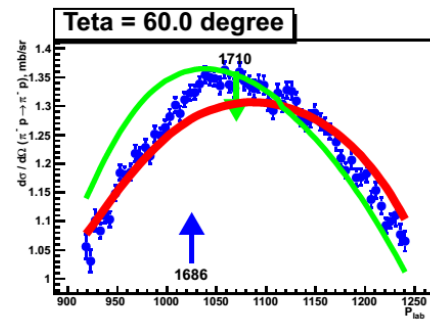
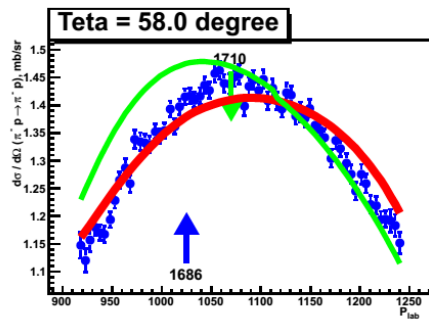
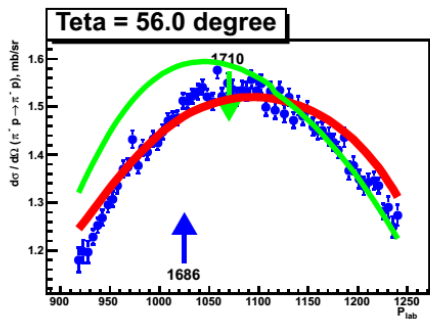
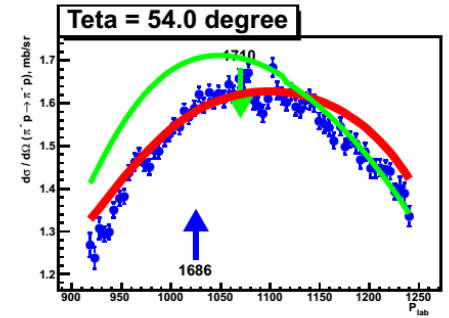
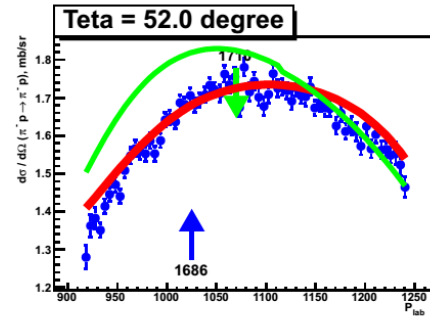
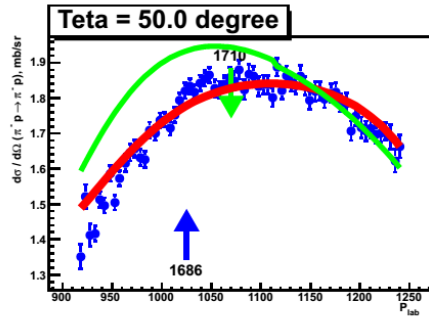
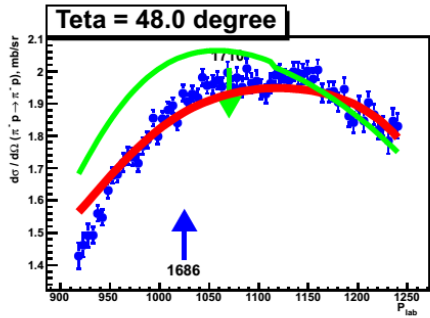
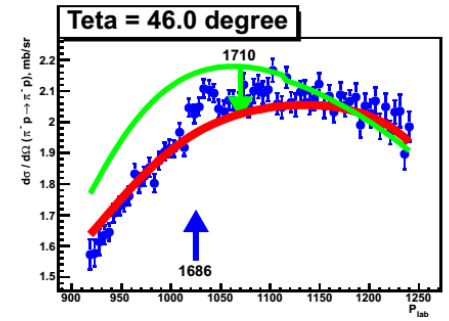
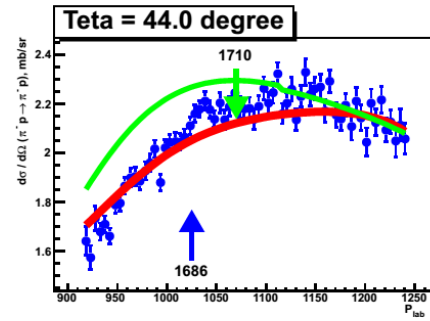
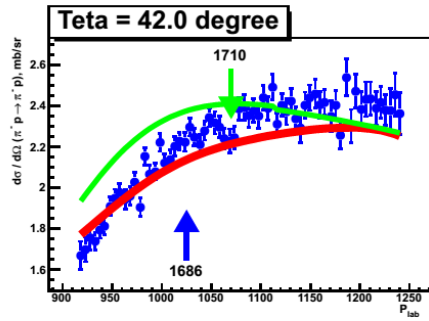
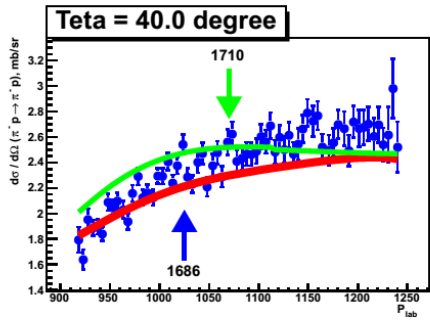
$K^0 \Sigma^0$ differential cross section

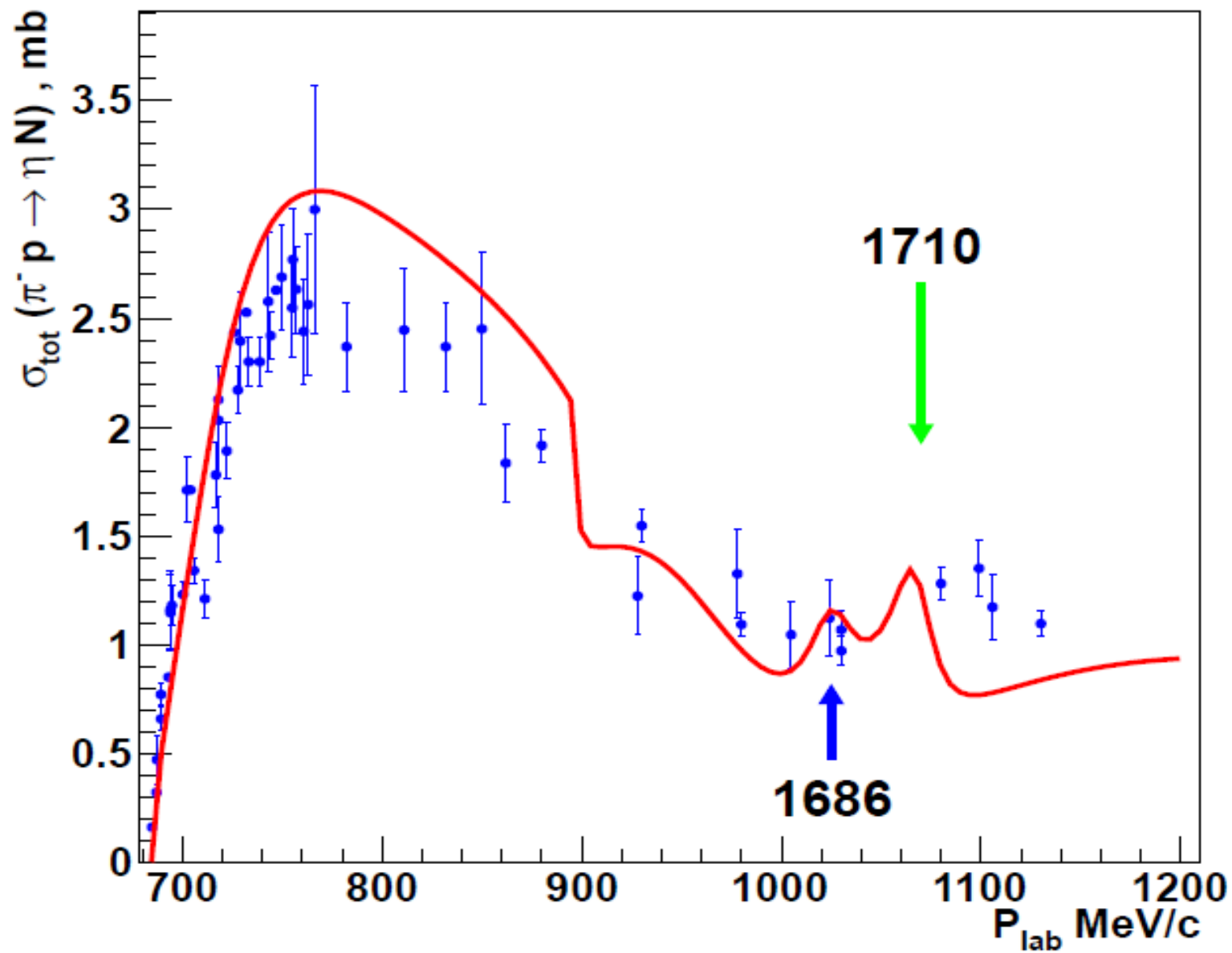
$K^+ \Sigma^+$ differential cross section

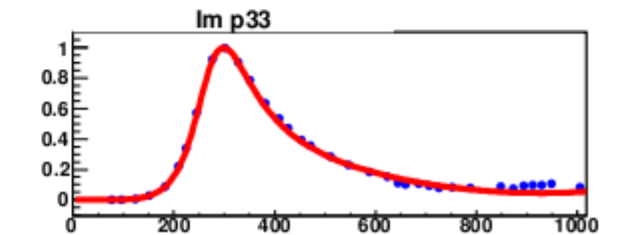
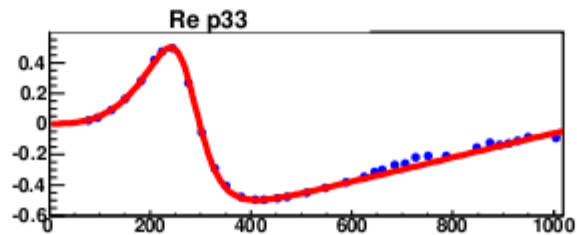
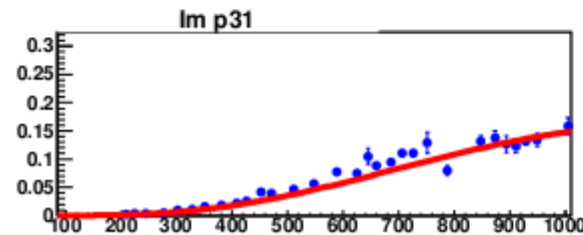
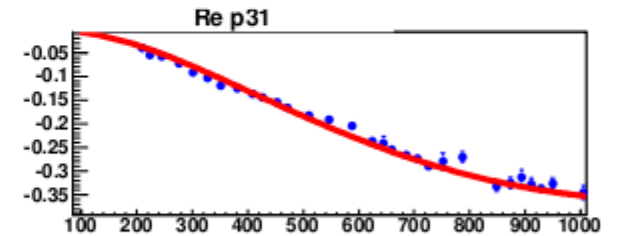
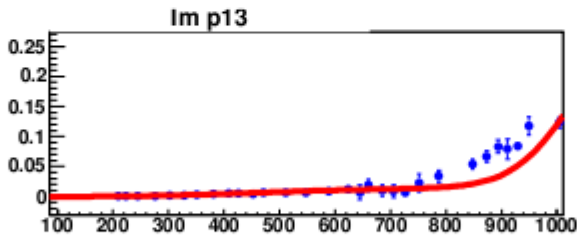
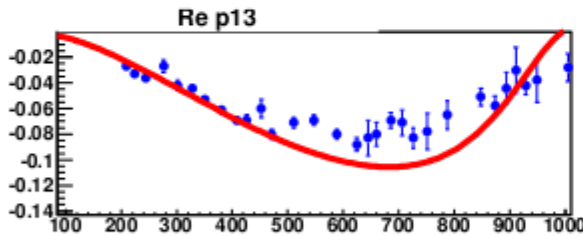
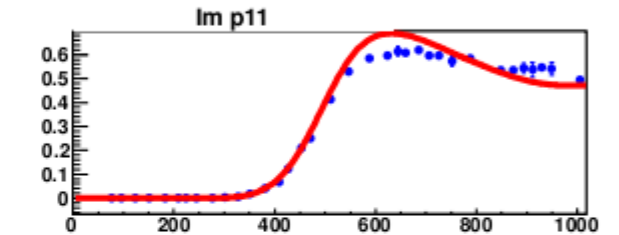
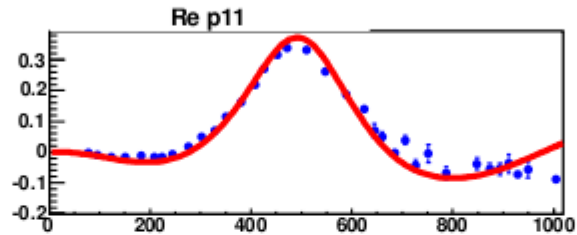
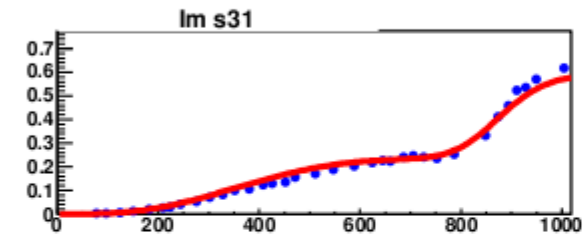
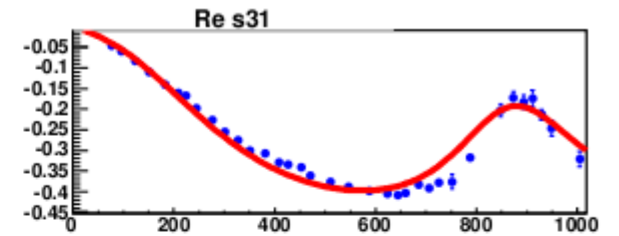
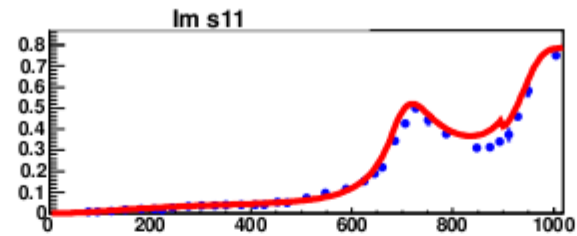
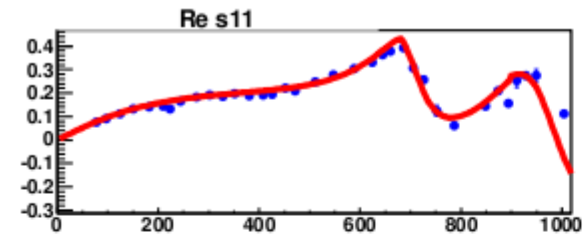
$K^+ \Sigma^-$ differential cross section

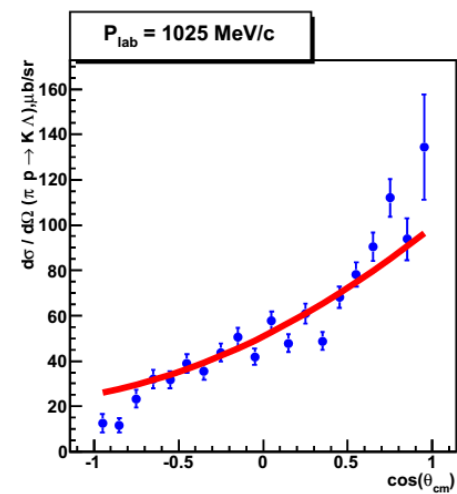
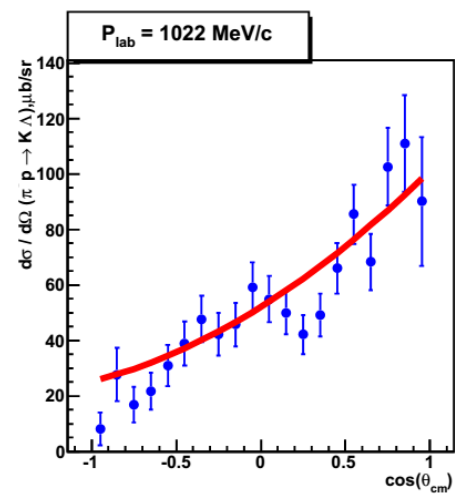
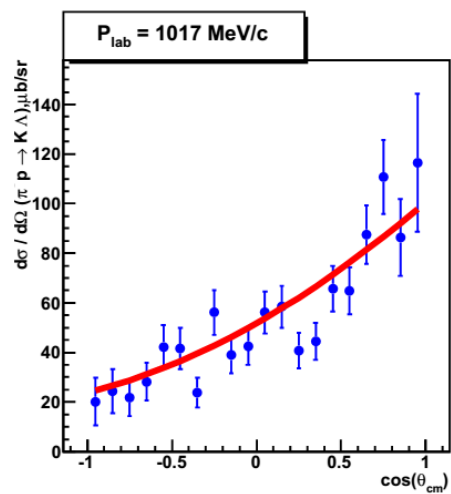
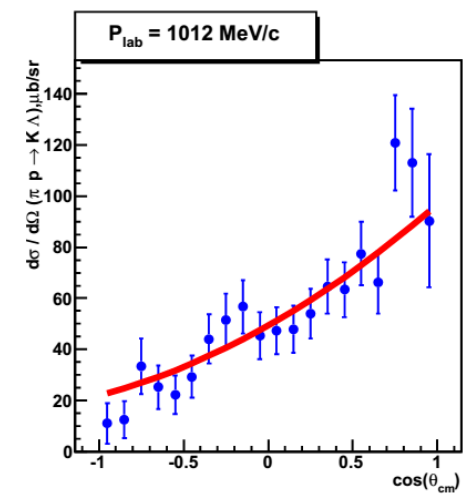
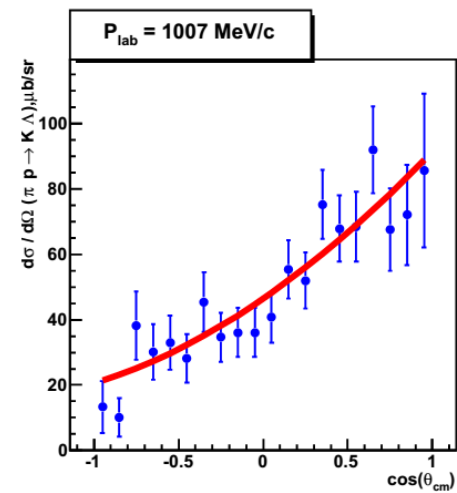
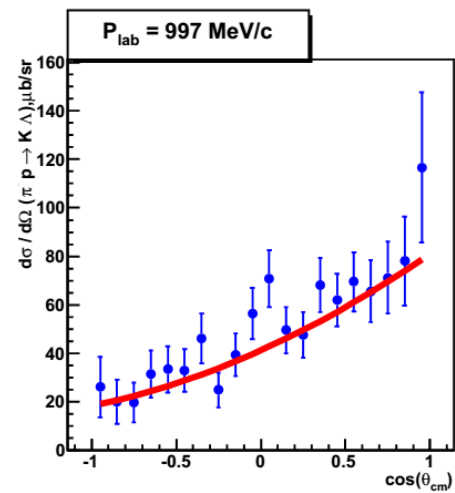
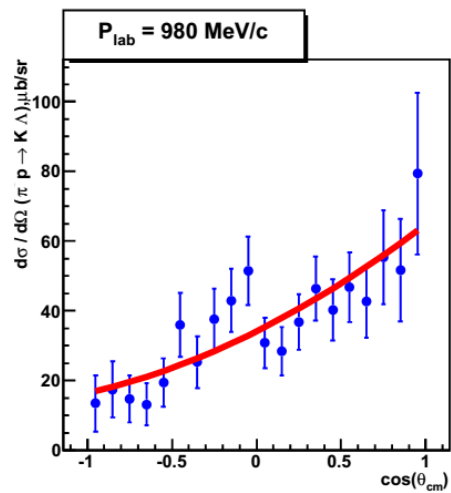
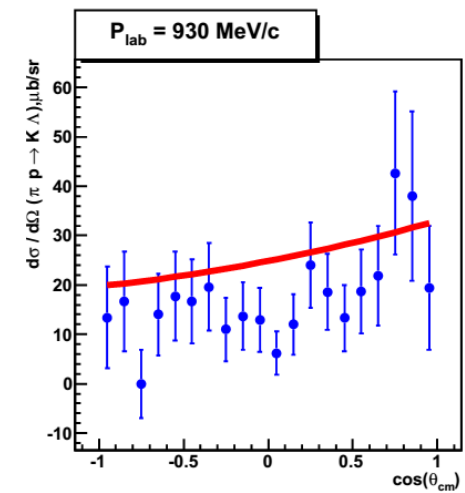


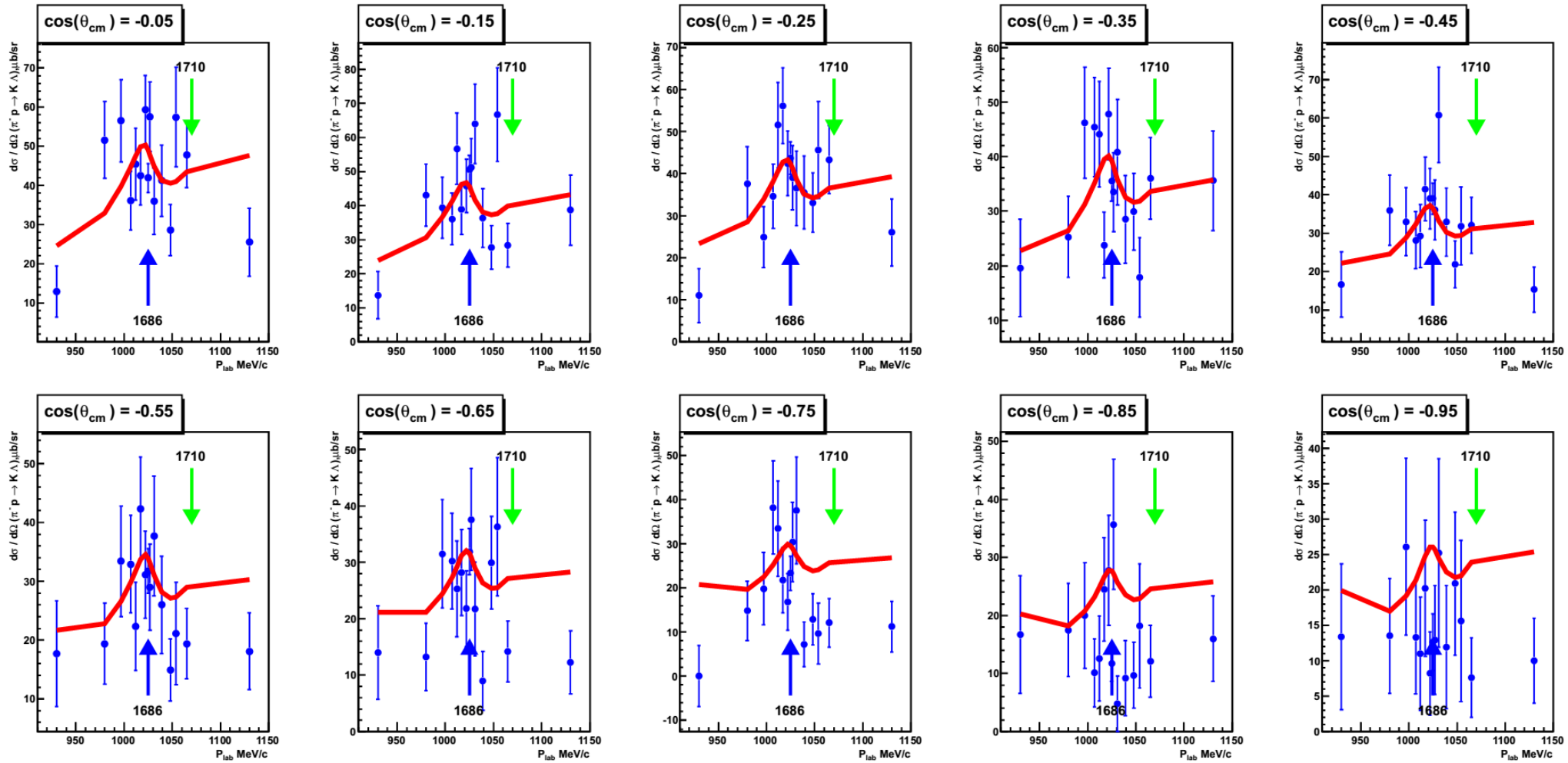


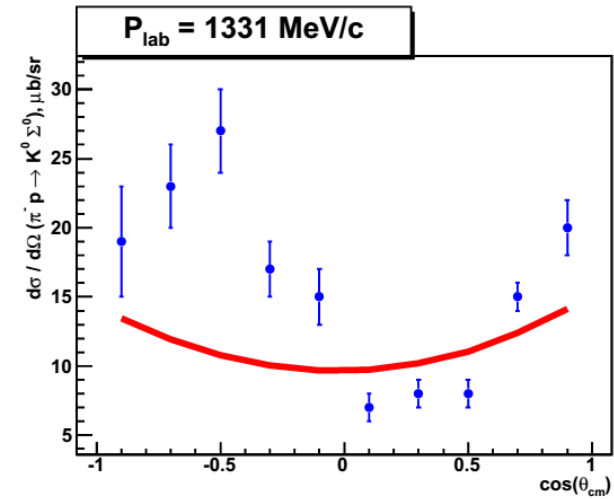
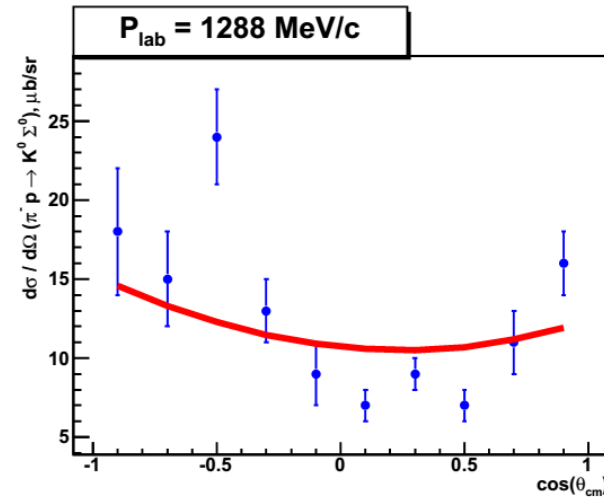
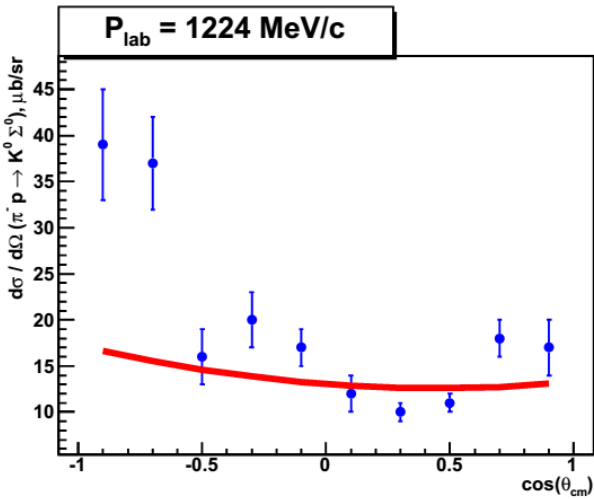
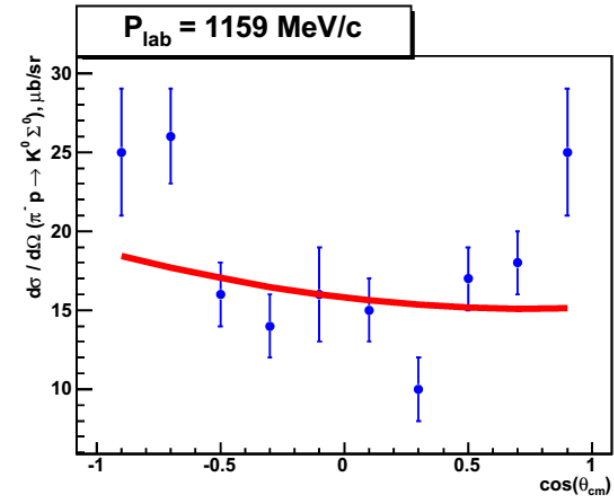
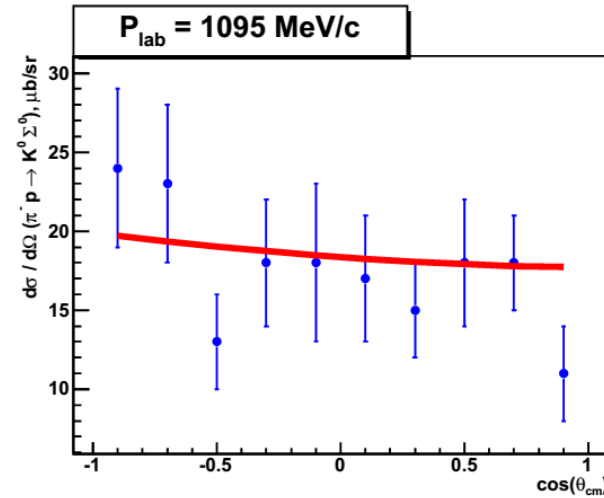
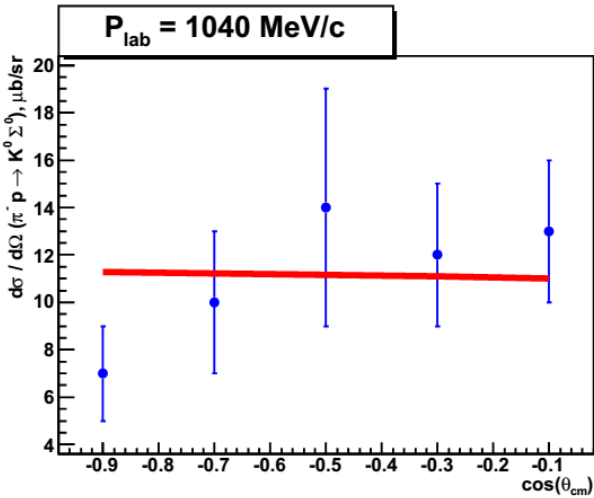


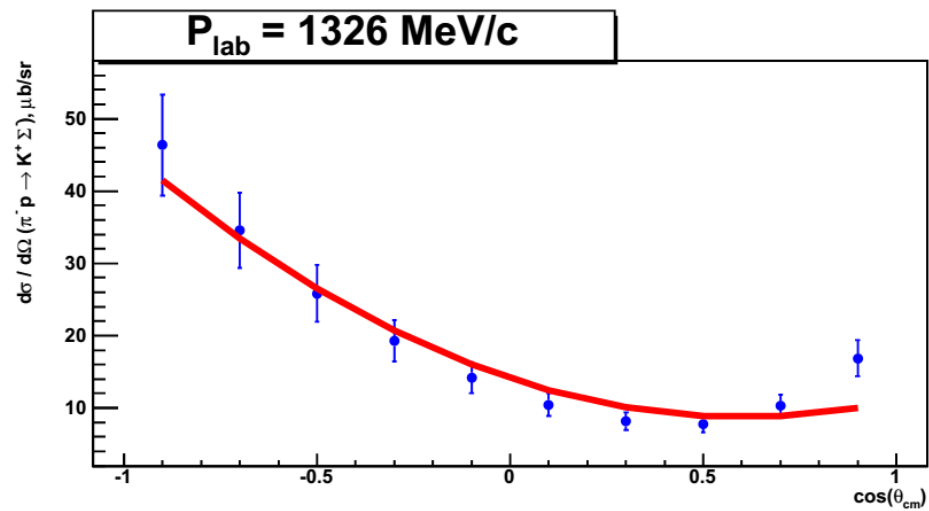
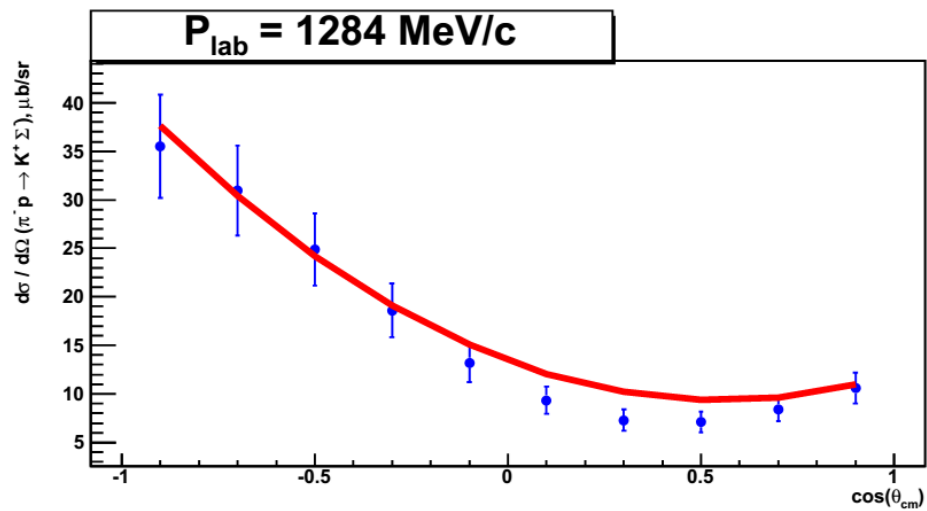
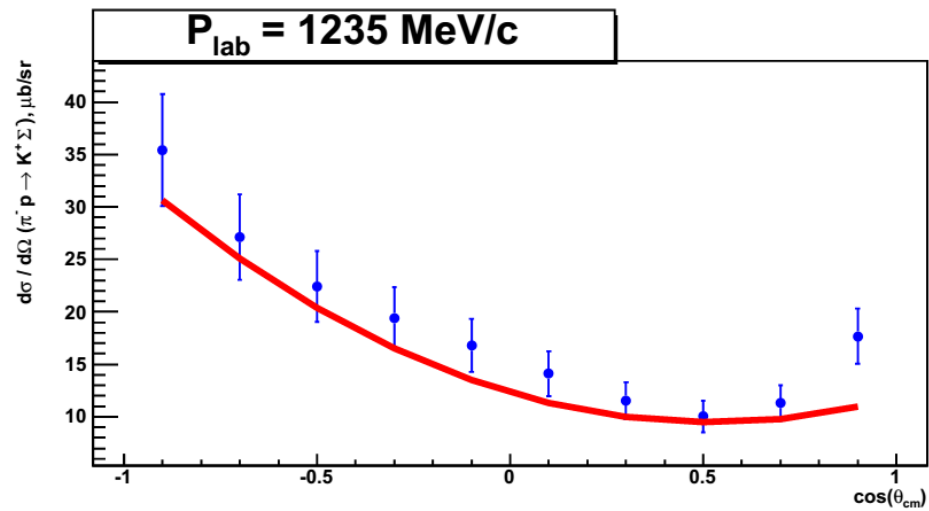
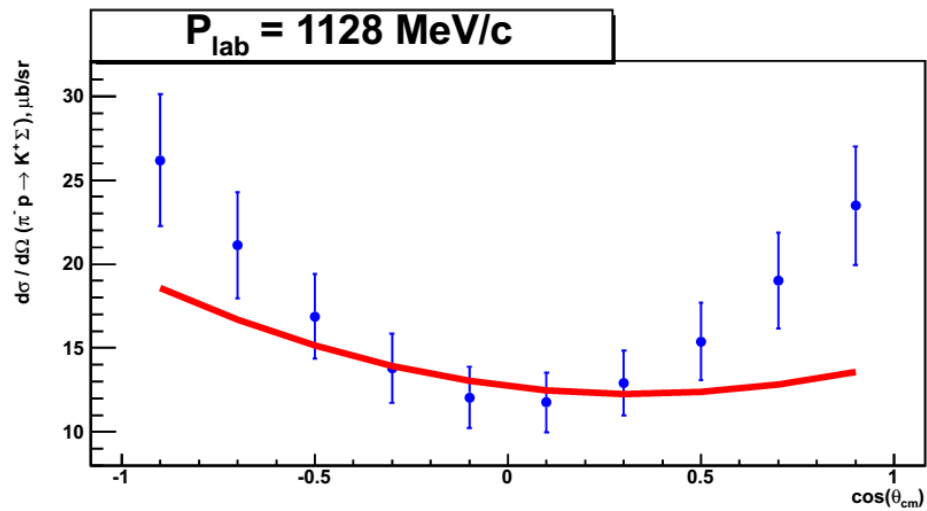


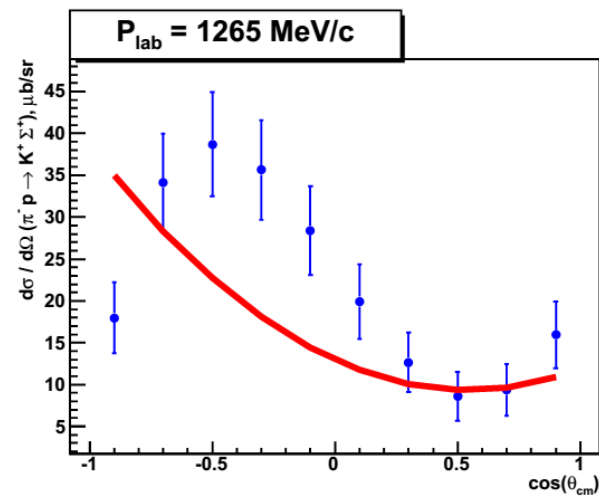
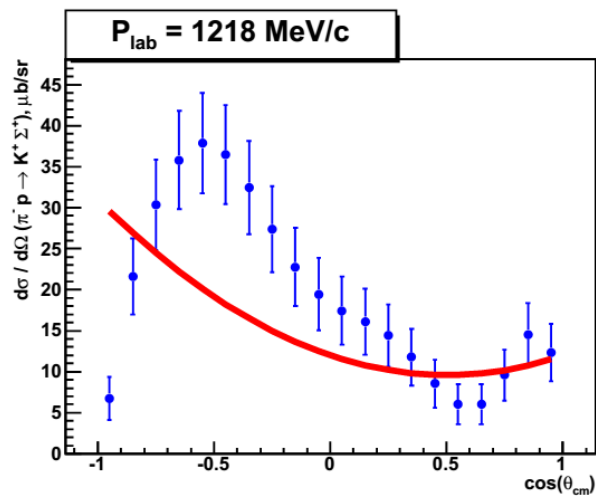
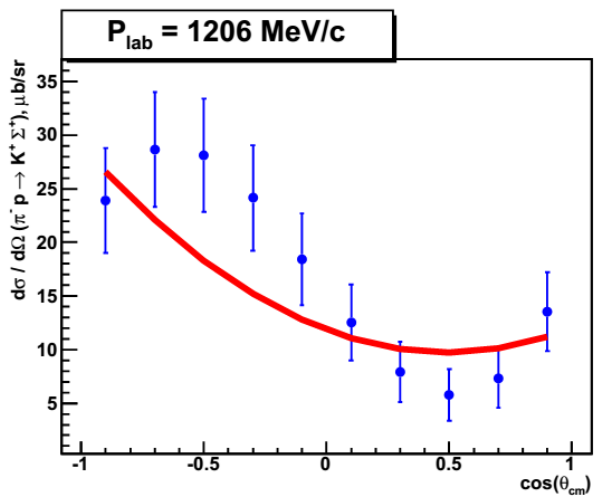
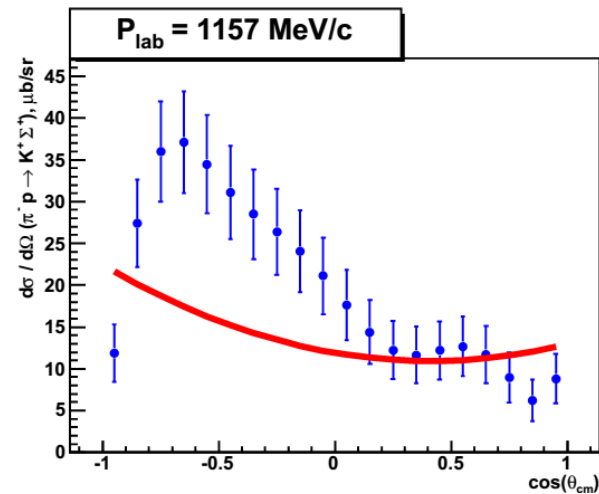
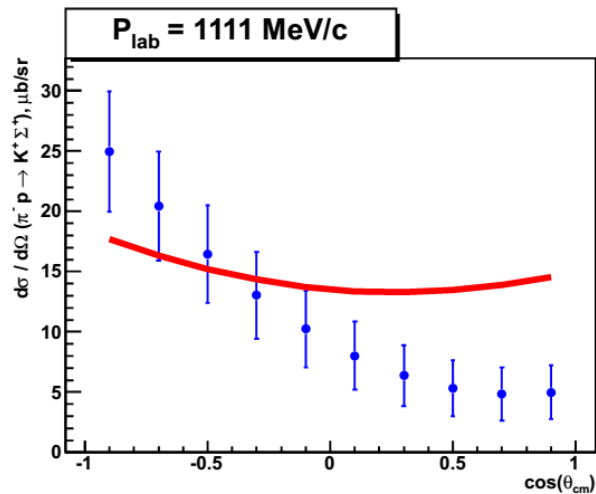
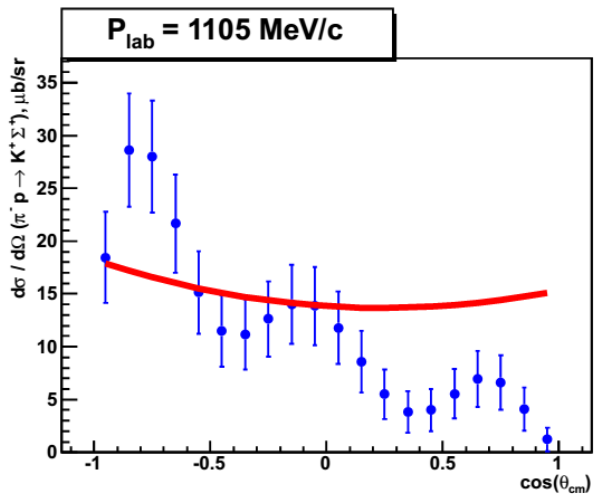












Very preliminary



S11

$$M=1686(1.5)$$

$$\Gamma_{tot}=18.2 \text{ MeV}$$

$$\Gamma_{el}=0.1 \text{ MeV}$$

$$\Gamma_{2\pi}=10.0 \text{ MeV}$$

$$\Gamma_{\eta n}=8.0 \text{ MeV}$$

$$\Gamma_{K\Lambda}=0.1 \text{ MeV}$$

P11

$$M=1710(2.0)$$

$$\Gamma_{tot}=25.0 \text{ MeV}$$

$$\Gamma_{el}=0.25 \text{ MeV}$$

$$\Gamma_{2\pi}=10.0 \text{ MeV}$$

$$\Gamma_{\eta n}=5.1 \text{ MeV}$$

$$\Gamma_{K\Lambda}=0.25 \text{ MeV}$$

$$\Gamma_{K\Sigma}=9.0 \text{ MeV}$$



$M=1686 \quad S11 \rightarrow P11 \rightarrow \chi^2 \uparrow 15\%$

$M=1710 \quad P11 \rightarrow S11 \rightarrow \chi^2 \uparrow 25\%$

Another explanations (for η photoproduction)

1. Interference effects.

Interference of well-known resonances

Interference of $S11(1650)$ and $P11(1710)$.

V. Shklyar, H. Lenske, U. Mosel, PLB650 (2007) 172 (Giessen group)

Interference of $S11(1535)$ and $S11(1650)$.

A. Anisovich et al. EPJA 41, 13 (2009) (Bonn-Gatchina group);

We not found such solution

2. Cusp effect

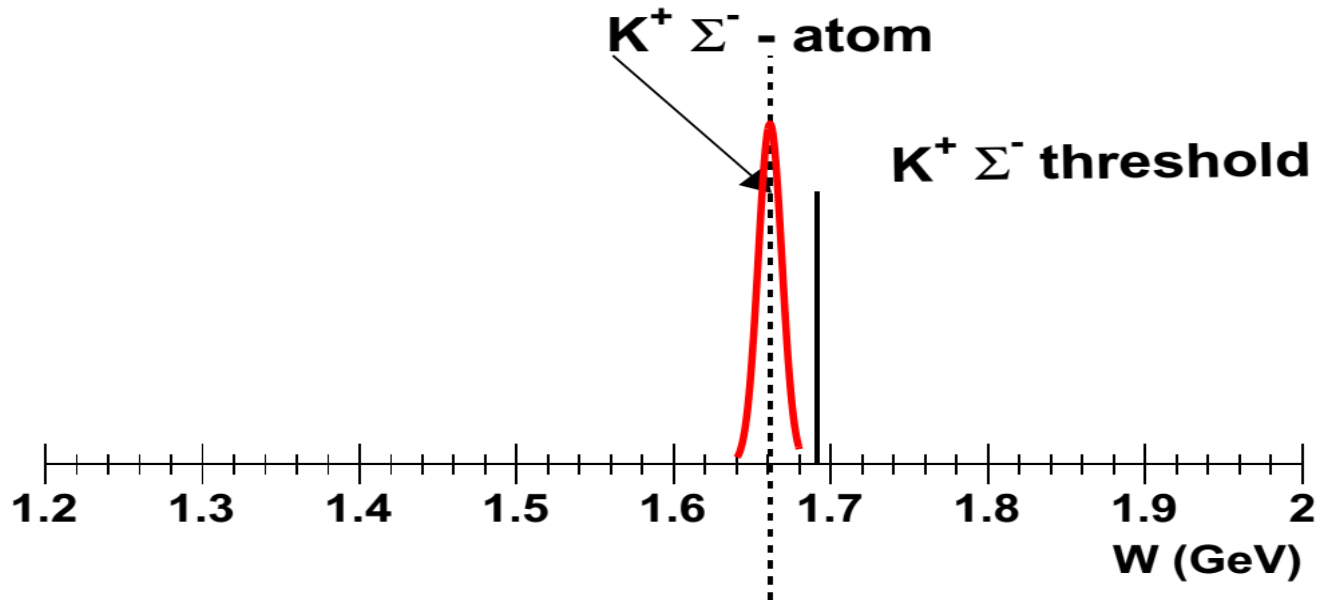
M. Doring, K. Nakayama, PLB B683:145 (2010)

We are working on this possibility



The electromagnetic effects.

$\pi^- p$ elastic , $\gamma n \rightarrow \eta n$



Experimental check:

$\pi^- p \rightarrow$ elastic, ηn , $k \Lambda$.

Isospin symmetry

$\pi^+ n \rightarrow$ elastic, ηn , $k \Lambda$.



Благодарю за внимание!