

Extra Physics with an Atomic Beam Source and a Lamb-Shift Polarimeter

22.06.2010

by Ralf Engels

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Introduction

1.) PIT@ANKE

2.) Bound β decay

3.) Precision Spectroscopy of H and D

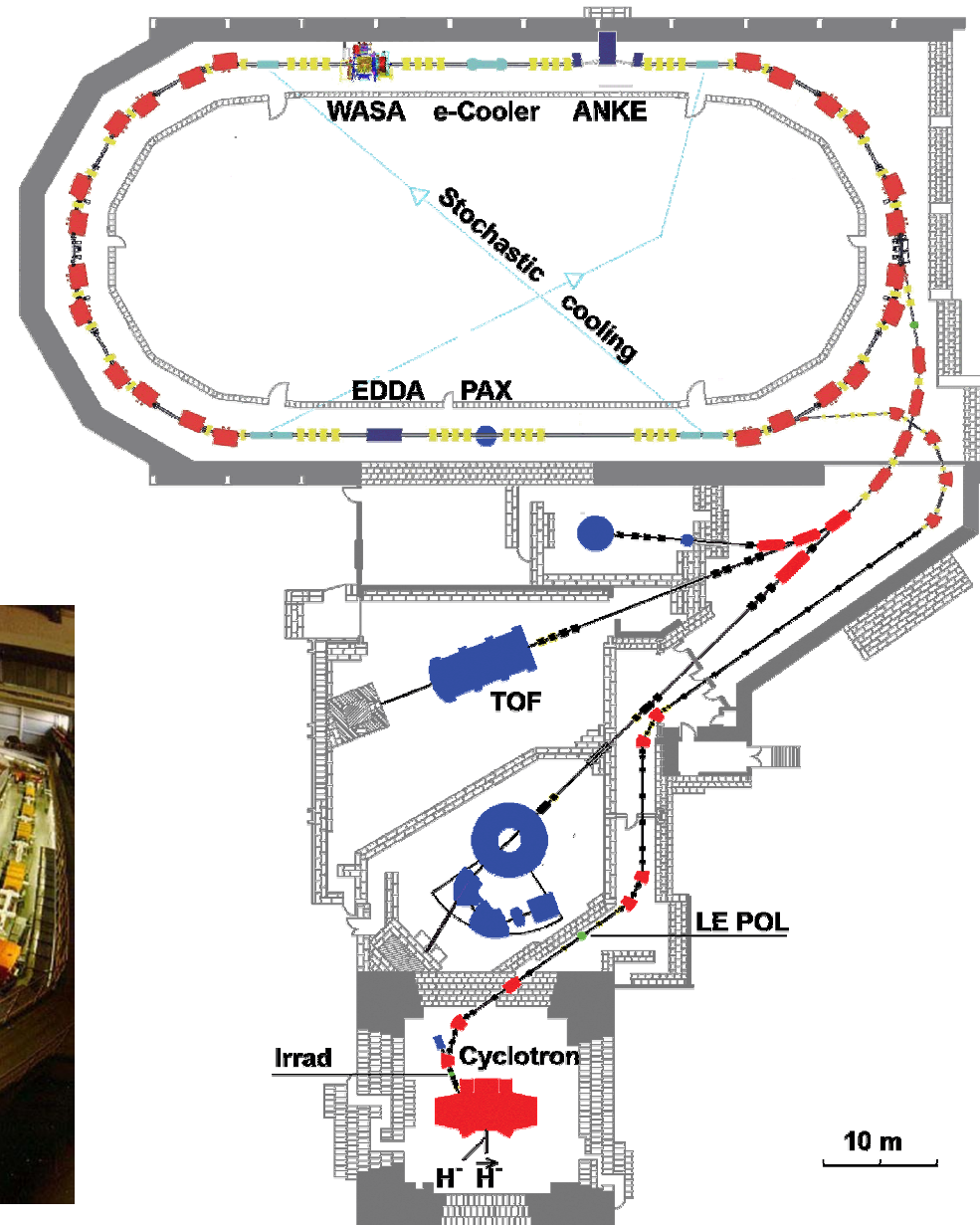
4.) Polarized Molecules

5.) PAX (Polarized Antiproton Experiment)

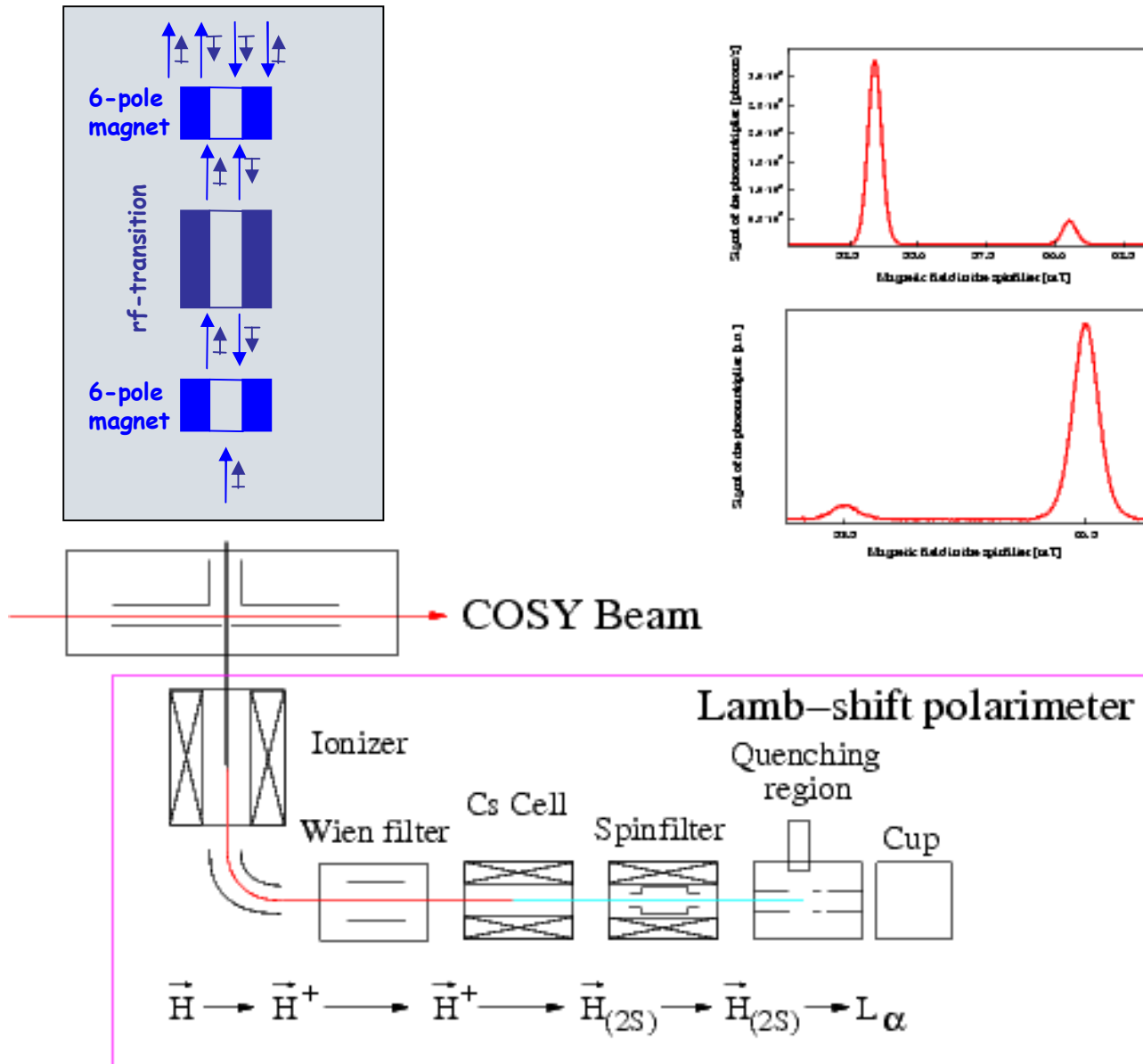
$$p, \vec{p}, d, \vec{d}$$

with momenta up to 3.7 GeV/c

- **internal experiments** – with the circulating beam
- **external experiments** – with the extracted beam



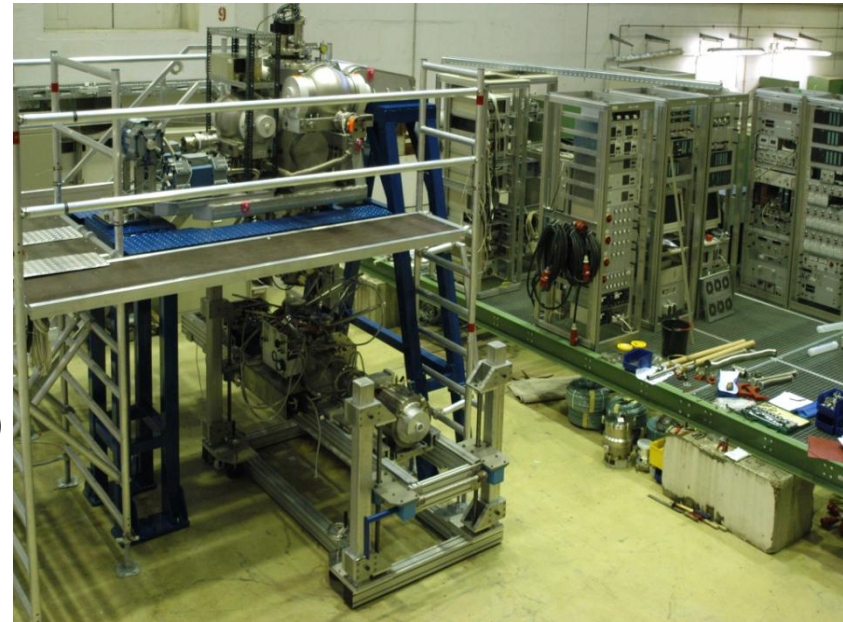
ABS and Lamb-shift polarimeter



PIT @ ANKE/COSY

Main parts of a PIT:

- **Atomic Beam Source**
 - Target gas
hydrogen or **deuterium**
 - H beam intensity (2 hyperfine states)
 $8.2 \cdot 10^{16}$ atoms / s
 - Beam size at the interaction point
 $\sigma = 2.85 \pm 0.42$ mm
 - Polarization for hydrogen atoms
 $P_Z = + 0.89 \pm 0.01$
 $P_Z = - 0.96 \pm 0.01$
- **Lamb-Shift Polarimeter**
- **Storage Cell**



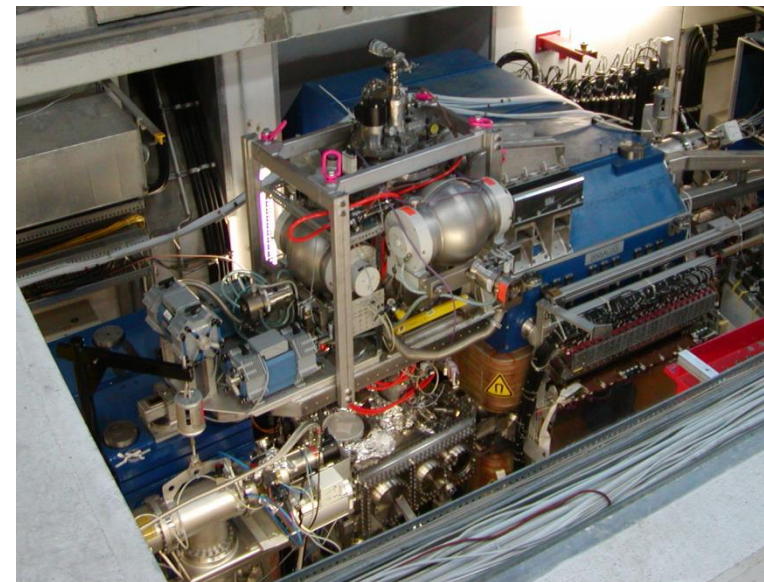
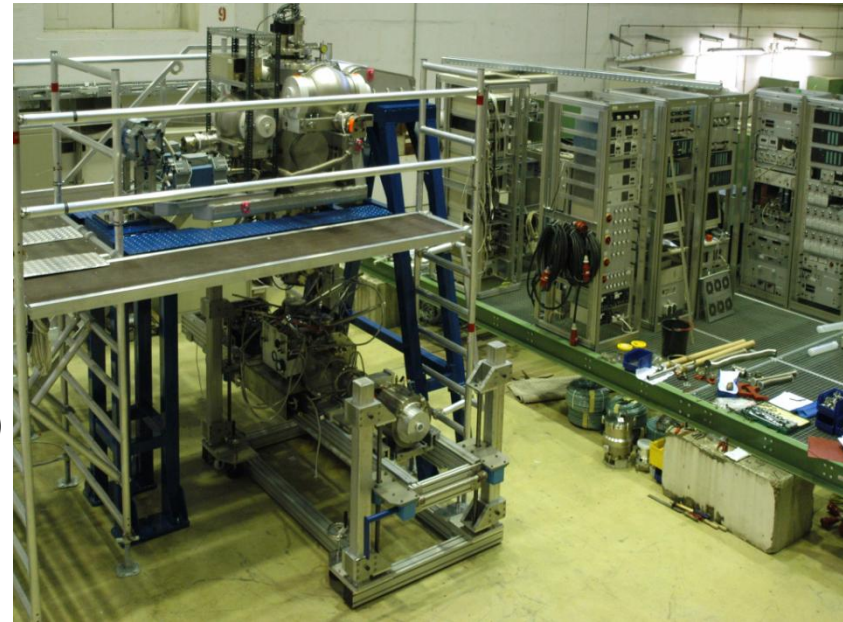
PIT @ ANKE/COSY



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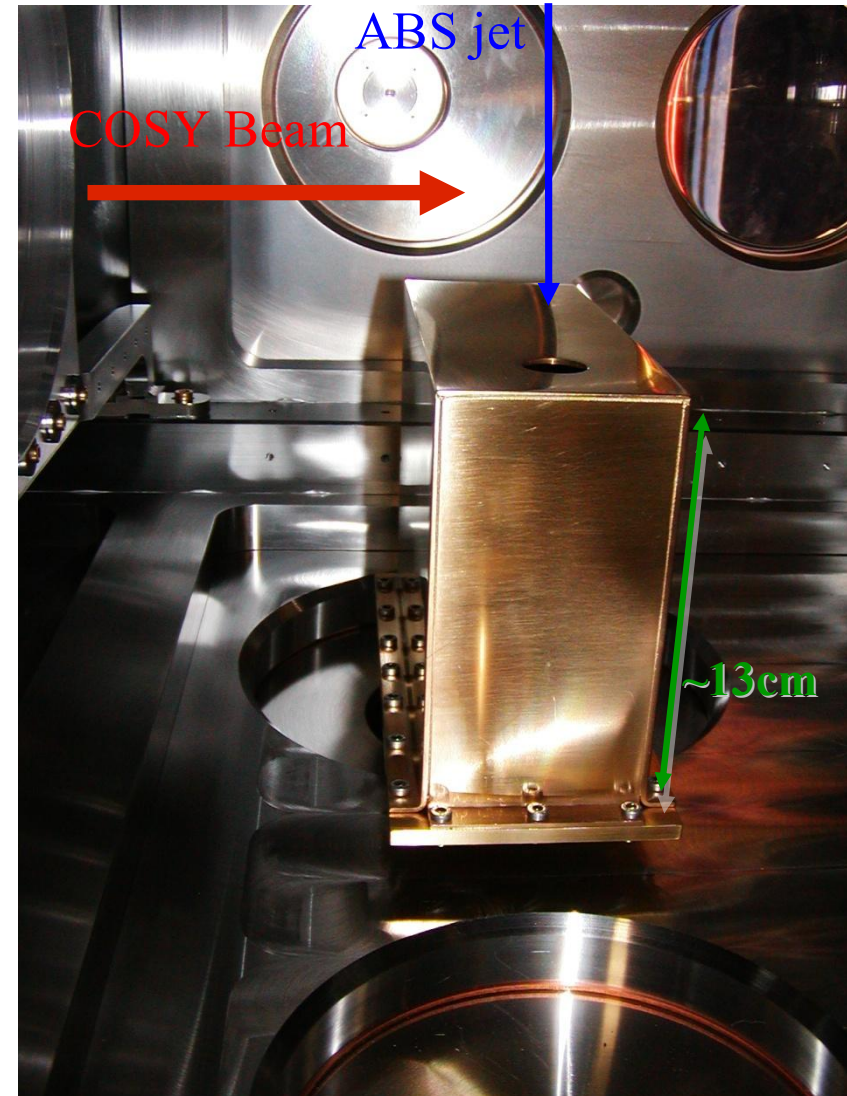
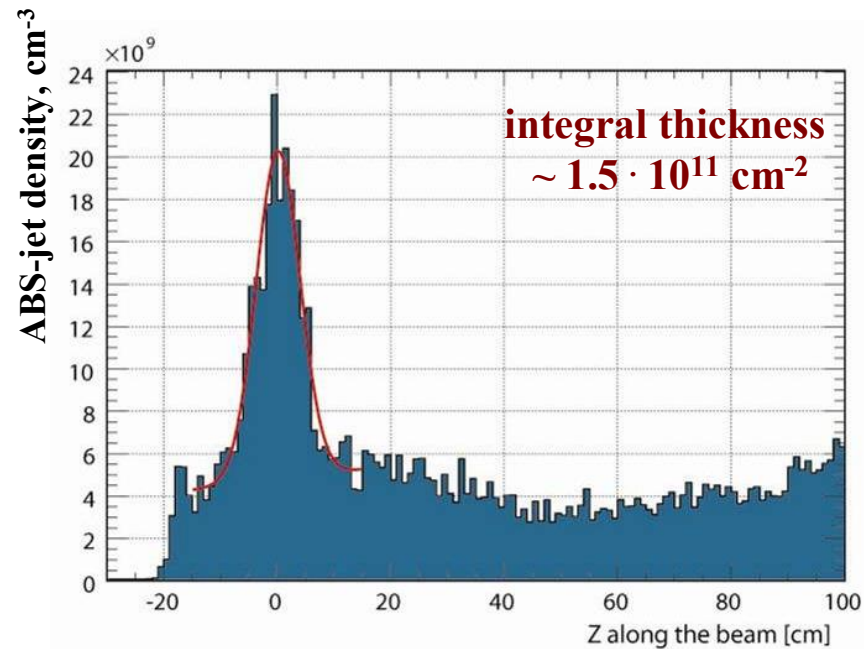
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The ABS jet target

Measured pressures in the target chamber, mbar

	Without ABS jet	With ABS jet
Without catcher	$4.0 \cdot 10^{-9}$	$3.0 \cdot 10^{-7}$
With catcher	$4.0 \cdot 10^{-9}$	$3.7 \cdot 10^{-8}$



ABS beam catcher

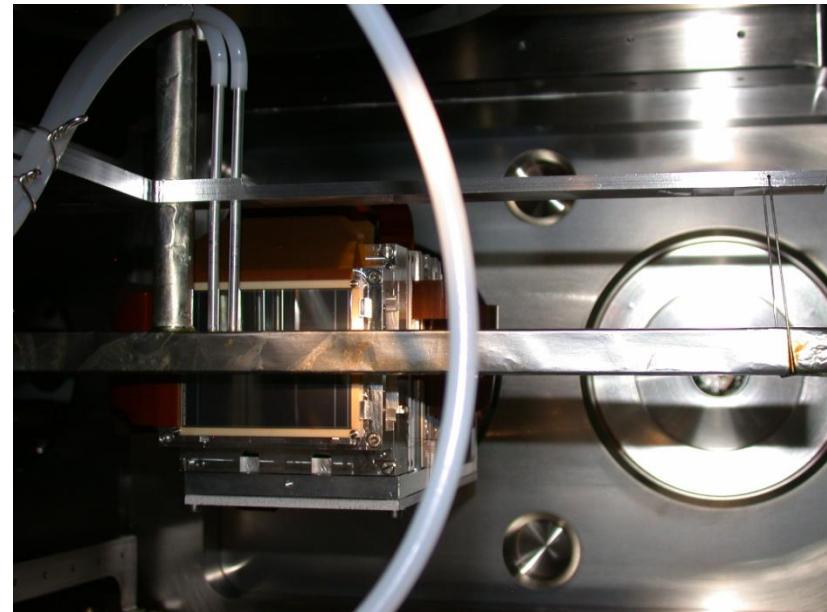
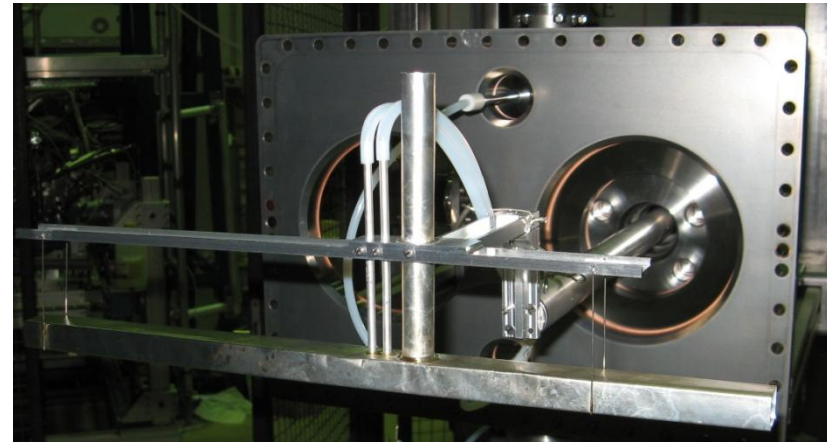
Storage cell: Final setup

Tools for the experiment

- New storage cell & support (20•20•390 mm)
 - > high target density
 - > unpolarized gas feeding system

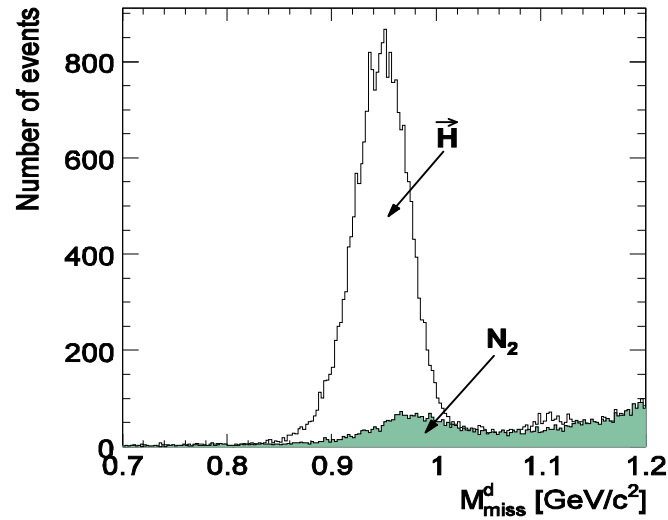
- LSP below the target chamber
 - > online measurement of the ABS beam polarization

- Silicon Tracking Telescope (STT)
 - > measurement of spectator protons nearby the storage cell center

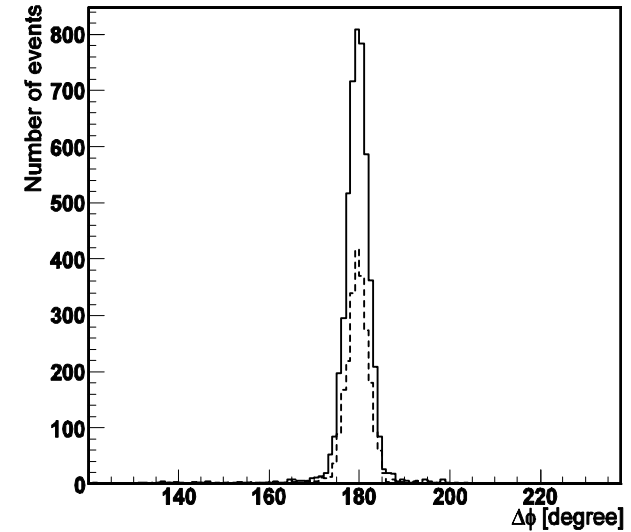
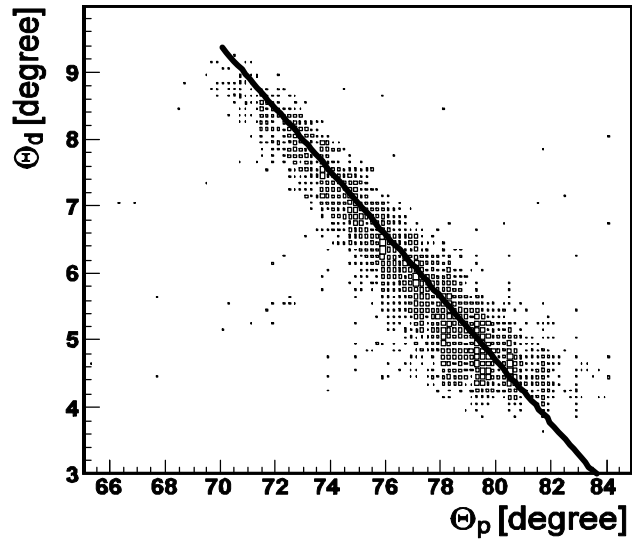
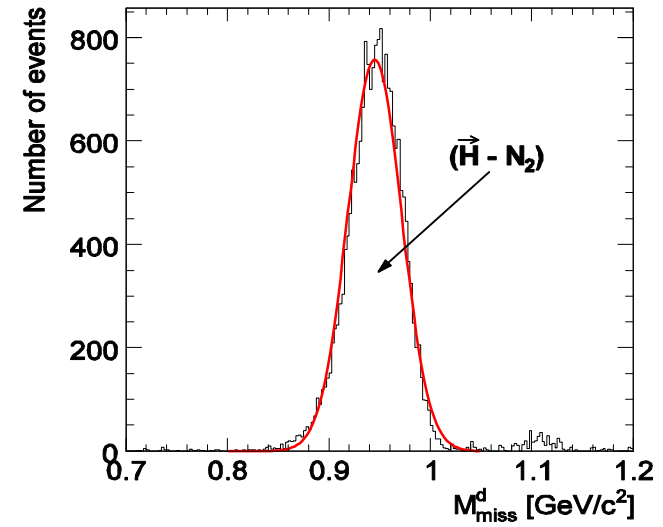


First Double Polarized Experiment: $\vec{dp} \rightarrow \vec{dp}$

\vec{dp} , $T_d = 1.2$ GeV, target H (N_2) gas

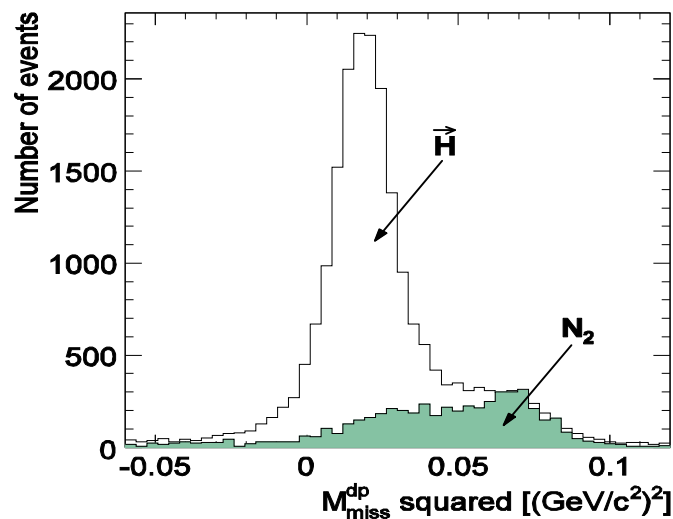


$dp \rightarrow dp$



First Double Polarized Experiment: $\vec{d}\vec{p}$

$\vec{d}\vec{p}$, $T_d = 1.2$ GeV, target H (N_2) gas

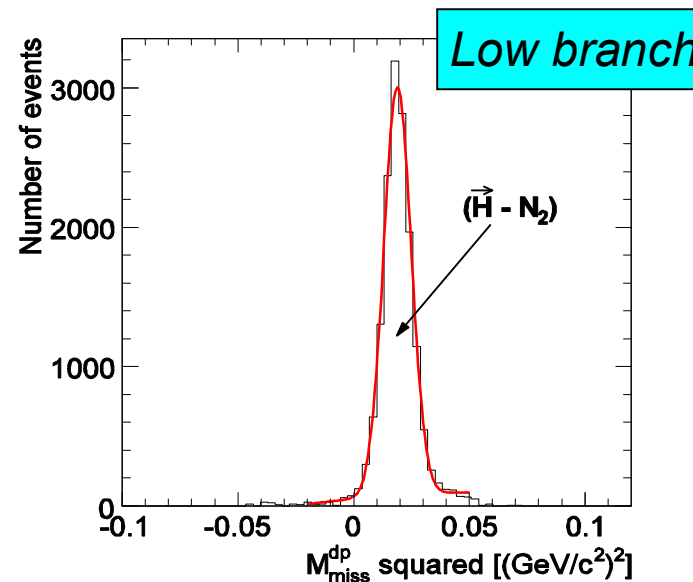
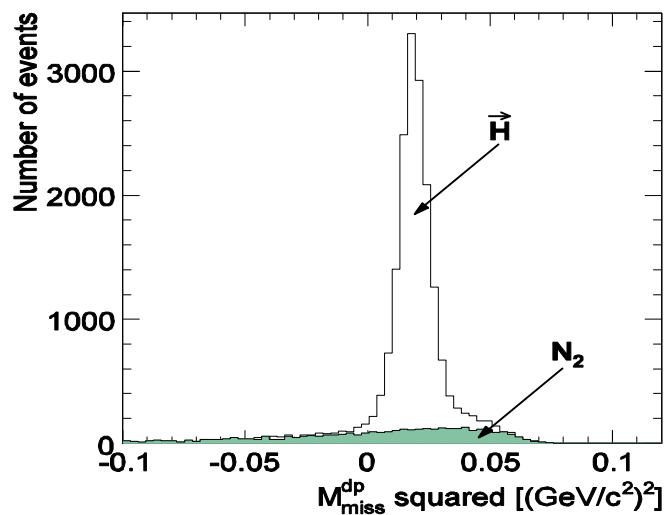
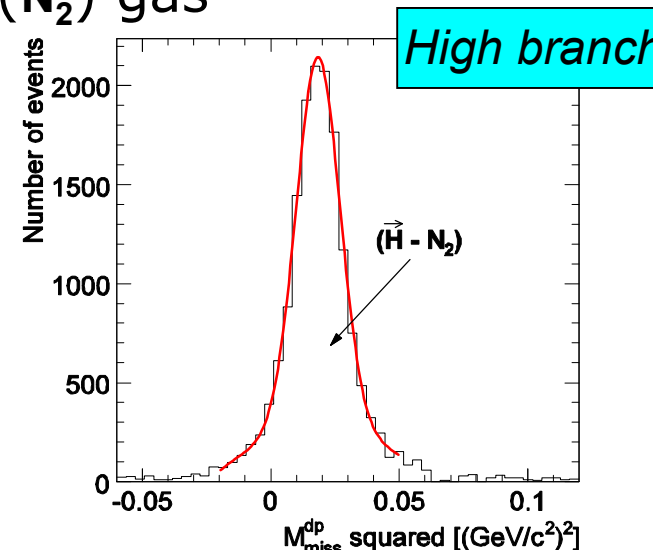


$dp \rightarrow (dp_{sp}) \pi^0$

quasi-free

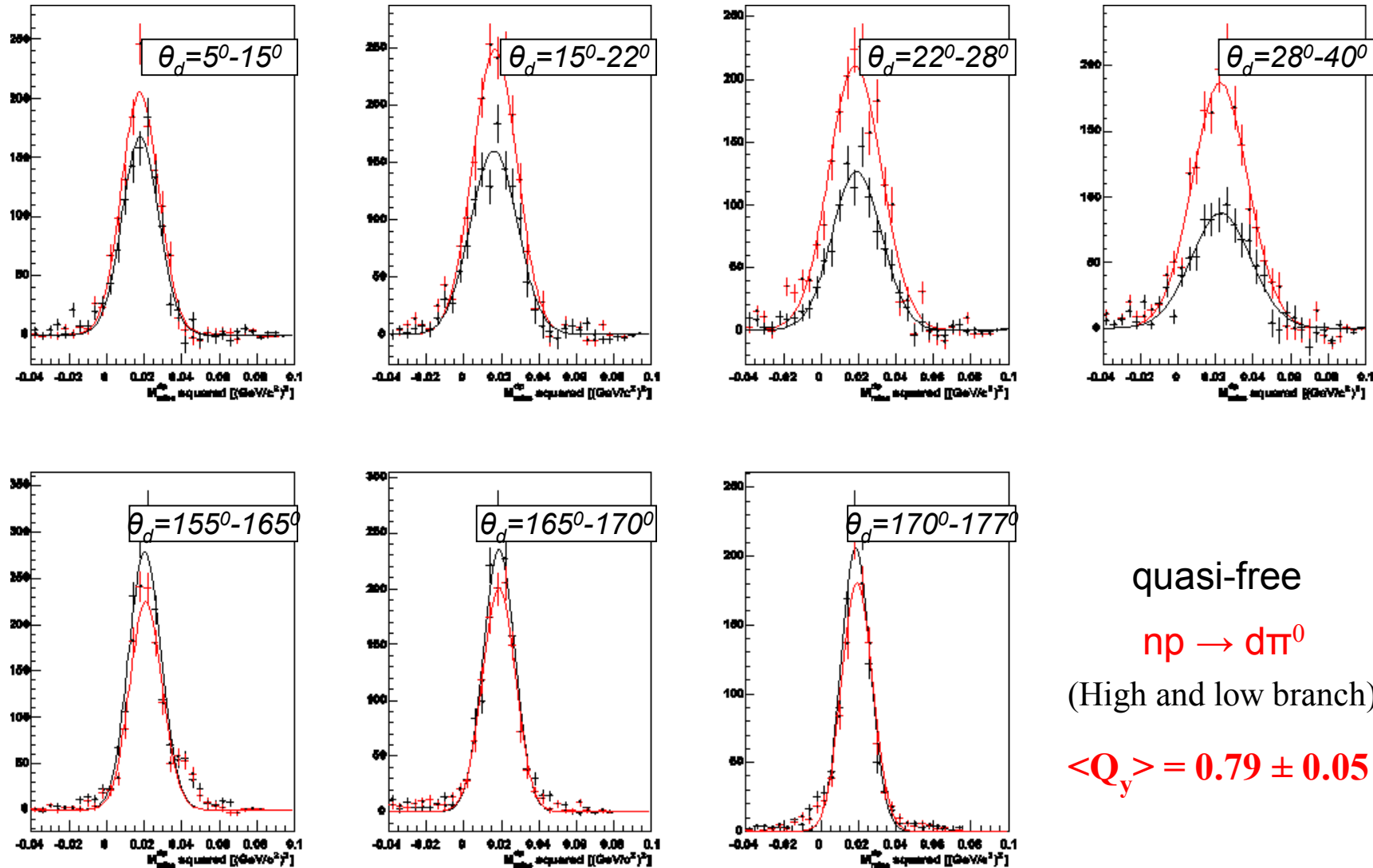
$np \rightarrow d \pi^0$

(High & Low branch)



Polarization Measurement of the Target

\vec{dp} , $T_d = 1.2$ GeV, polarized H gas



quasi-free

$np \rightarrow d\pi^0$

(High and low branch)

$\langle Q_y \rangle = 0.79 \pm 0.05$

More Experiments with ABS and LSP

2.) Bound β Decay

3.) Precision Spectroscopy of Hydrogen / Deuterium

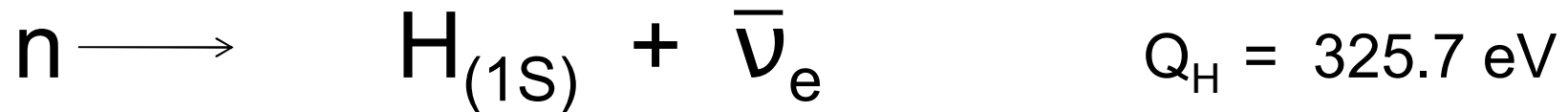
4.) Polarized H₂ Molecules

5.) PAX (**P**olarized **A**ntiproton **E**xperiment)

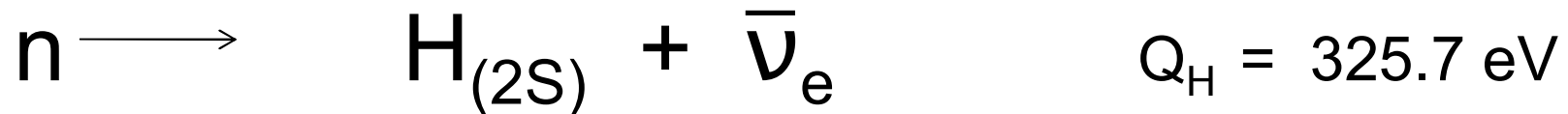
The Bound- β Decay (Tech. Uni. Munich)



Efficiency: $4 \cdot 10^{-6}$



Efficiency: $\sim 10^{-1}$



L.L. Nemenov,
Sov. J. Nucl. Phys **31** (1980)

The Neutron Decay

Helicity of the Antineutrino: right-handedness

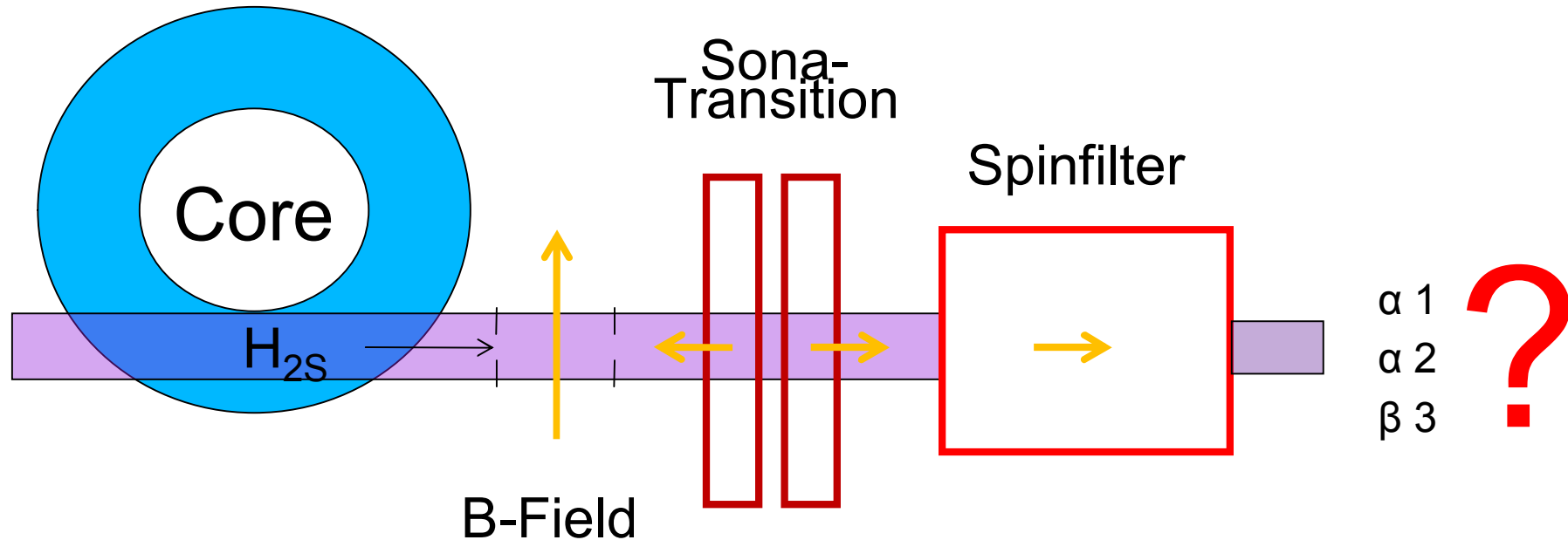
$\bar{\nu}$	n	p	e ⁻	W_i (%)	F	mF	HFS
←	←	←	→	44.14	0,1	0	α_2, β_4
←	←	→	←	55.24	0,1	0	β_4, α_2
←	→	→	→	0.62	1	1	α_1
→	←	←	←	0	1	-1	β_3
→	→	→	←	0	0,1	0	β_4, α_2
→	→	←	→	0	0,1	0	α_2, β_4

→ left handed admixtures ?

→ scalar or tensor contributions to the weak force ?

The Neutron Decay

Reactor: FRM II



Problem:

How to register single metastable hydrogen atoms?

(Count rates: H_{1S} : 3 s^{-1} / H_{2S} : 0.3 s^{-1} / $H_{2S(\alpha 2)}$: $\leq 0.1 \text{ s}^{-1}$)

The Neutron Decay

1. Selective Ionization with 2 Laser (Hänsch et al.)

Advantage: - Efficiency ~50 % expected

Problems: - Not used before

- Very expensive and difficult (Resonator: 20 kW !!!)

- Background free ?

2. Lyman- α detection with PM

Advantage: - Used before (PM: sensitive to 110-135 nm)

Problems: - QE: ~10 % / LSP: Efficiency: 10^{-3}

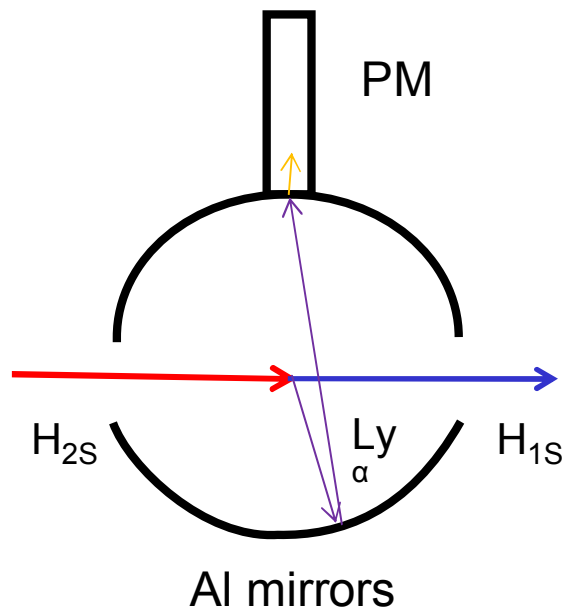
- Ideas: a.) Mirrors for 121 nm \rightarrow 5 % possible ?

b.) Photoeffect-Chamber \rightarrow 80 % possible ?

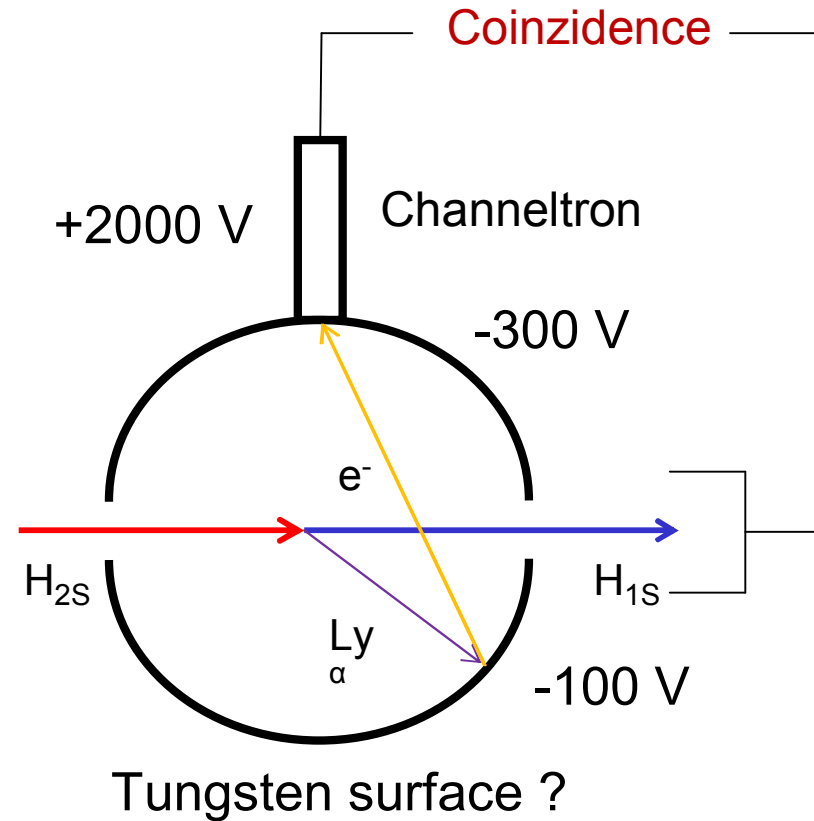
The Neutron Decay

2. Lyman- α detection with PM

Mirrors



Photoeffect



The Neutron Decay

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- Ideas: a.) Mirrors for 121 nm \rightarrow 5 % possible ?
b.) Photoeffect-Chamber \rightarrow 80 % possible ?

3. Charge Exchange at Ar (K)

Advantage: - Used before (pol. Lamb-shift source)

Problems: - Efficiency: ~10 %, can not be increased much

- Background free ? \rightarrow Energy-separation necessary

The Neutron Decay

3. Charge Exchange at Ar (K)

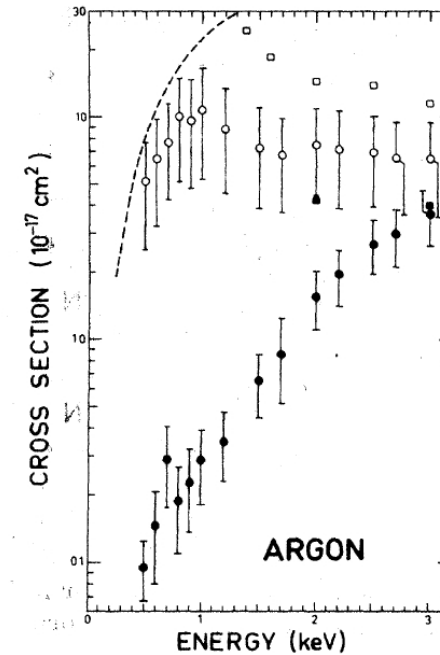
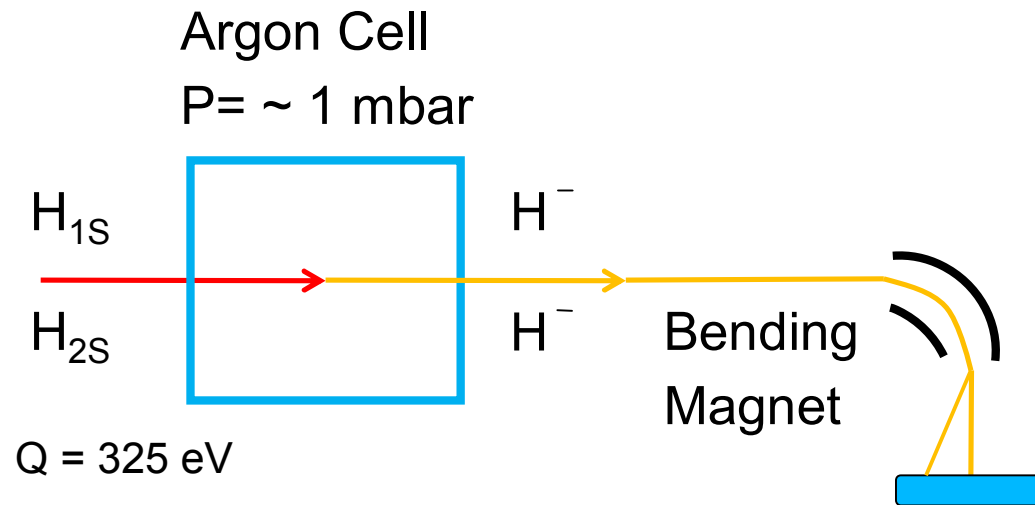
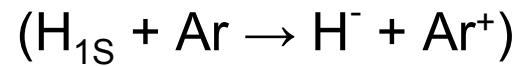
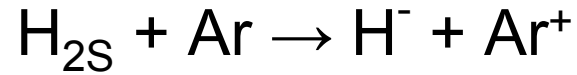
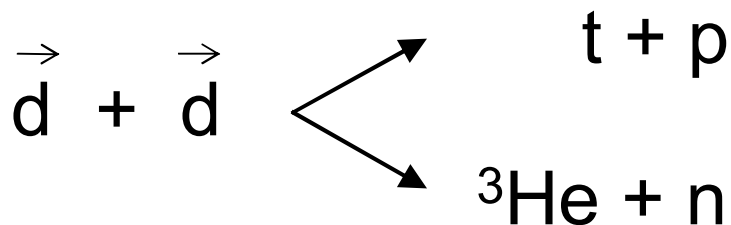
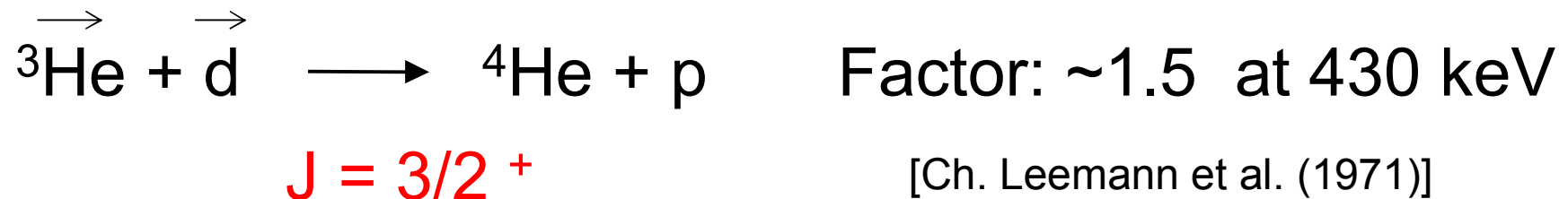


FIG. 10. Electron-capture cross sections for H(1^2S) and H(2^2S) in argon (55-mrad detector's acceptance angle). σ_g : \bullet , present work; \blacksquare , Williams.⁵ $\sigma_{m\pi}$: \circ , present work; \square , Dose and Gunz⁷ recalibrated; ---, theoretical calculation by Olson.¹⁴

Paper is in preparation

Double Polarized Fusion

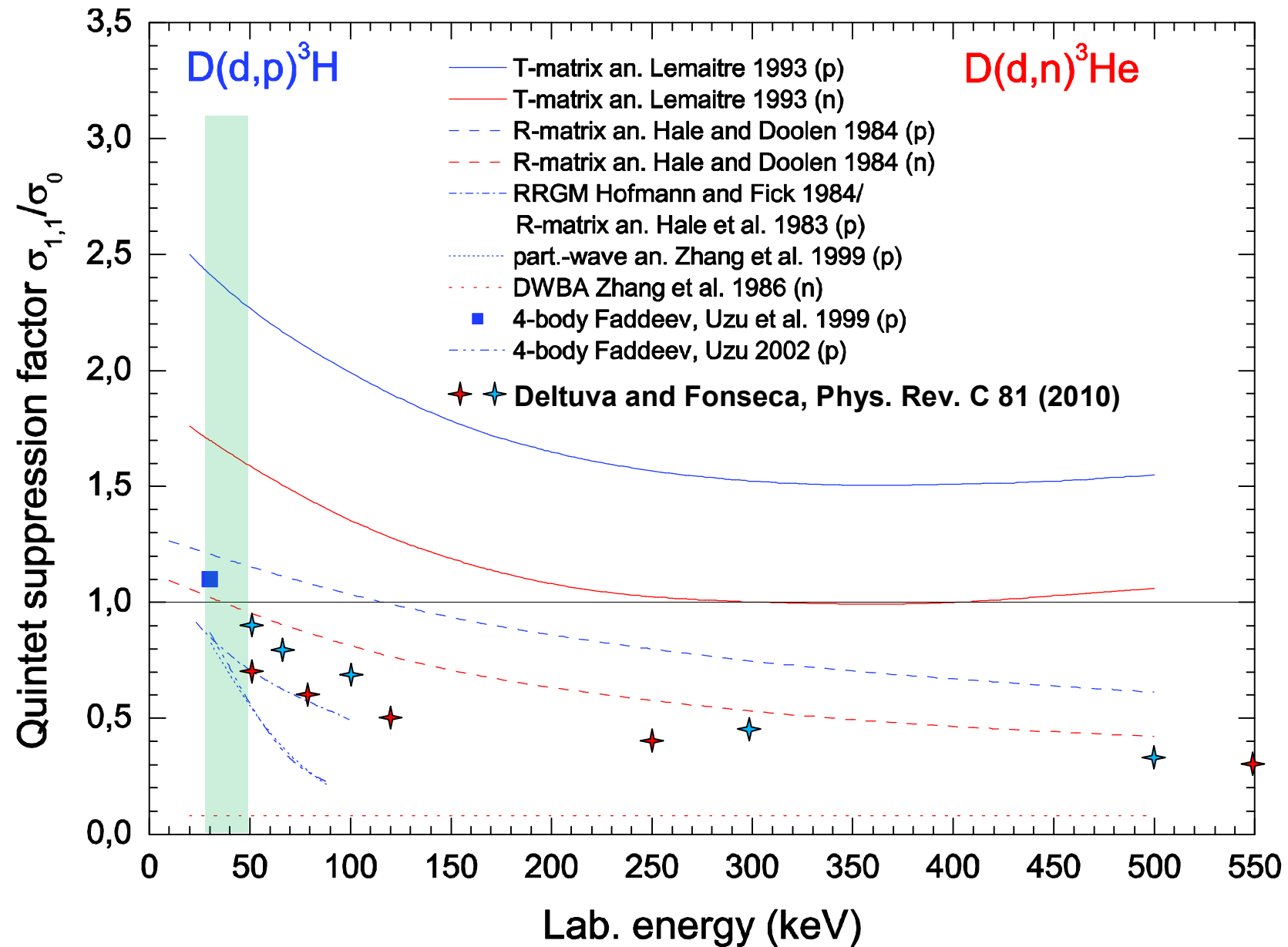
Can the total cross section of the fusion reactions be increased by using polarized particles ?



Can cross sections be increased ?
 Can neutrons be suppressed ?
 Can the trajectories of the neutrons be controlled?

H. Paetz gen. Schieck, Eur. Phys. J. A **44**, 321-354 (2010)

The Quintet Suppression Factor



The Formula

$$\begin{aligned}
 \sigma(\Theta, \Phi) = \sigma_0(\Theta) \{ & 1 + \frac{3}{2} [A_y^{(b)}(\Theta)p_y + A_y^{(t)}q_y] + \frac{1}{2} [A_{zz}^{(b)}(\Theta)p_{zz} + A_{zz}^{(t)}(\Theta)q_{zz}] \\
 & + \frac{1}{6} [A_{xx-yy}^{(b)}(\Theta)p_{xx-yy} + A_{xx-yy}^{(t)}(\Theta)q_{xx-yy}] \\
 & + \frac{2}{3} [A_{xz}^{(b)}(\Theta)p_{xz} + A_{xz}^{(t)}(\Theta)q_{xz}] \\
 & + \frac{9}{4} [C_{y,y}(\Theta)p_yq_y + C_{x,x}(\Theta)p_xq_x + C_{x,z}(\Theta)p_xq_z \\
 & + C_{z,x}(\Theta)p_zq_x + C_{z,z}(\Theta)p_zq_z] \\
 & + \frac{3}{4} [C_{y,zz}(\Theta)p_yq_{zz} + C_{zz,y}(\Theta)p_{zz}q_y] \\
 & + C_{y,xz}(\Theta)p_yq_{xz} + C_{xz,y}(\Theta)p_{xz}q_y + C_{x,yz}(\Theta)p_xq_{yz} \\
 & + C_{y,zx}(\Theta)p_yq_{zx} + C_{z,yz}(\Theta)p_zq_{yz} + C_{y,z,z}(\Theta)p_{yz}q_z \\
 & + \frac{1}{4} [C_{y,xx-yy}(\Theta)p_yq_{xx-yy} + C_{xx-yy,y}(\Theta)p_{xx-yy}q_y \\
 & + C_{zz,zz}(\Theta)p_{zz}q_{zz}] \\
 & + \frac{1}{3} [C_{zz,xz}(\Theta)p_{zz}q_{xz} + C_{xz,zz}(\Theta)p_{xz}q_{zz}] \\
 & + \frac{1}{12} [C_{zz,xx-yy}(\Theta)p_{zz}q_{xx-yy} + C_{xx-yy,zz}(\Theta)p_{xx-yy}q_{zz}] \\
 & + \frac{4}{9} [C_{xz,xz}(\Theta)p_{xz}q_{xz} + C_{yz,yz}(\Theta)p_{yz}q_{yz}] \\
 & + \frac{8}{9} [C_{xy,yz}(\Theta)p_{xy}q_{yz} + C_{yz,xy}(\Theta)p_{yz}q_{xy}] \\
 & + \frac{16}{9} C_{xy,xy}(\Theta)p_{xy}q_{xy} \\
 & + \frac{1}{9} [C_{xz,xx-yy}(\Theta)p_{xz}q_{xx-yy} + C_{xx-yy,xz}(\Theta)p_{xx-yy}q_{xz}] \\
 & + \frac{1}{36} C_{xx-yy,xx-yy}(\Theta)p_{xx-yy}q_{xx-yy} \\
 & + \frac{1}{2} [C_{x,xy}(\Theta)p_xq_{xy} + C_{xy,x}(\Theta)p_{xy}q_x + C_{z,xy}(\Theta)p_zq_{xy} \\
 & + C_{xy,z}(\Theta)p_{xy}q_z] \}
 \end{aligned}$$

Spins of both deuterons
are aligned:

Only $p_z(q_z)$ and $p_{zz}(q_{zz}) \neq 0$

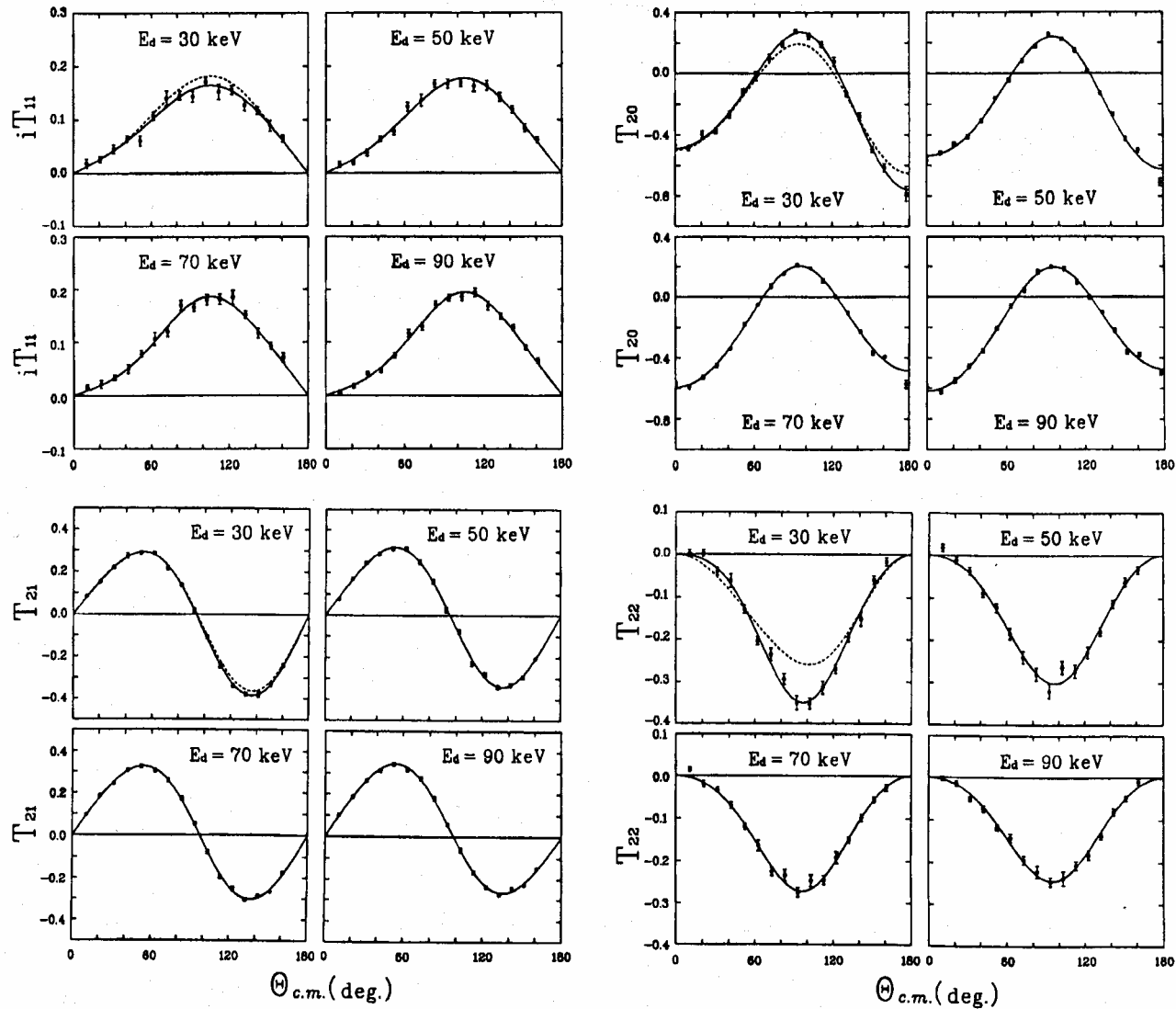
$$\sigma(\Theta, \Phi) = \sigma_0(\Theta) \left\{ 1 + \frac{3}{2} [A_{zz}^{(b)}(\Theta)p_{zz} + A_{zz}^{(t)}(\Theta)q_{zz}] + \frac{9}{4} C_{z,z}(\Theta)p_zq_z + \frac{1}{4} C_{zz,zz}(\Theta)p_{zz}q_{zz} \right\}$$

Only beam is polarized:
($p_{i,j} \neq 0, q_{i,j} = 0$)

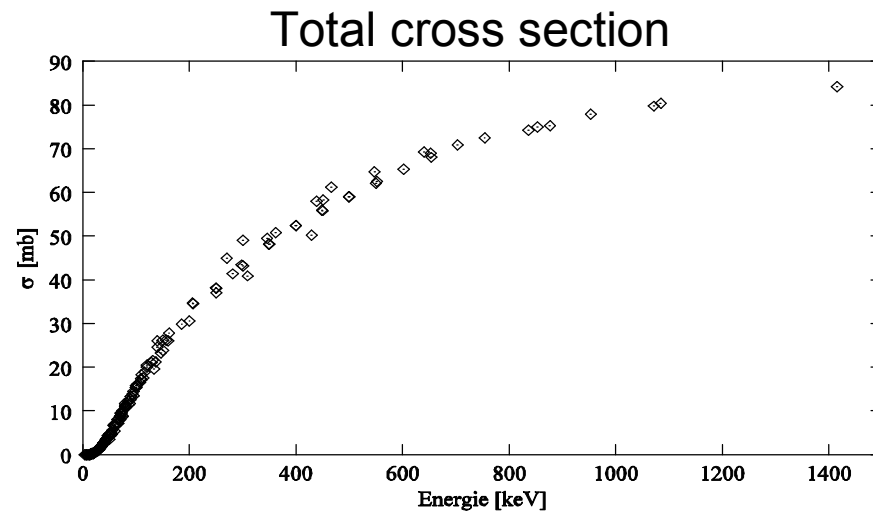
$$\begin{aligned}
 \sigma(\Theta, \Phi) = \sigma_0(\Theta) \cdot \{ & 1 + 3/2 A_y(\Theta) p_y \\
 & + 1/2 A_{xz}(\Theta) p_{xz} \\
 & + 1/6 A_{xx-yy}(\Theta) p_{xx-yy} \\
 & + 2/3 A_{zz}(\Theta) p_{zz} \}
 \end{aligned}$$

The Analysing Powers

Tagishi et al.; Phy. Rev. C **46** (1992) 1155-1158



The Experimental Setup



Q: 4.033 MeV

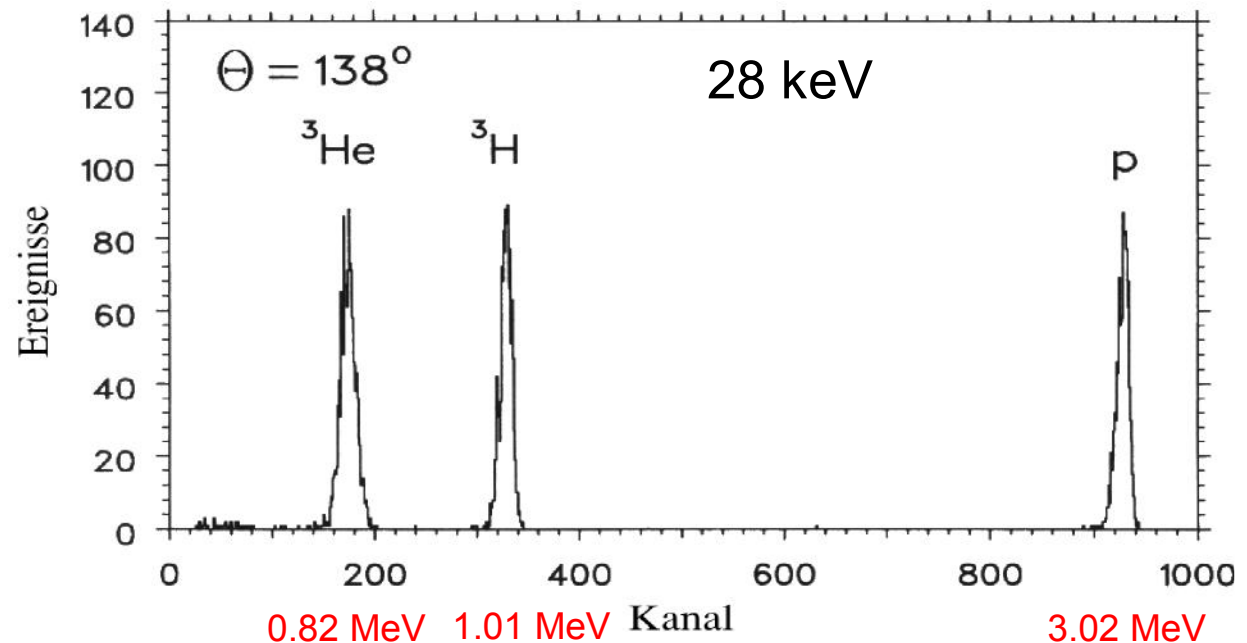


Q: 3.268 MeV

Becker et al.
Few Body Sys. **13** (1992)
[Analysing Powers]

Tagishi et al.
Phys.Rev. C **46** (1992)
[Analysing Powers]

Imig et al.
Phys.Rev. C **73** (2006)
[Spin-Transfer Koeff.]



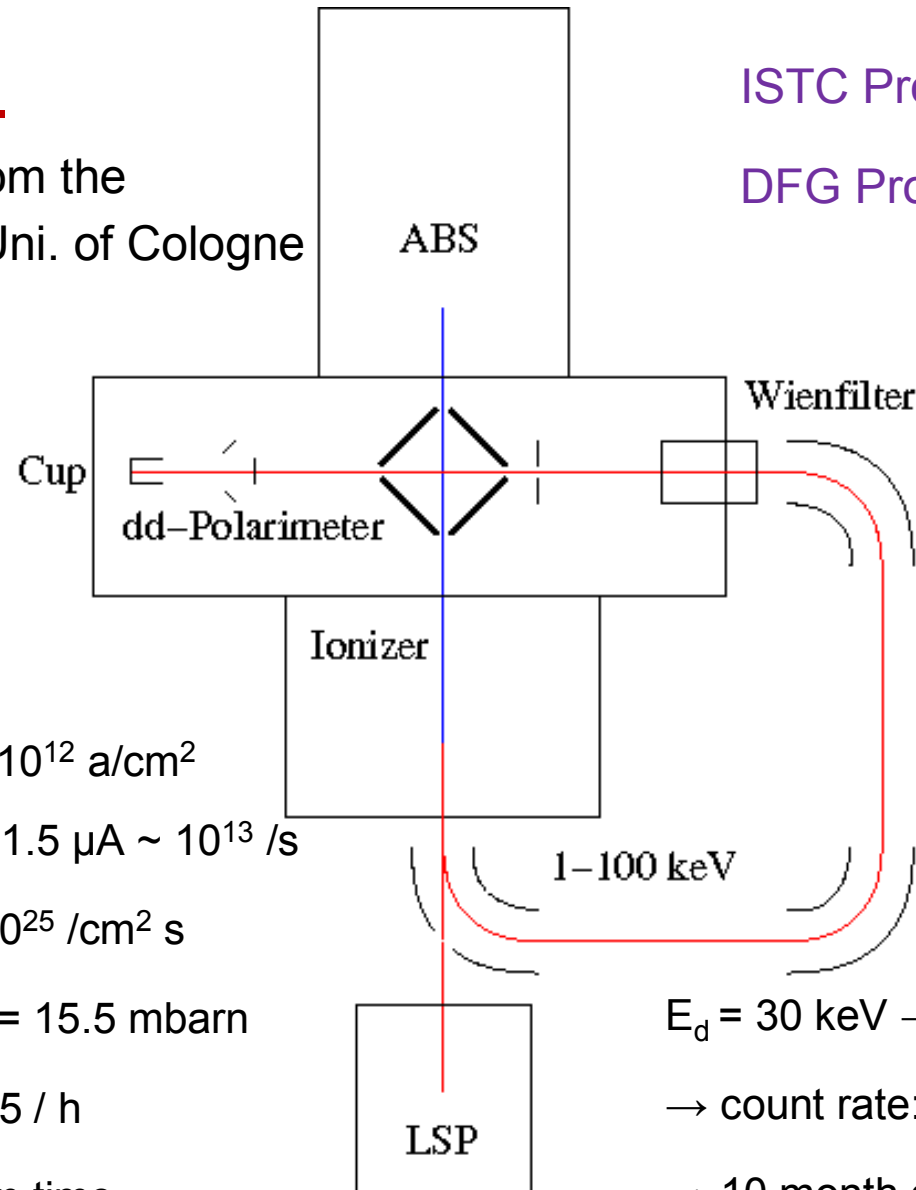
The Experimental Setup in St. Petersburg

1. Setup:

ABS and LSP from the
SAPIS Project, Uni. of Cologne

ISTC Project # 3881

DFG Project: EN 902/1-1



Target Density: $< 10^{12} \text{ a/cm}^2$

Beam Intensity: $> 1.5 \mu\text{A} \sim 10^{13} / \text{s}$

→ Luminosity: $\leq 10^{25} / \text{cm}^2 \text{ s}$

$E_d = 100 \text{ keV} \rightarrow \sigma = 15.5 \text{ mbarn}$

→ count rate: $\sim 155 / \text{h}$

→ 1 month of beam time

$E_d = 30 \text{ keV} \rightarrow \sigma = 1.2 \text{ mbarn}$

→ count rate: $\sim 12 / \text{h}$

→ 10 month of beam time

The Experimental Setup in St. Petersburg

New Setup:

ABS from the
SAPIS project:
(after upgrade)
 $\sim 4 \cdot 10^{16}$ a/s
 $\rightarrow \sim 2 \cdot 10^{11}$ a/cm²

Detector Setup:

- 4 π covered by
- large pos. sens. Detectors
 - (~300 single PIN diodes ?)

POLIS (KVI, Groningen)
Ion beam: $I \leq 20$ μ A
 $\rightarrow 1.5 \cdot 10^{14}$ d/s
($E_{\text{beam}} \leq 32$ keV)

dd-fusion polarimeter

LSP from POLIS

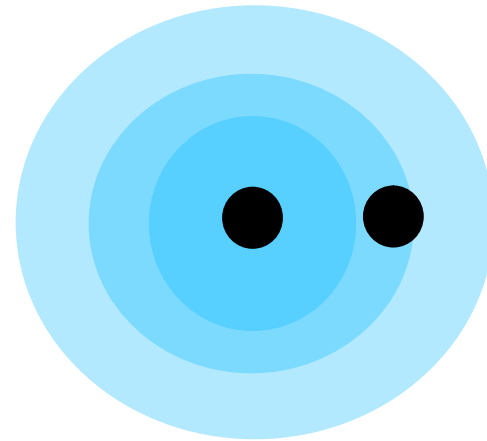
Luminosity: $3 \cdot 10^{25}$ /cm² s

\rightarrow count rate: ~ 40 /h

\rightarrow 2 month of beam time

LSP from the SAPIS project

The Electron Screening Effect



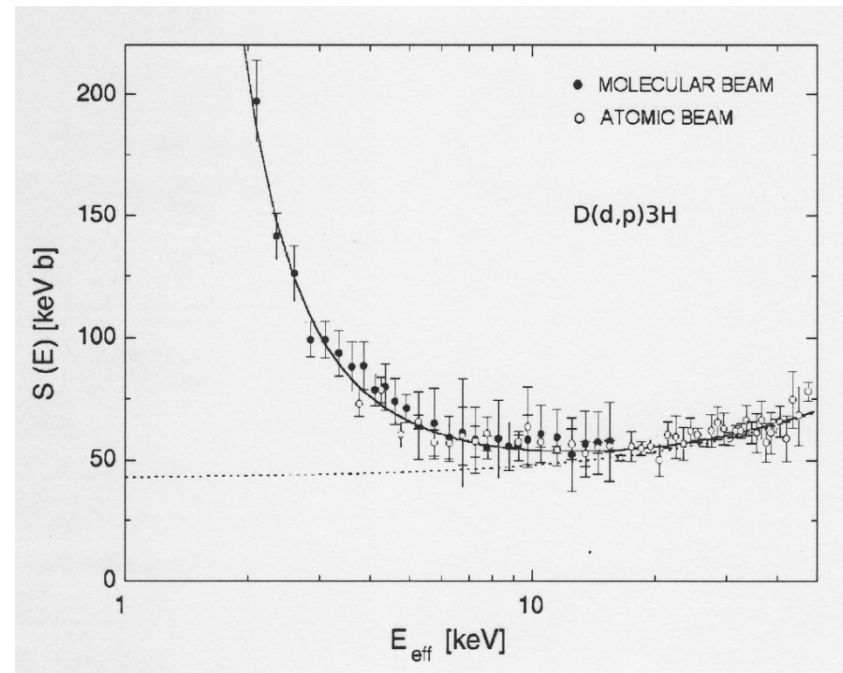
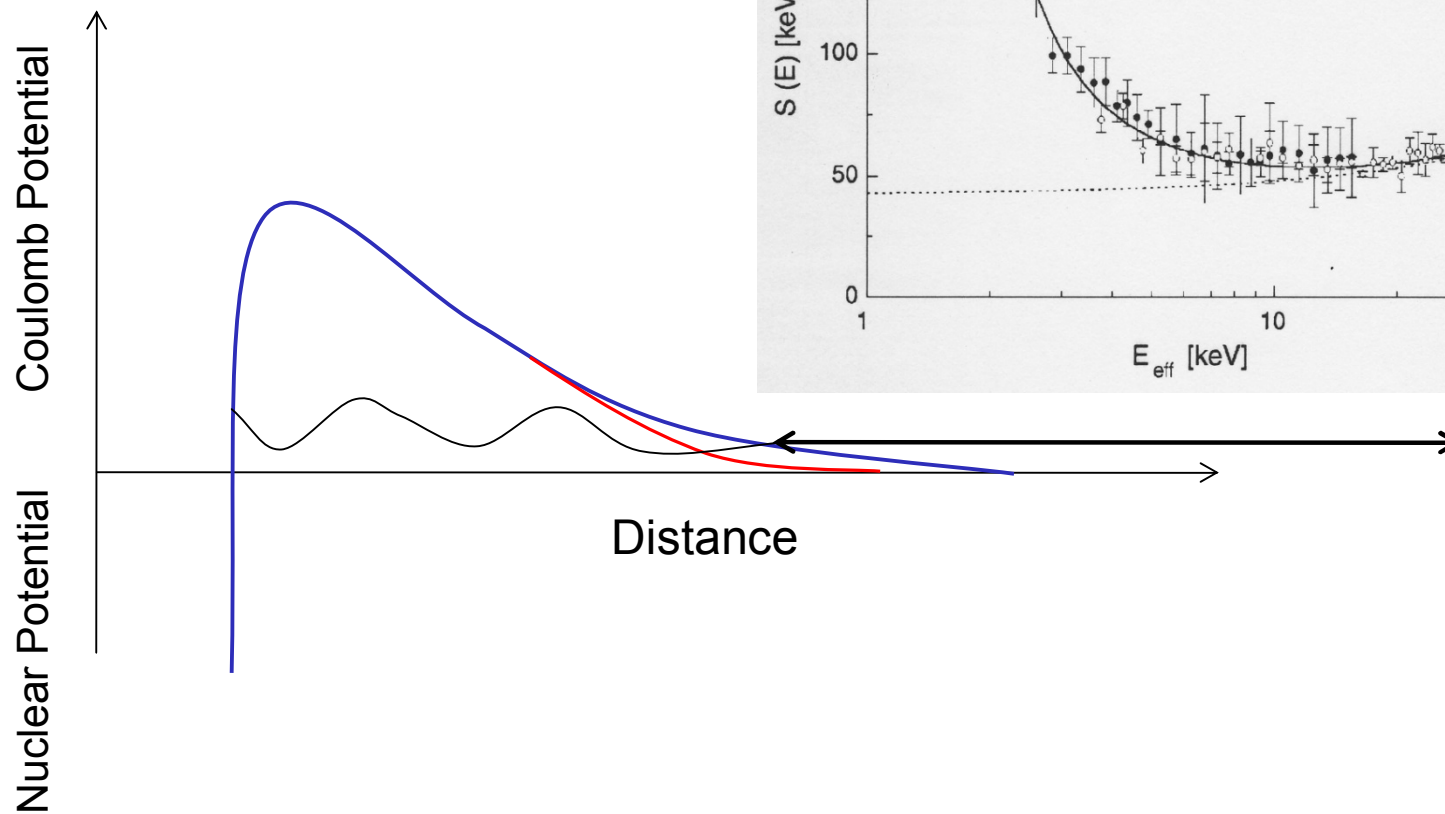
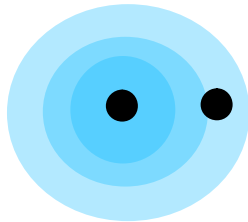
Projectile: Deuteron

Target: Deuterium Atom
(Deuteron + Electron)

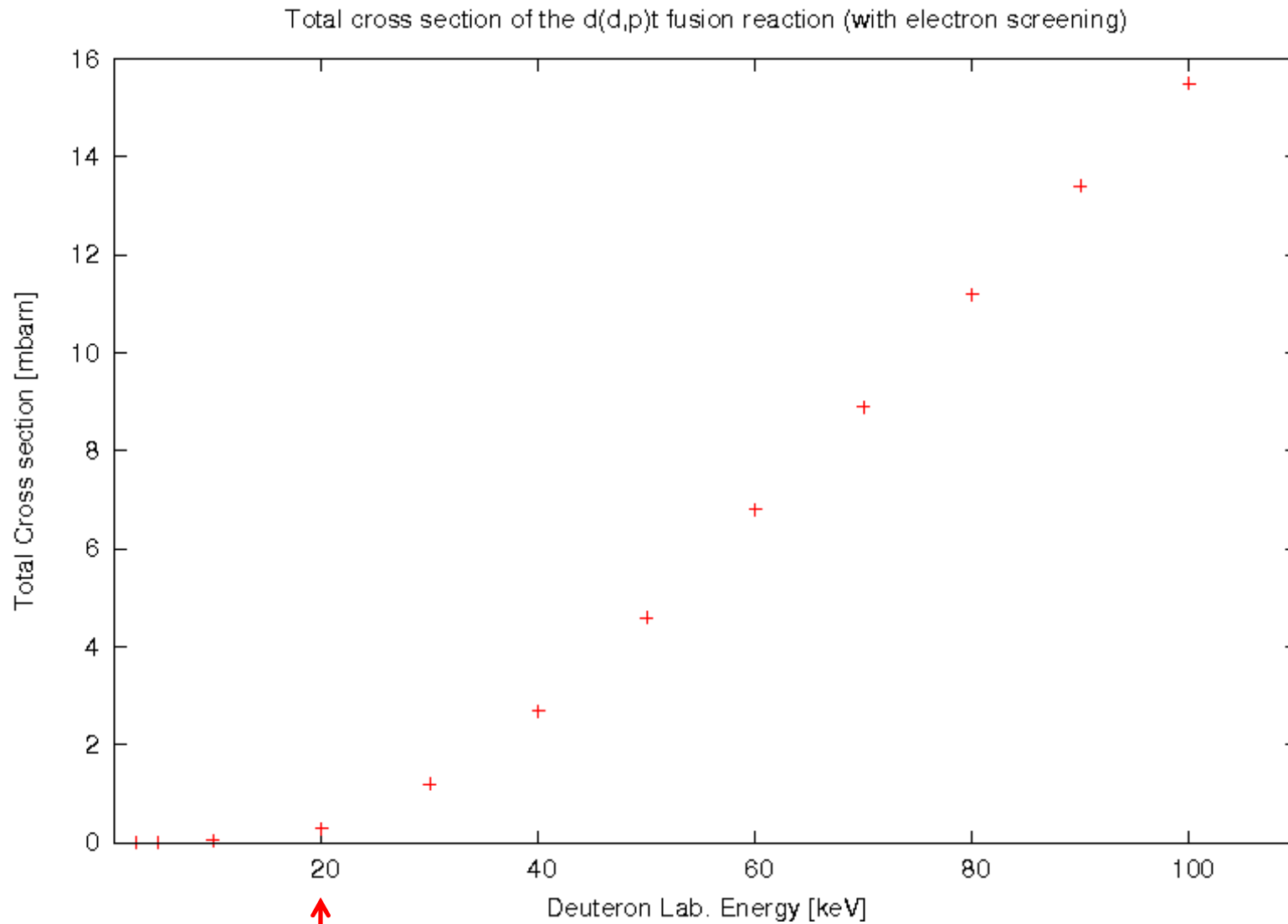
The Electron Screening Effect

Astrophysical S-Factor:

F. Raiola et al.; Eur. Phys. J. **A 13**, 377 (2002)

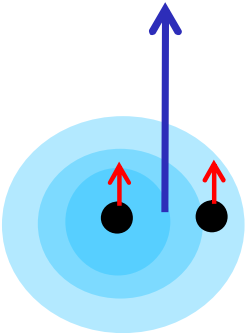


The Total Cross Section



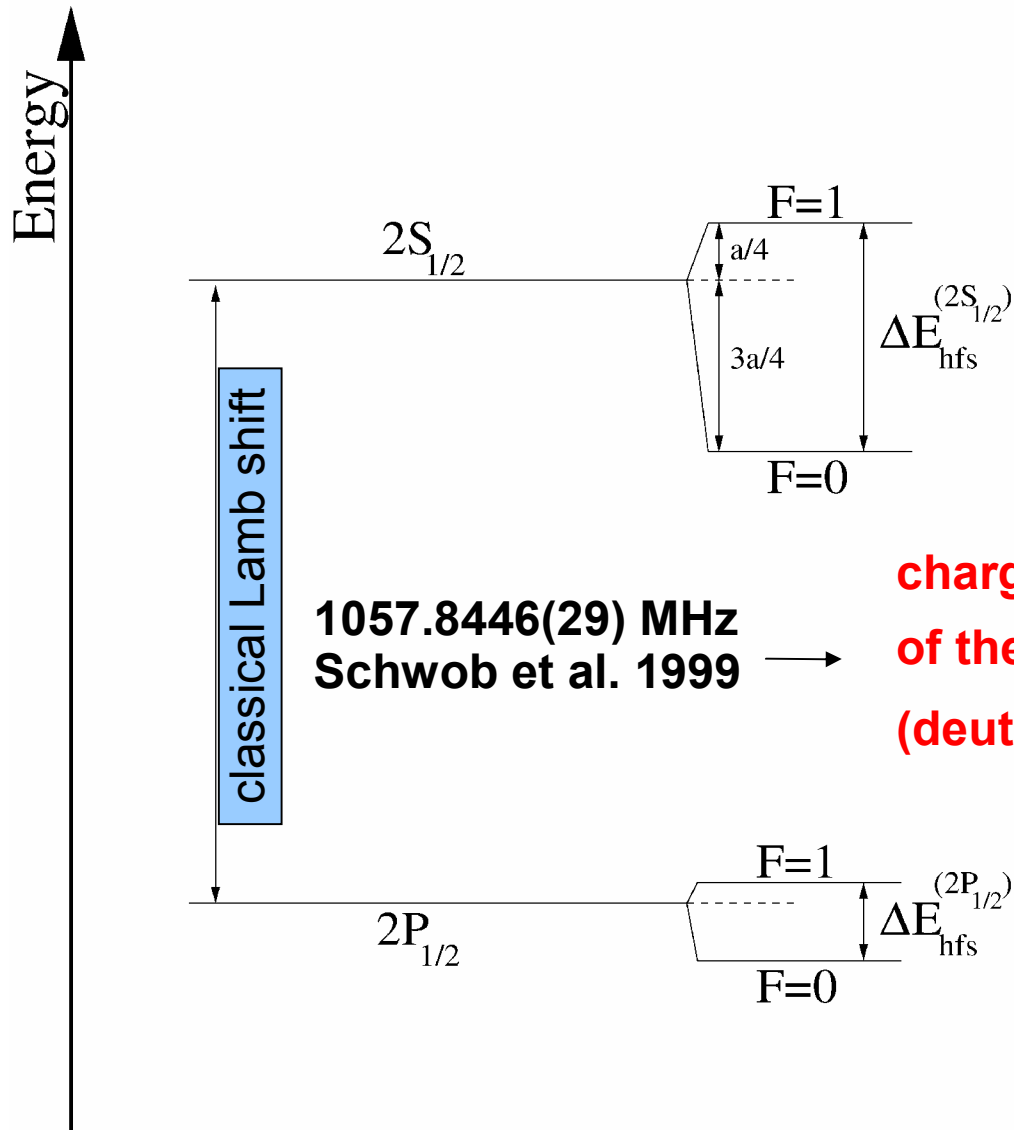
Electron Screening Effects starts

The Electron Screening Effect



Precision Spectroscopy of H and D for n=2:

Actual Status



177556834.3 (6.7) Hz
Kolachevsky et al. 2009

177556785 (29) Hz
Rothery and Hessels 2000

1057.8446(29) MHz
Schwob et al. 1999

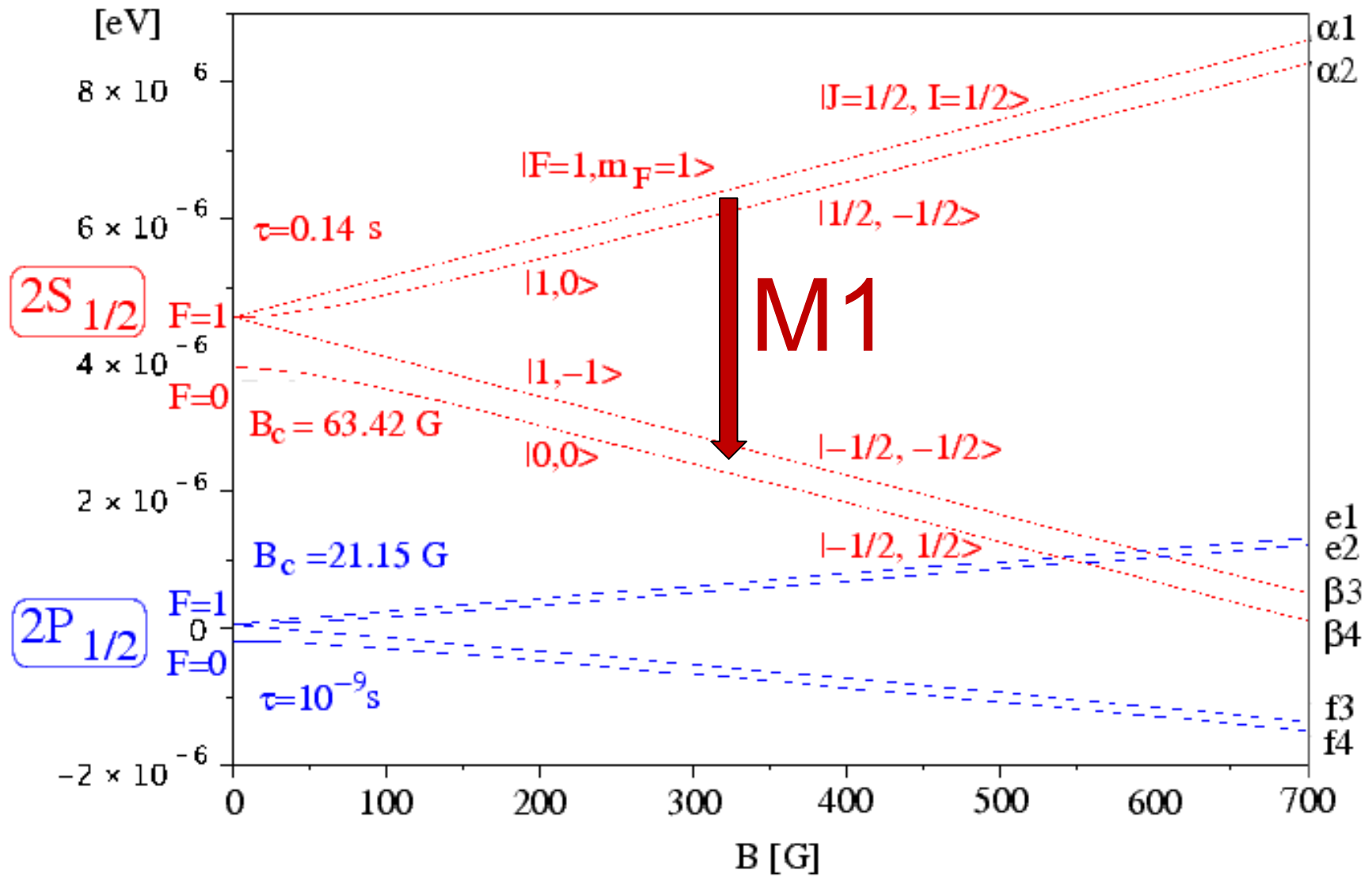
**charge distribution
of the proton
(deuteron)**

1057.842(4) MHz
Pachucki 2001

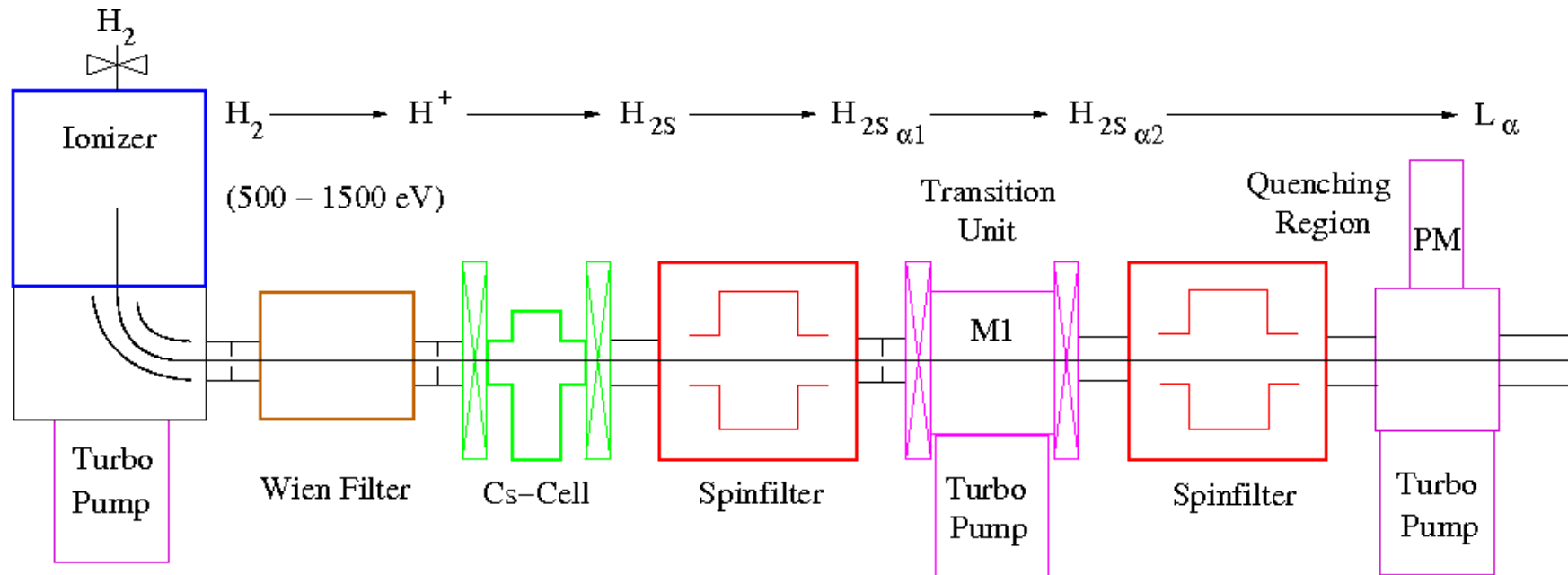
59.22(14) MHz
Lundeen et al. 1975

59.2212(?) MHz
Moskovkin et al. 2007

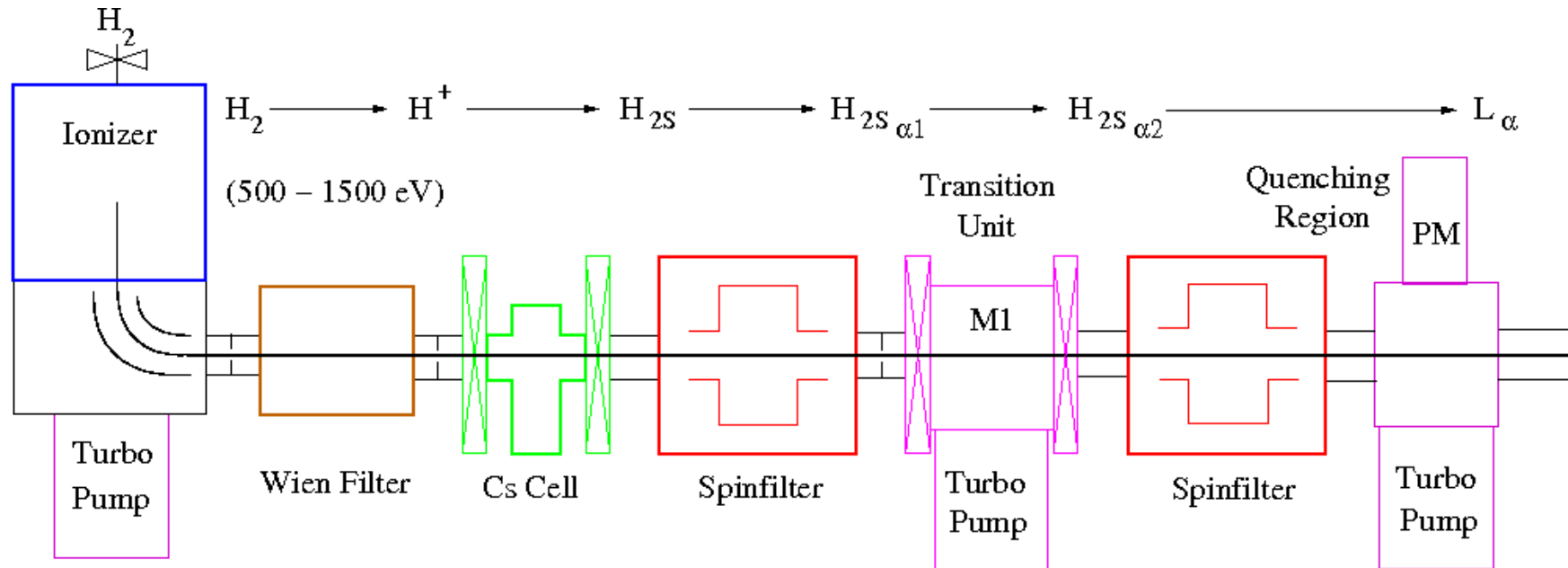
Breit-Rabi Diagram



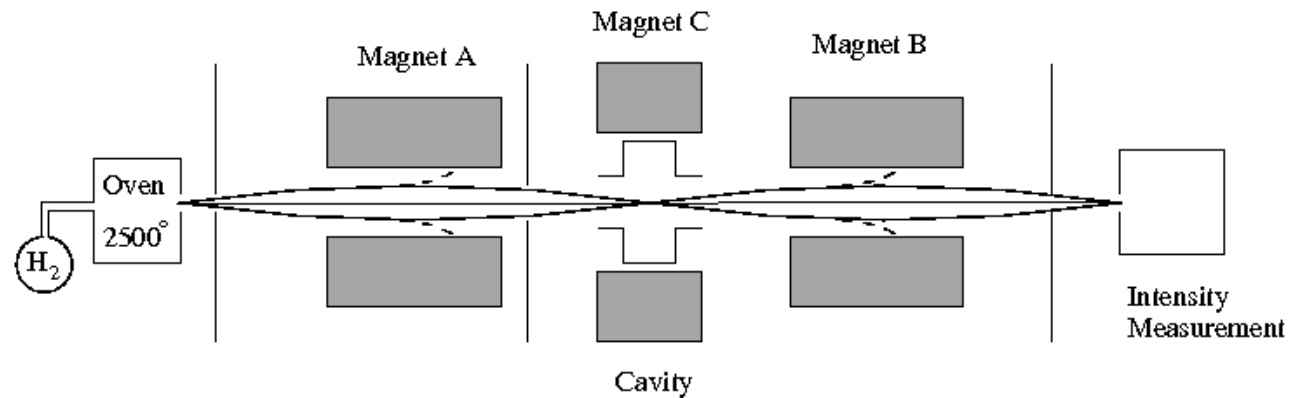
How to measure the HFS of the $2S_{1/2}$ state



How to measure the HFS of the $2S_{1/2}$ state

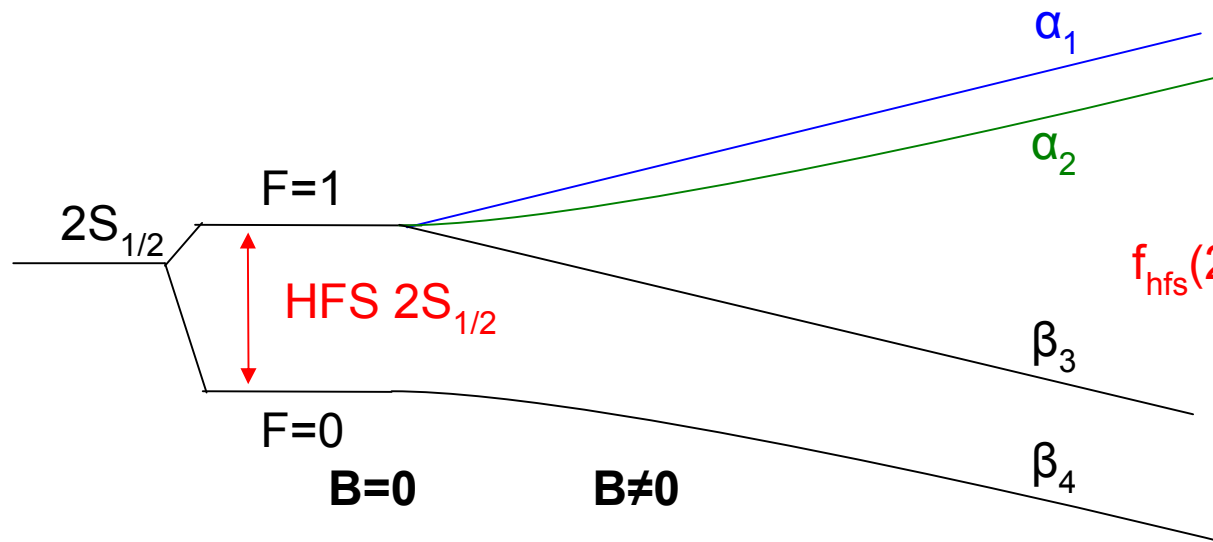


Rabi Apparatus



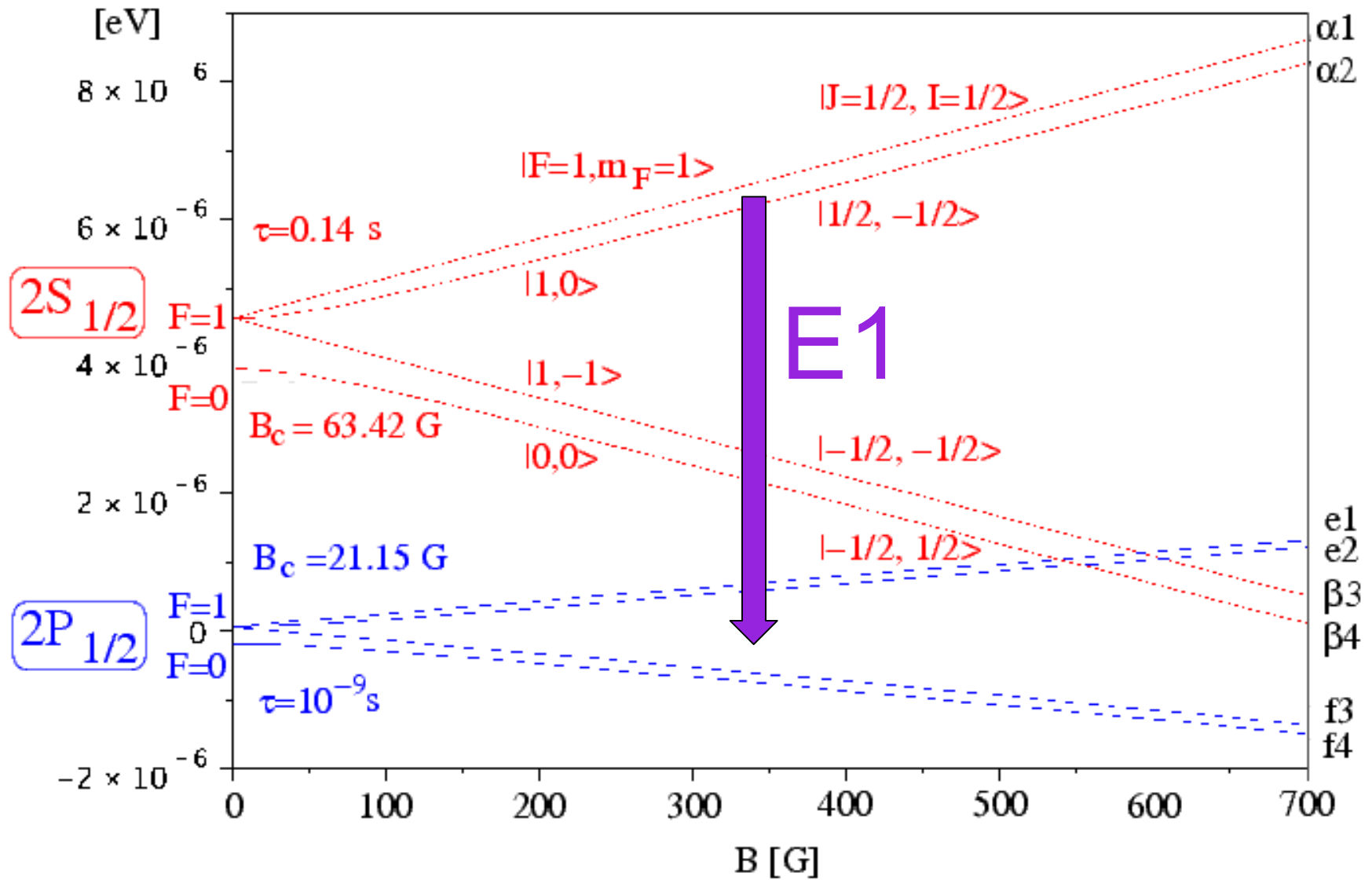
Advantages of this method

- natural linewidth: 1.1 Hz
- High Intensity: 1.5 μA protons $\rightarrow 10^{11} \text{ H}_{(\alpha_1)} / \text{s} \rightarrow 10^5 \text{ photons/s}$
- direct measurement
- Doppler free (2nd order can be measured)
- monochromatic beam energy can be changed: rel. /quadratic Doppler measurable
- **HFS $2S_{1/2}$ can be measured without knowing the absolute strength of the magnetic field B in the interaction region**
- relative magnetic field strength can be measured with the $\alpha_1 \leftrightarrow \alpha_2$ transition

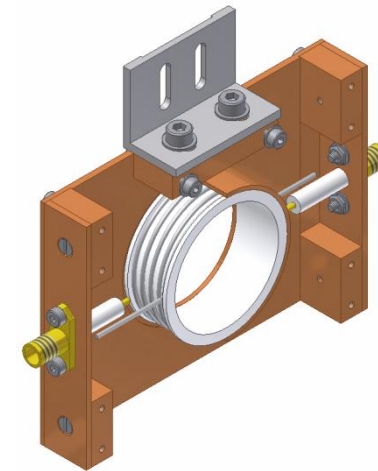
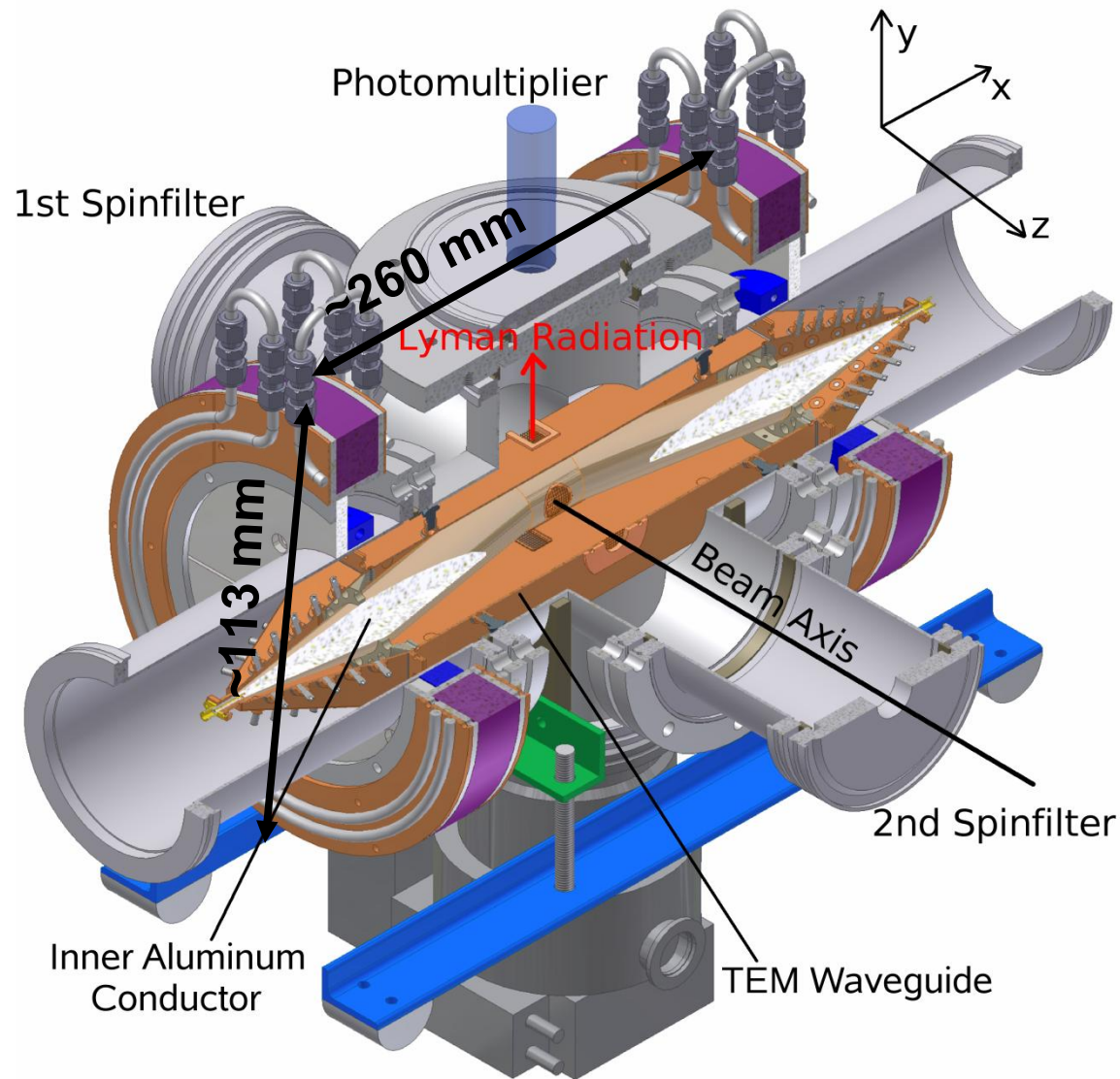


$$\begin{aligned}
 f_{\text{hfs}}(2S_{1/2}) &= f_{(\alpha_1 - \beta_4)} - f_{(\alpha_2 - \beta_3)} \\
 &= f_{(\alpha_1 - \alpha_2)} + f_{(\beta_3 - \beta_4)}
 \end{aligned}$$

Breit-Rabi Diagram



How to measure the classical Lamb shift and the HFS of the $2P_{1/2}$ state



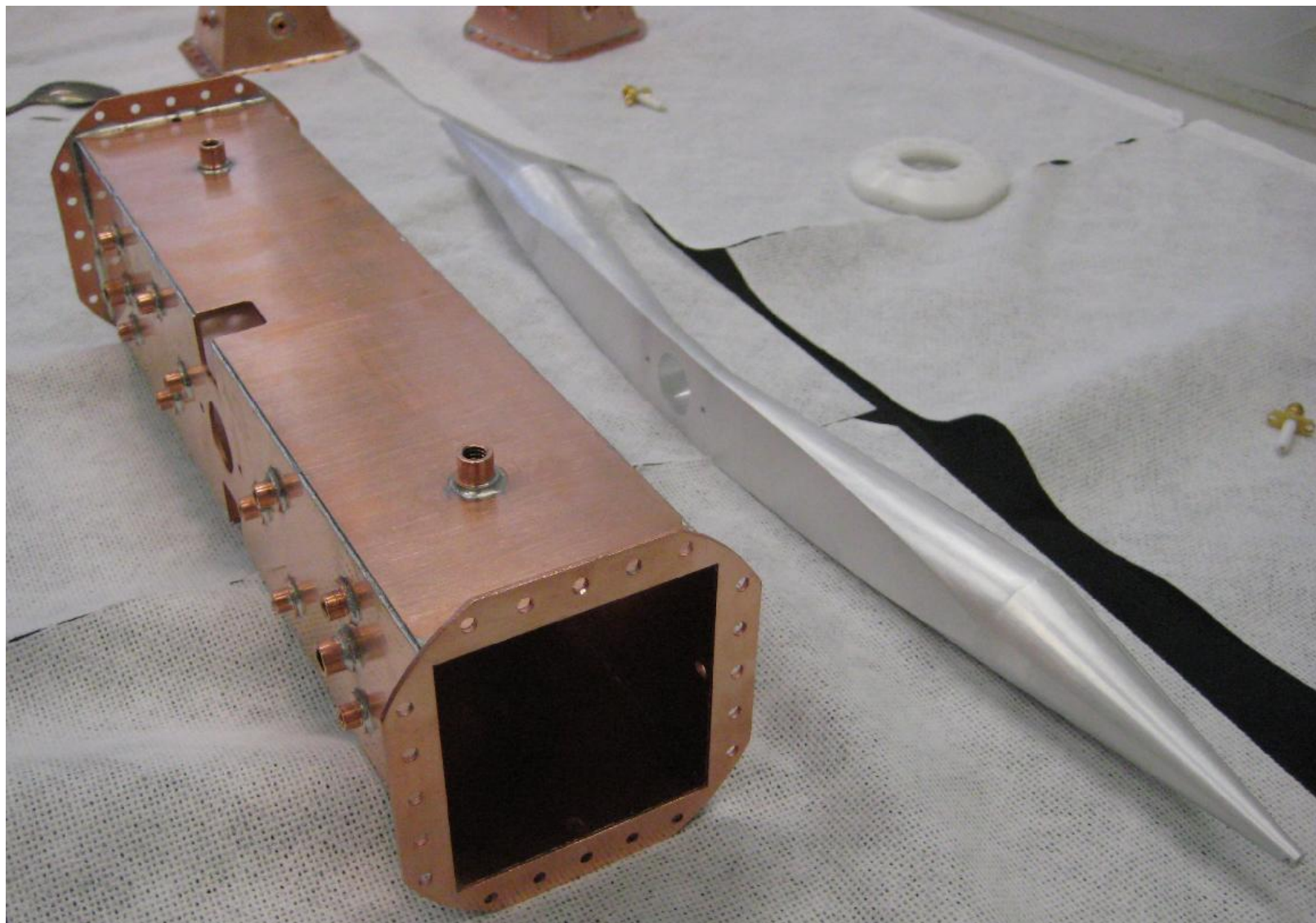
Possible Transitions:

$B \perp v$: $\alpha 1 \rightarrow f 4$ ($\alpha 1 \rightarrow e 2$)
: $\alpha 2 \rightarrow f 3$ ($\alpha 2 \rightarrow e 1$)

$B \parallel v$: $\alpha 1 \rightarrow e 1$
: $\alpha 2 \rightarrow e 2$ ($\alpha 2 \rightarrow f 4$)

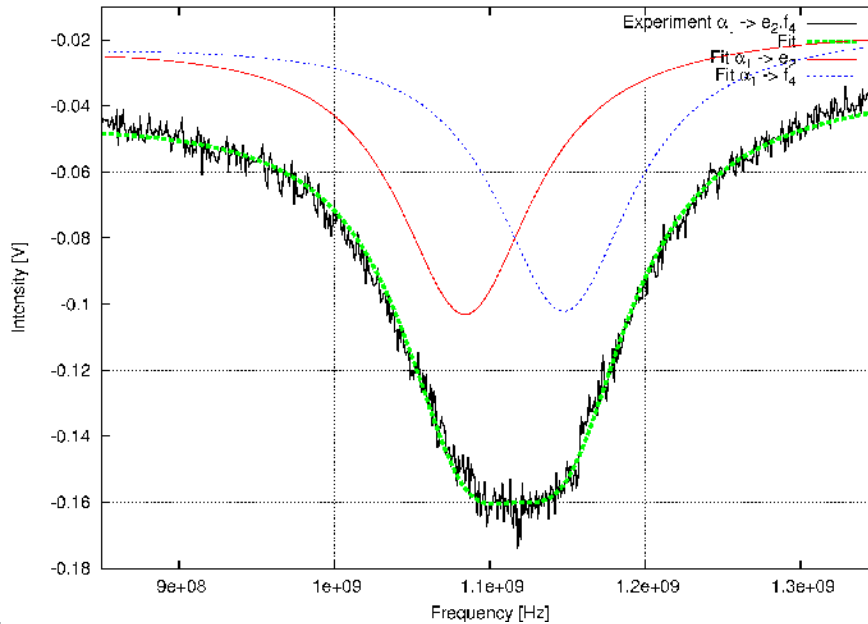
FWHM ~ 100 MHz

The TEM waveguide



Preliminary Results (B~0 G)

B ⊥ v: α1



2P_{1/2} Hyperfine Splitting

Experiment

59.98(2.03) MHz
Westig (Cologne Uni.)

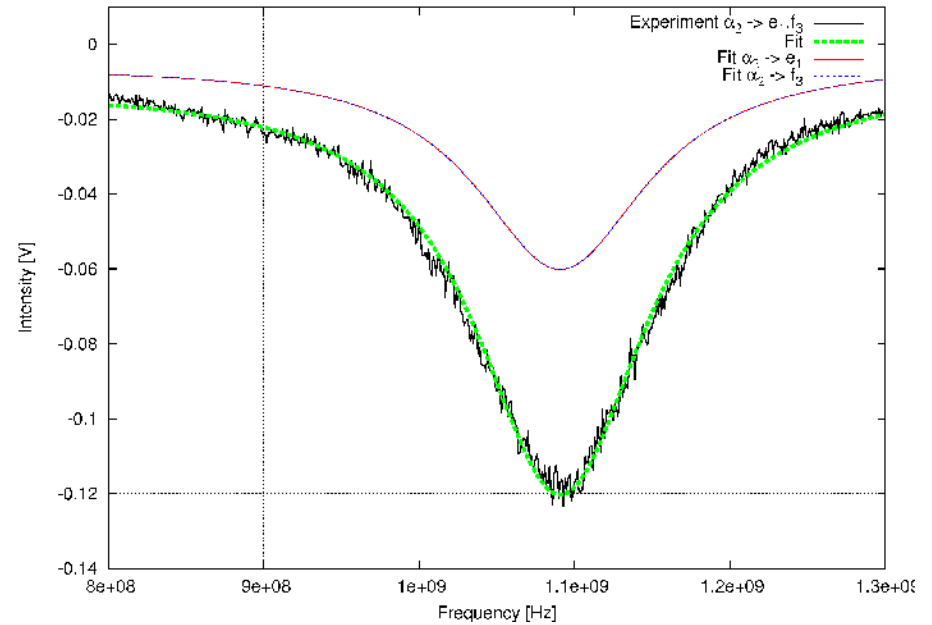
59.22(14) MHz
Lundeen et al. 1975

Theory

59.2212 MHz
Moskovkin et al. 2007

Westig et al.; Eur. Phys. J. D **57**, 27-32 (2010)

B ⊥ v: α2



classical Lamb shift

Experiment

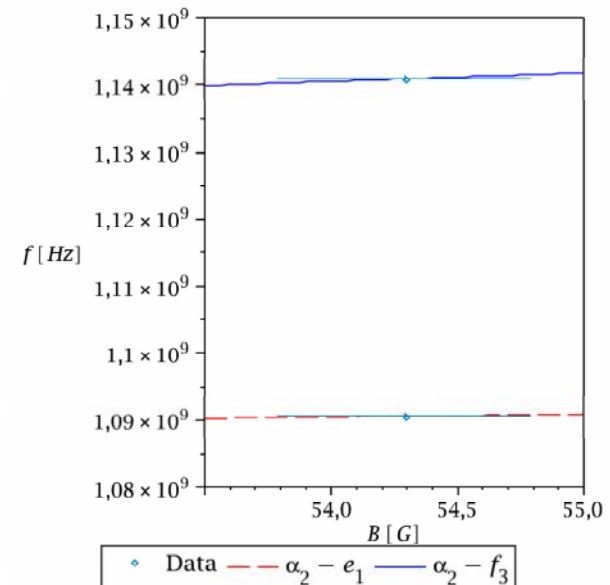
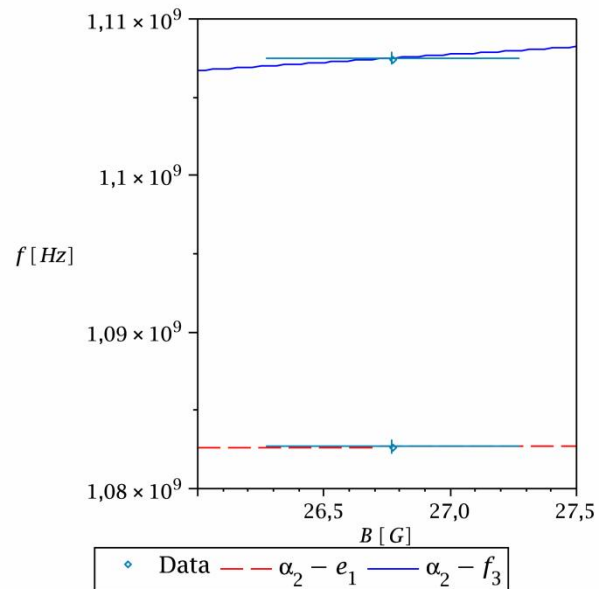
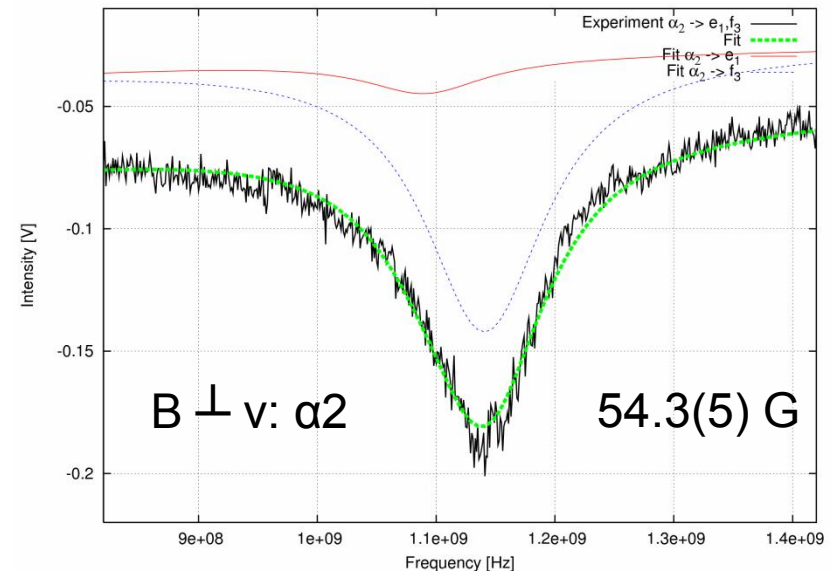
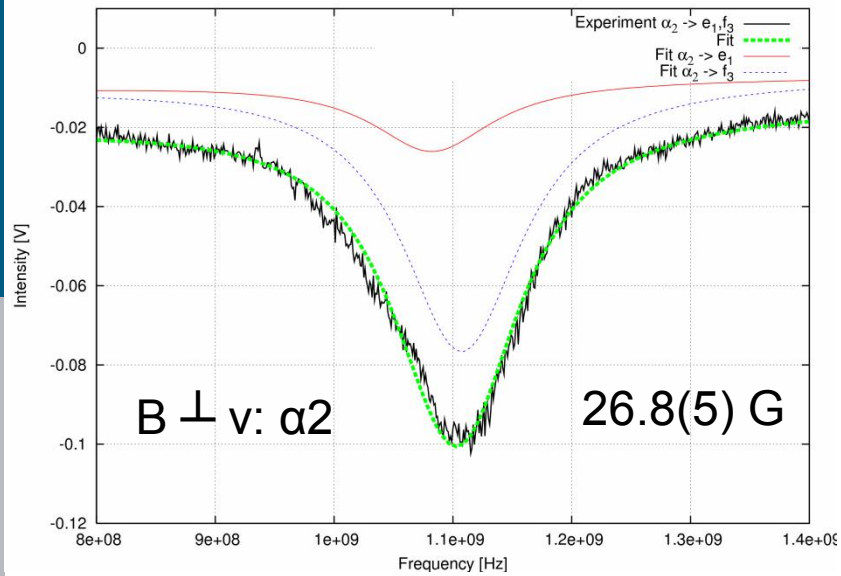
1057.34(1.11) MHz
Westig (Cologne Uni.)

1057.8446(29) MHz
Schwob et al. 1999

Theory

1057.842(4) MHz
Pachucki 2001

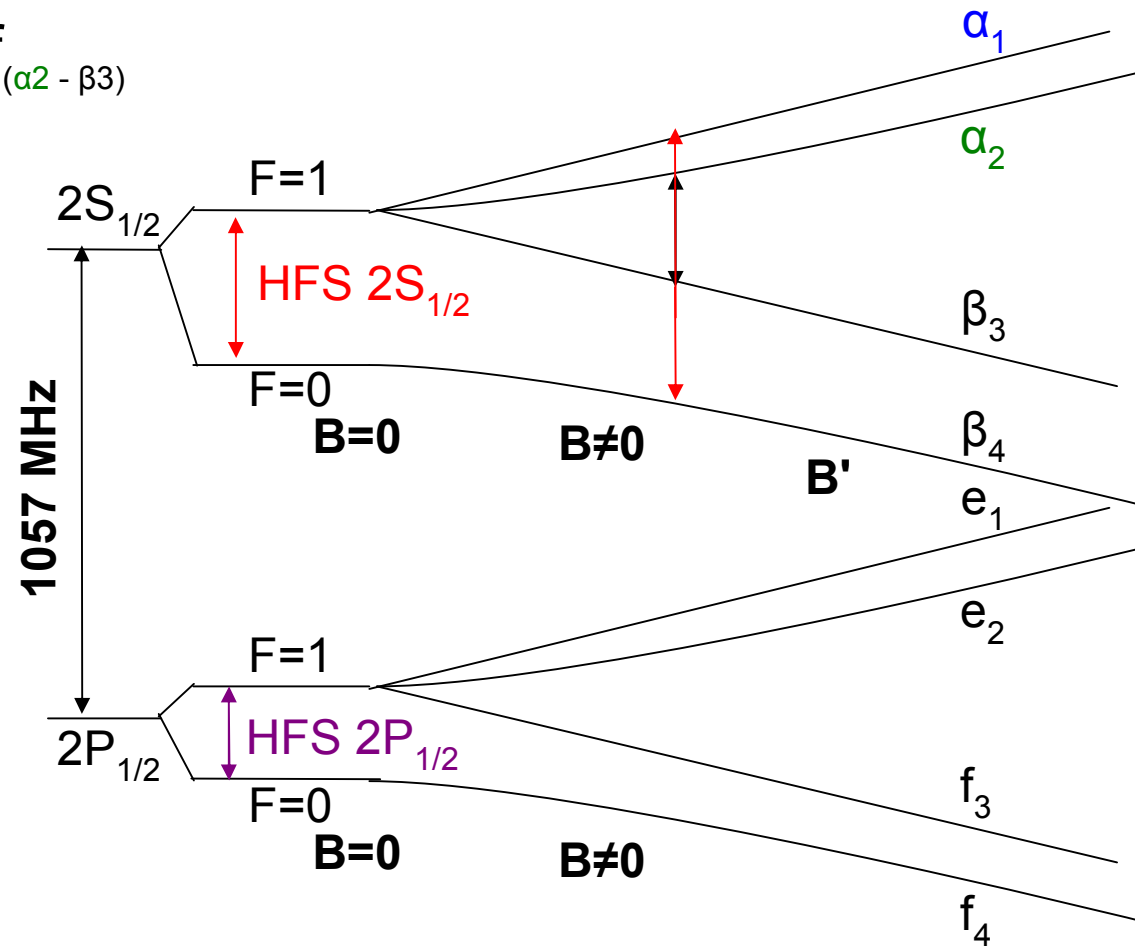
Preliminary Results (B= 26.8 and 54.3 G)



How to measure the HFS of the $2P_{1/2}$ state and the Lamb shift

$$\text{HFS } 2S_{1/2} = f_{(\alpha_1 - \beta_4)} - f_{(\alpha_2 - \beta_3)}$$

Lamb shift



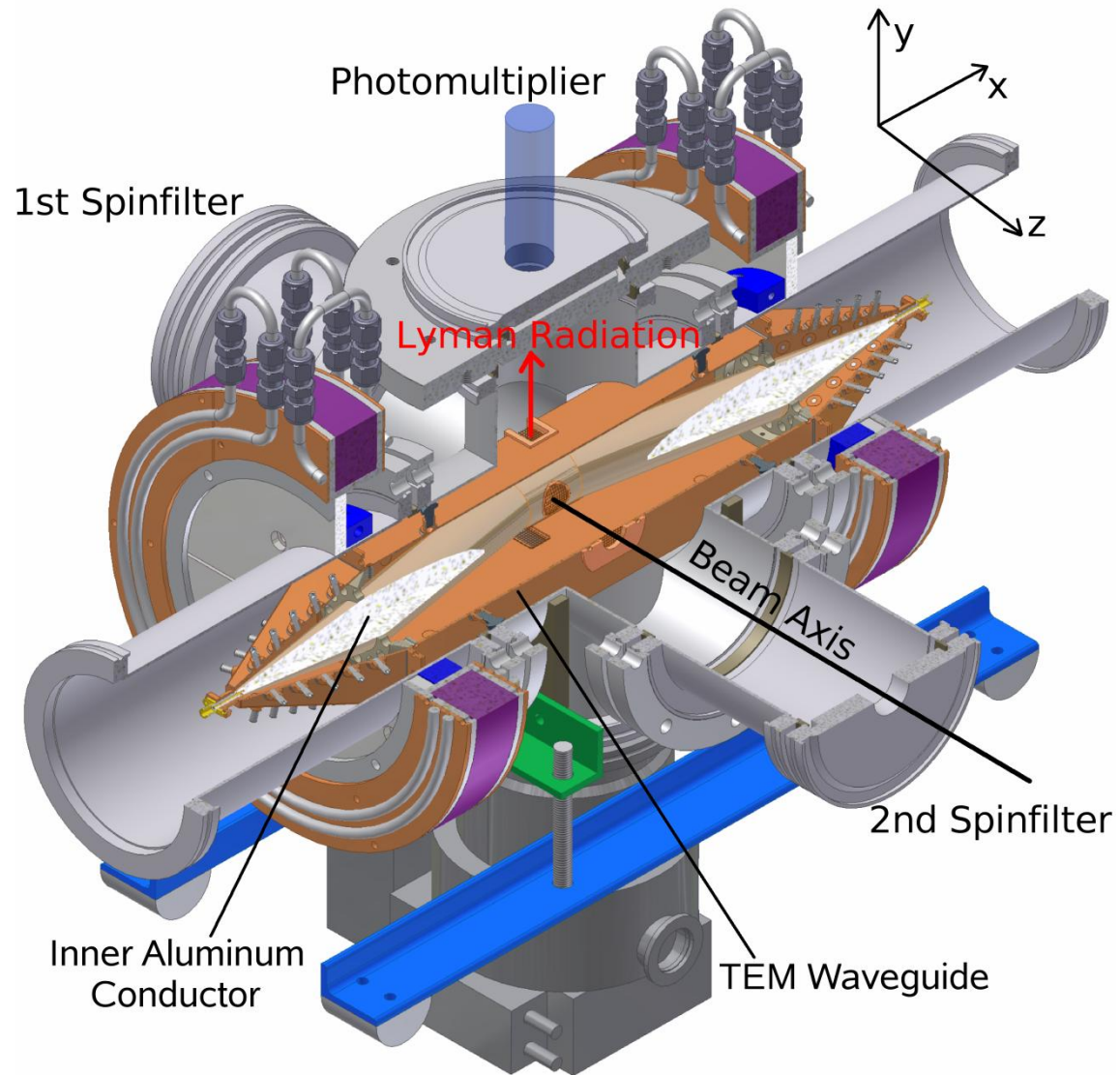
$$\text{HFS } 2P_{1/2} = [f_{(\alpha_1 - f_4)} - f_{(\alpha_2 - f_3)}] - [f_{(\alpha_1 - e_1)} - f_{(\alpha_2 - e_2)}]$$

$B \perp v$ $B \parallel v$

Uncertainty

- **Statistics:** E1 transitions: 10^7 photons/s \rightarrow in 20 min: $\Delta \sim 1$ kHz
M1 transitions: 10^5 photons/s \rightarrow in 20 min: $\Delta < 1$ Hz
- **Doppler effect:** - longitudinal suppressed ($k \perp v$)
- relativistic and transversal can be measured
- **Magnetic field:** - Homogeneity up to now: ± 0.5 G $\rightarrow \Delta \sim 1$ MHz !!!
- Magnetic field direction not well fixed ($B \sim 0$ G) !!!
- **Power measurement of the rf:** - up to now: $\Delta P \sim 3\%$ $\rightarrow \Delta \sim 100$ kHz (E1)
(New device: $\Delta P \sim 0.002\%$ are possible !)
- **Heisenberg:** $1 \sim \Delta t * \Delta f \rightarrow$ HWHH ~ 40 MHz (for 1 keV proton beam)
- Electric Fields, Motional Stark Effect, ...
- ?

Uncertainty



Uncertainty

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- ?

Summary

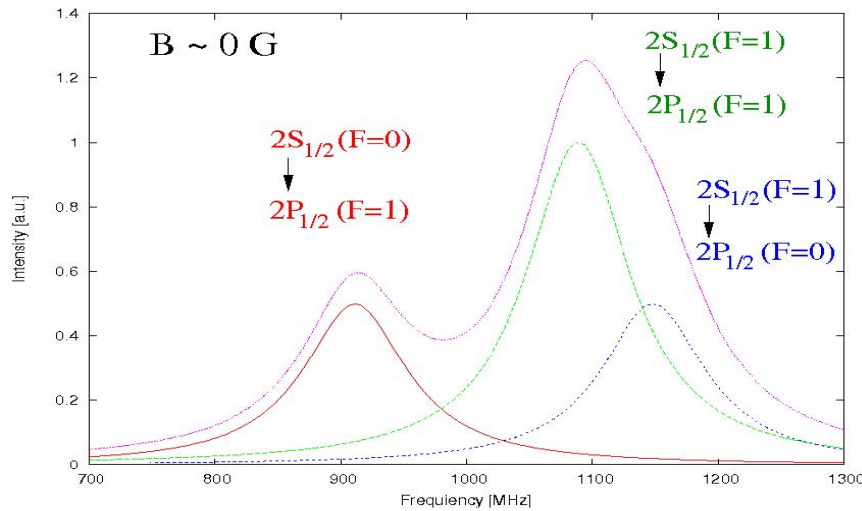
Measurement of the complete Breit-Rabi Diagram for n=2 for H and D

	1. Step (magnetic)	2. Step (SOF)	3. Step (ABS)
2S _{1/2} state:	g-factor: 10 ⁻⁵ HFS: 1 kHz	g-factor: 10 ⁻⁶ HFS: 100 Hz	g-factor: 10 ⁻⁸ HFS: 10 – 1 Hz
2P _{1/2} state:	g-factor: 10 ⁻⁴ HFS: 10 kHz	g-factor: 10 ⁻⁵ HFS: 1 kHz	g-factor: 10 ⁻⁶ HFS: 100 Hz
Lambshift:	10 kHz	?	100 Hz
2P _{3/2} state	g-factor: 10 ⁻⁴ HFS: 10 kHz	g-factor: 10 ⁻⁵ HFS: 10-1 kHz	

An alternative method to measure g-factors, the Lamb shift and the hyperfine splittings for Anti-Hydrogen at FAIR ?

Ramsey: SOF-Method (1950)

Lundeen, Jessop and Pipkin, Phys. Rev. Lett. **34** (1975)



HFS $2P_{1/2} = 59.22 (14) \text{ MHz}$

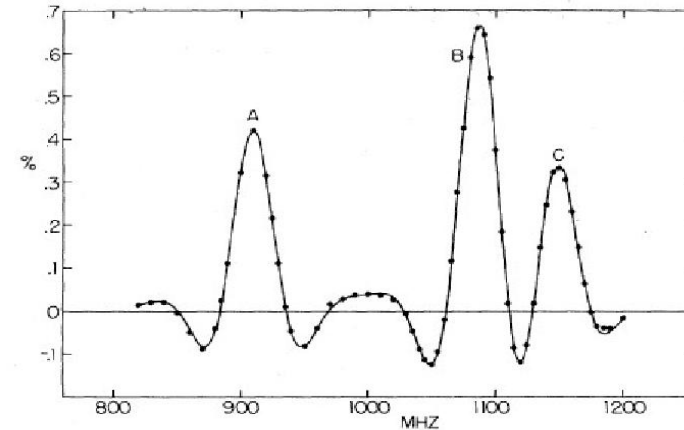


FIG. 2. The line profile observed for the $2^2S_{1/2} \rightarrow 2^2P_{1/2}$ transition with separated oscillatory fields.

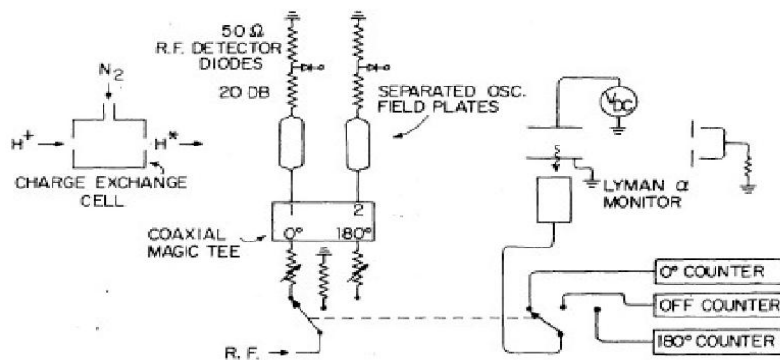


FIG. 3. A schematic diagram of the essential parts of the apparatus. The coaxial magic *T* divides the rf power equally between the two rf plates with a relative phase of 0° or 180° depending upon the input port used.

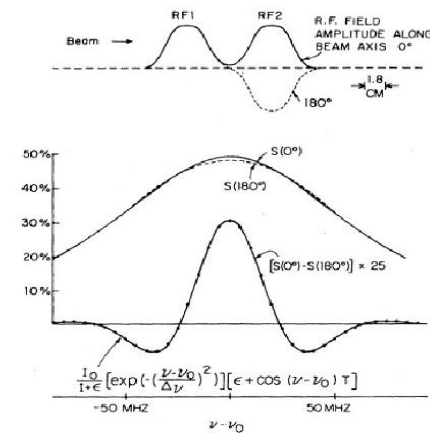


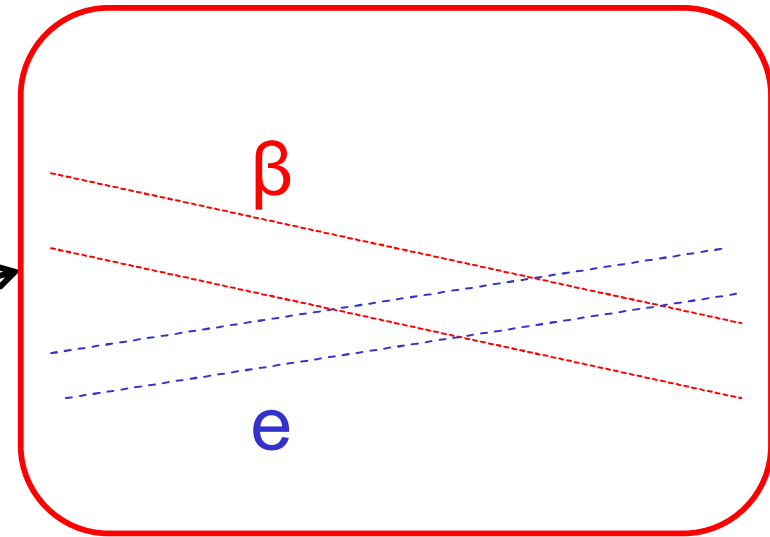
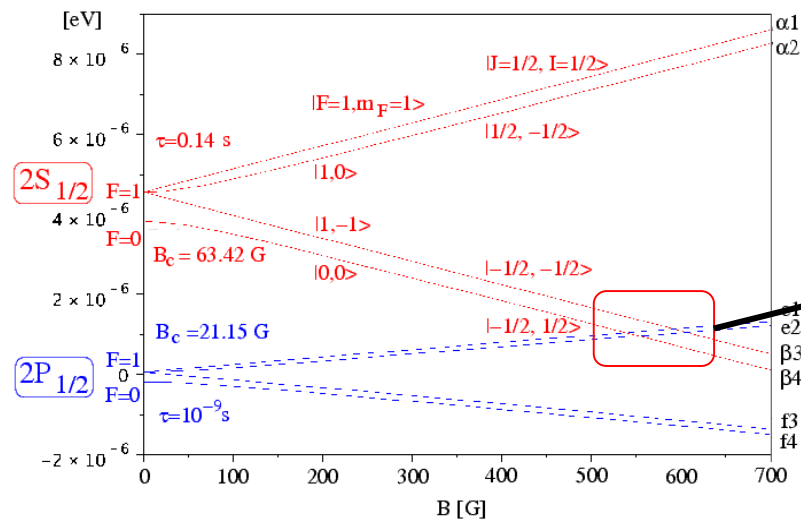
FIG. 4. A diagram showing the formation of the interference signal and the empirical fitting function.

What else ?

- Deuterium (tritium) is possible (α_1 , α_2 and α_3) [D_{21} Theory]
- Other atoms are possible: ^3He , ...
- **Method will work for Antihydrogen:**
 - During recombination of antiproton and positron up to 30% of the antihydrogen ends up as metastable atoms in the $2S_{1/2}$ state !!!
 - PM registers single photons
 - Huge range of beam energies is possible: $1/40 \leftrightarrow 2000$ eV (smaller energy \rightarrow smaller spinfilter \rightarrow higher intensity)
 - Intensity: count rate 1 Photon/s $\leftrightarrow 10^{3-4} \text{H}_{(2S)}/\text{s}$ now !

$\bar{\text{H}}$ Hyperfine Splittings of the $2S_{1/2}$, $2P_{1/2}$ ($2P_{3/2}$) and the classical Lamb shift can be measured with the spinfilter.

Outlook: Parity Violation



Direct transitions between β and e states are not allowed !!!
 (**Parity conservation**)

Weak force is part of the binding energy of the S states !!!

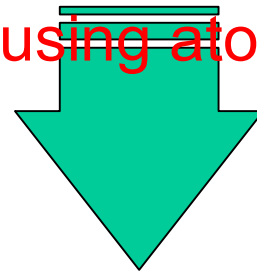
Weak force parity violation \rightarrow β - e transitions are possible

(R.W. Dunford and R.J. Holt; J. Phys. G **34** (2007) 2099-2118)

Polarized H₂ Molecules

- Beam intensities of conventional ABS barely reach $\sim 10^{17}$ at/s
⇒ target density $d_t \sim 10^{14}$ at/cm² (typical T-shaped storage cell)
- Depolarization at low T of storage cell don't allow further cooling

Performance of PIT using atomic beams saturates!



New storage cell materials

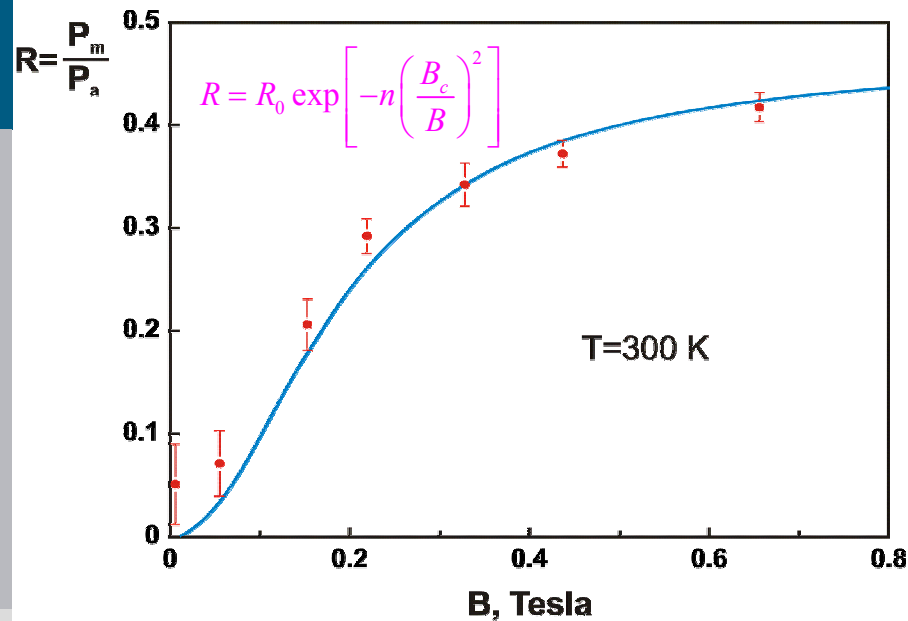
Polarized molecules?

Polarized H₂ Molecules

- **Sticking time** of molecules at the surface is much **smaller** compared to atoms
 - ⇒ cell can be cooled down to much **lower T**
 - ⇒ **higher** target density ($d_t \sim T^{-1/2}$)
- Polarized molecules is an interesting object for **atomic physics** which has never been deeply investigated (e.g. depolarization on the surface)
- Recombination of polarized atoms in different hyperfine states is interesting to **astrophysics** (e.g. formation of molecular hydrogen in cold clouds)

Polarized H₂ Molecules

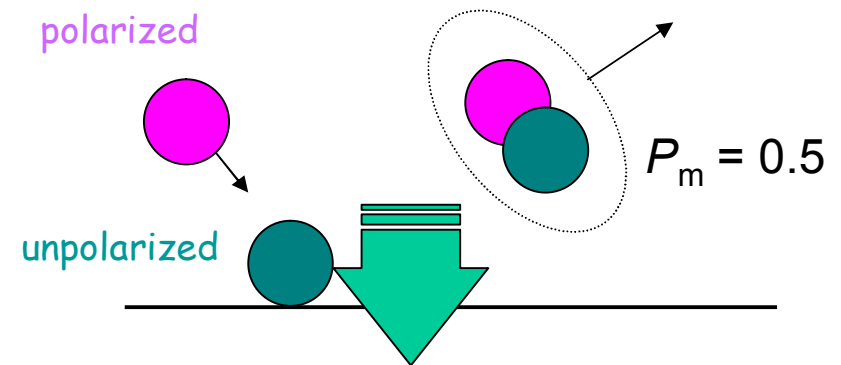
Measurements from NIKHEF, IUCF, HERMES show that recombined molecules retain fraction of initial nuclear polarization of atoms!



Nuclear Polarization of Hydrogen Molecules from Recombination of Polarized Atoms
T.Wise et al., Phys. Rev. Lett. 87, 042701 (2001).

$$\lim_{B \rightarrow \infty} R = 0.5$$

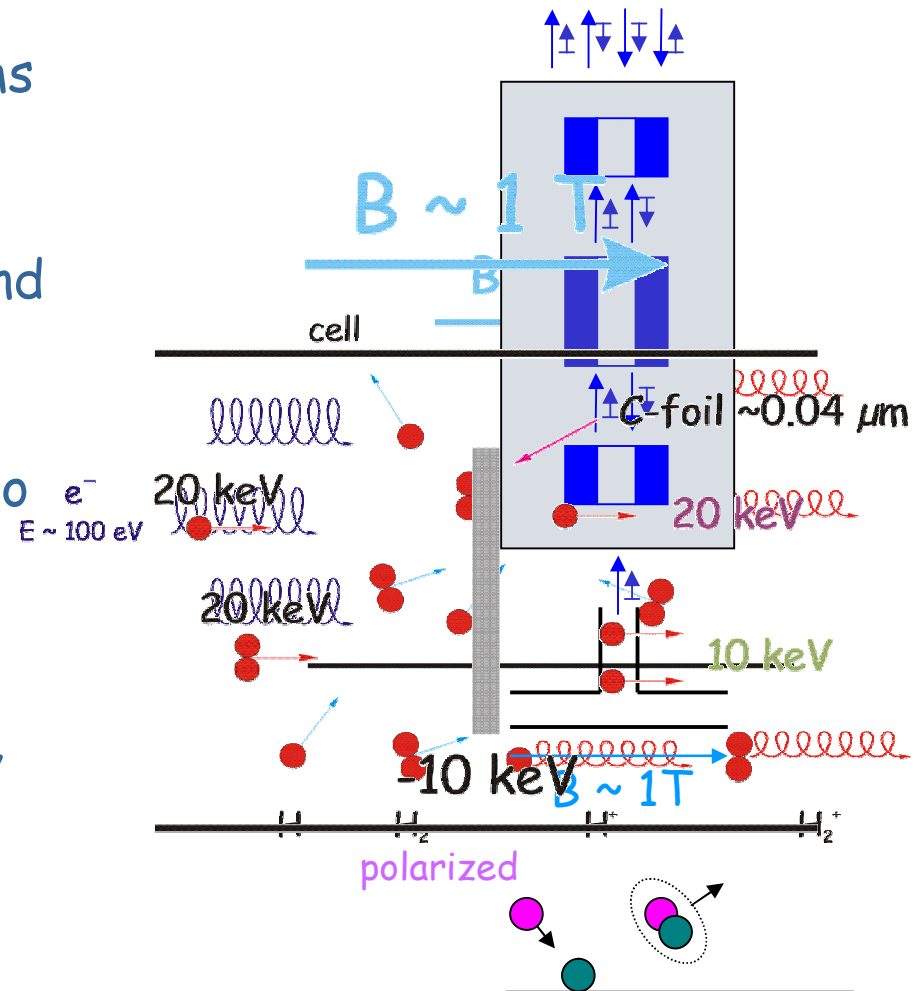
Naive model



Is there a way to increase (surface material, T, B etc)?

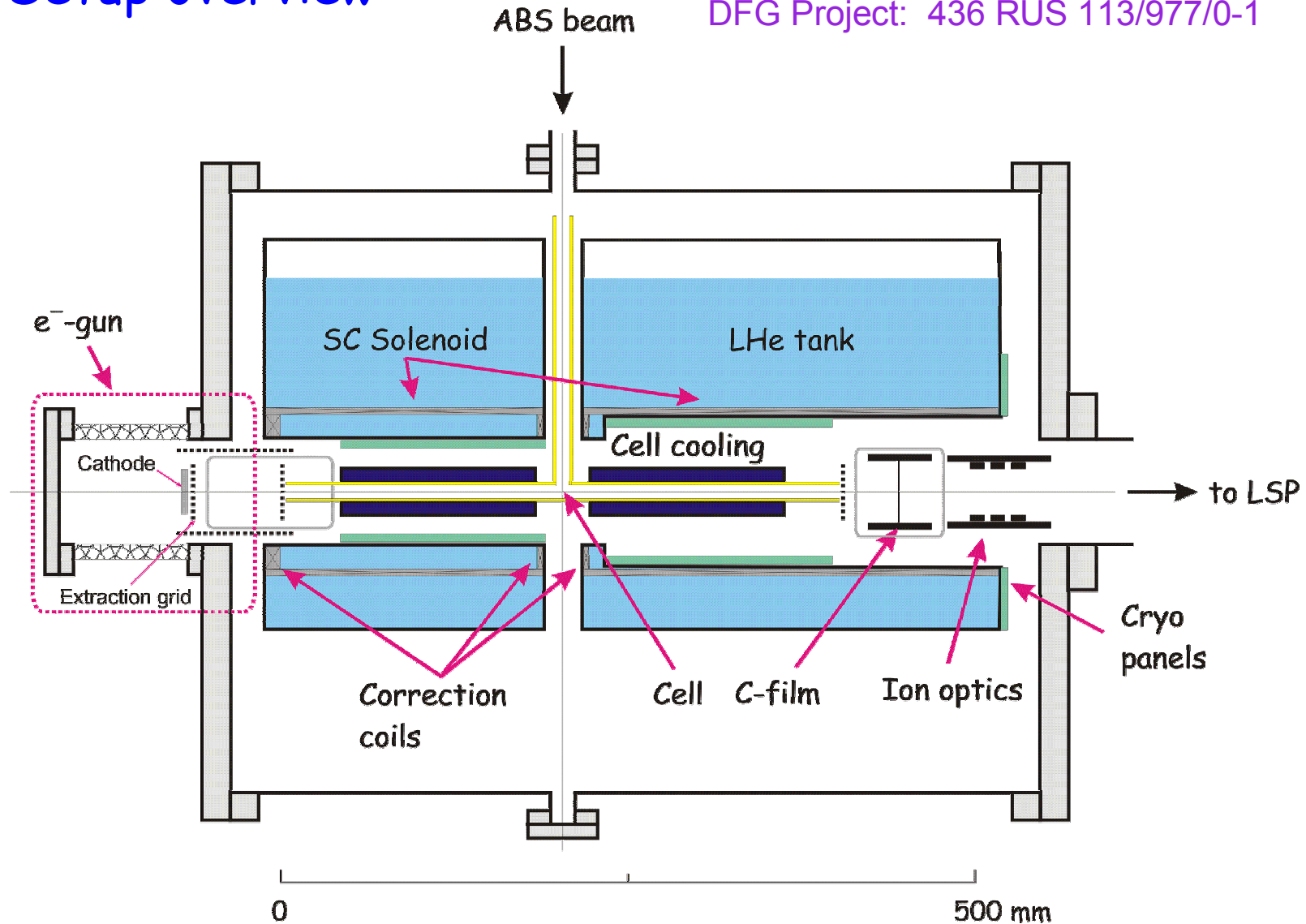
Polarized H₂ Molecules

- Recombination of polarized atoms into molecules
- Conversion of polarized atoms and molecules into ions
- Conversion of H₂⁺ and H⁺ ions into protons with different energy (suggested by W.Haerberli)
- Separation of protons by energy
- Measurement of proton polarization in LSP

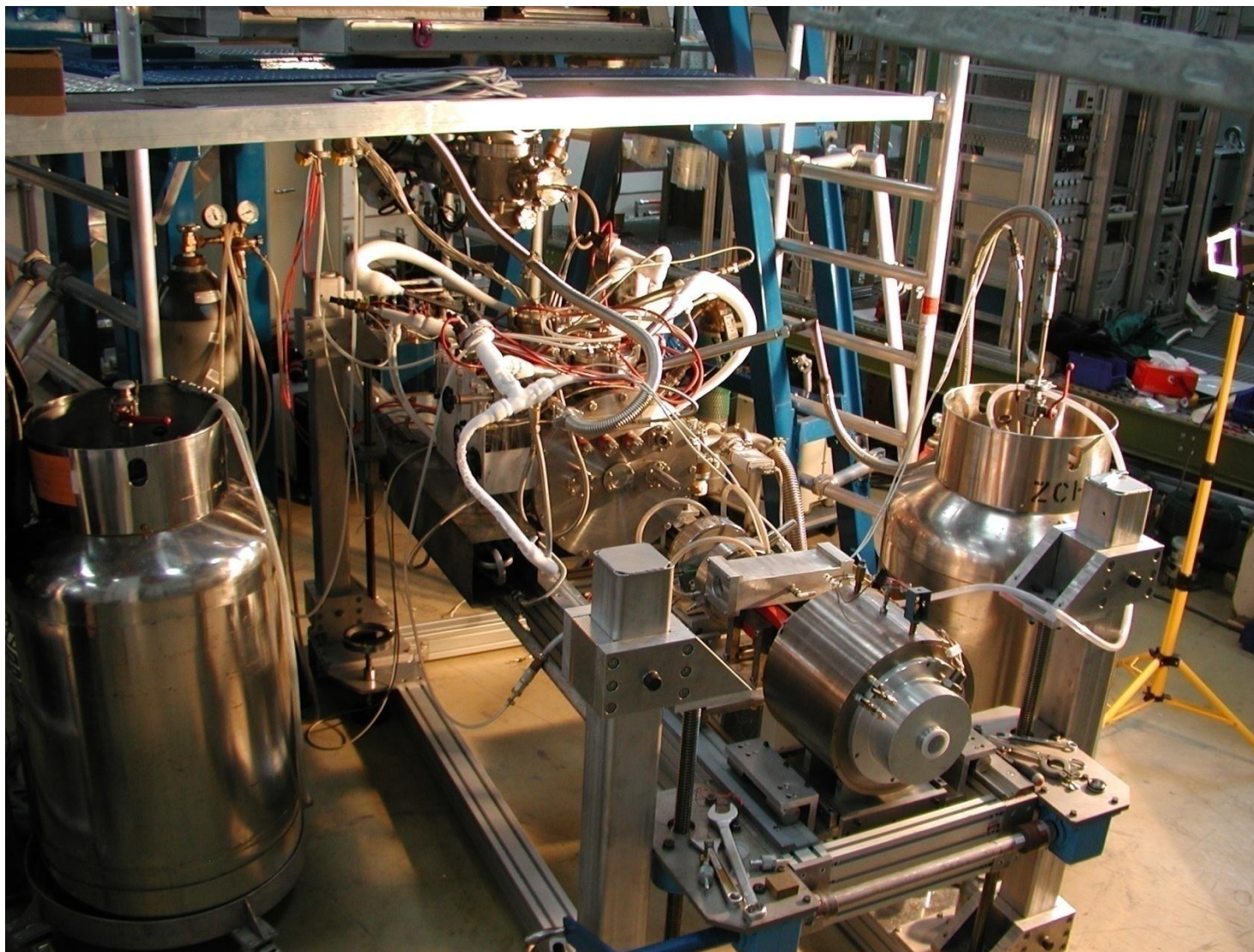


Polarized H₂ Molecules

Setup overview

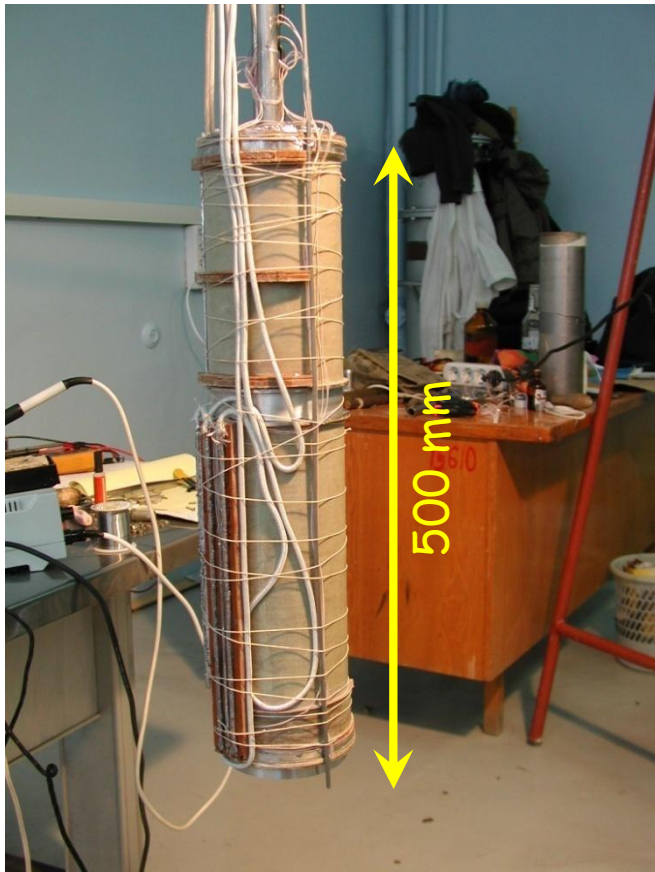


Polarized H₂ Molecules

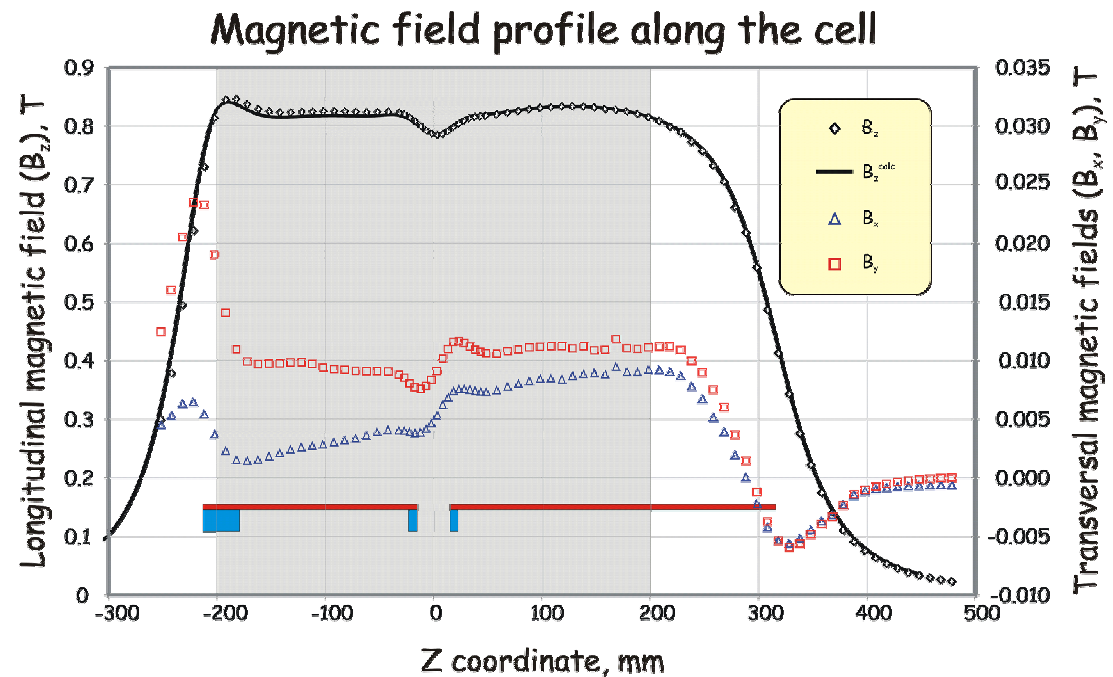


Polarized H₂ Molecules

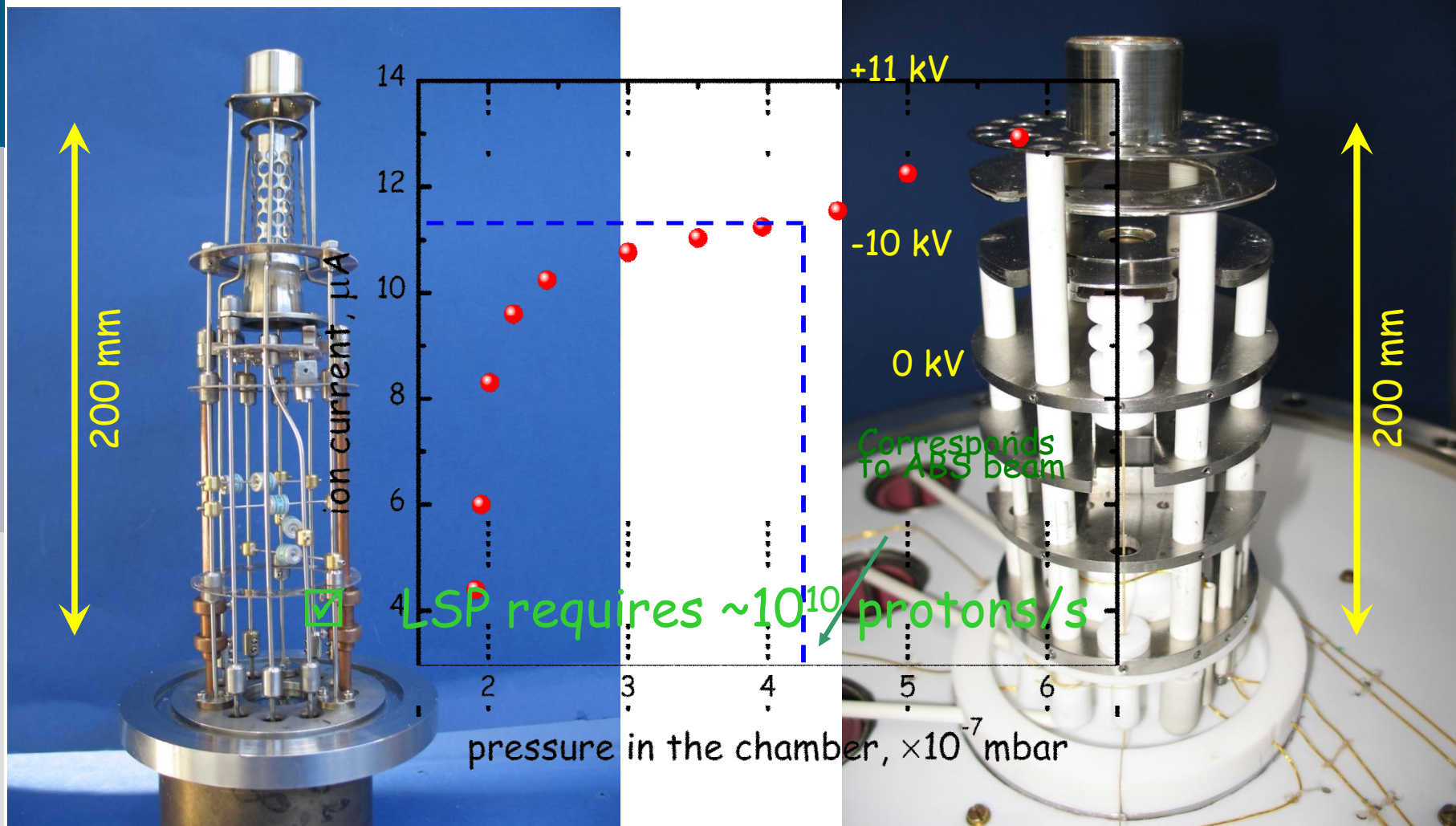
Superconducting Solenoid



- SC wire NiTi+Cu (Ø 0.5 mm)
- Nominal current 50 A \Rightarrow B ~ 1 T
- Degradation of frozen field \leq 0.1% per 5 hrs
- LHe consumption ~ 8 l/h



Polarized H₂ Molecules e⁻-gun and ion optics



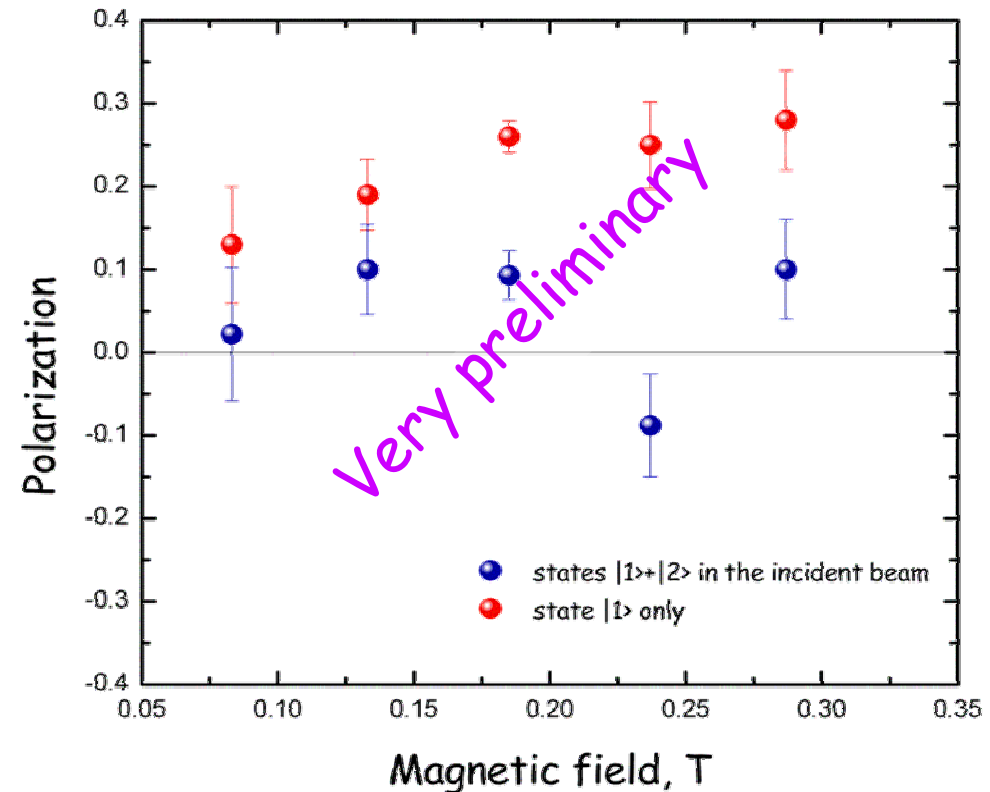
Polarized H₂ Molecules

Magnetic field dependence

- Cell coating: Au
- Cell temperature: 80K

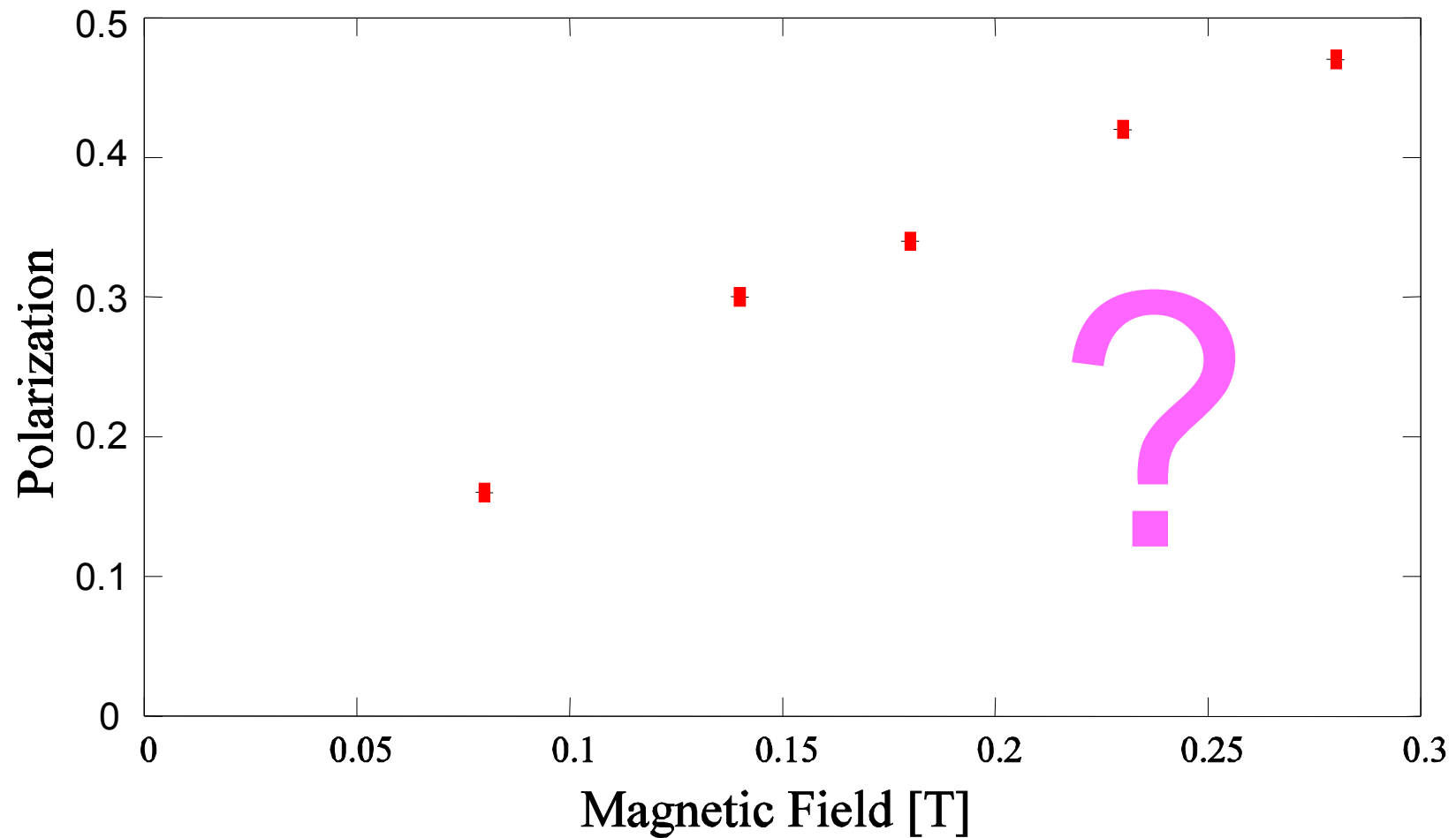
Low polarization???

- misalignment of quantization axis?
- T too low?
- Au is not an appropriate material?



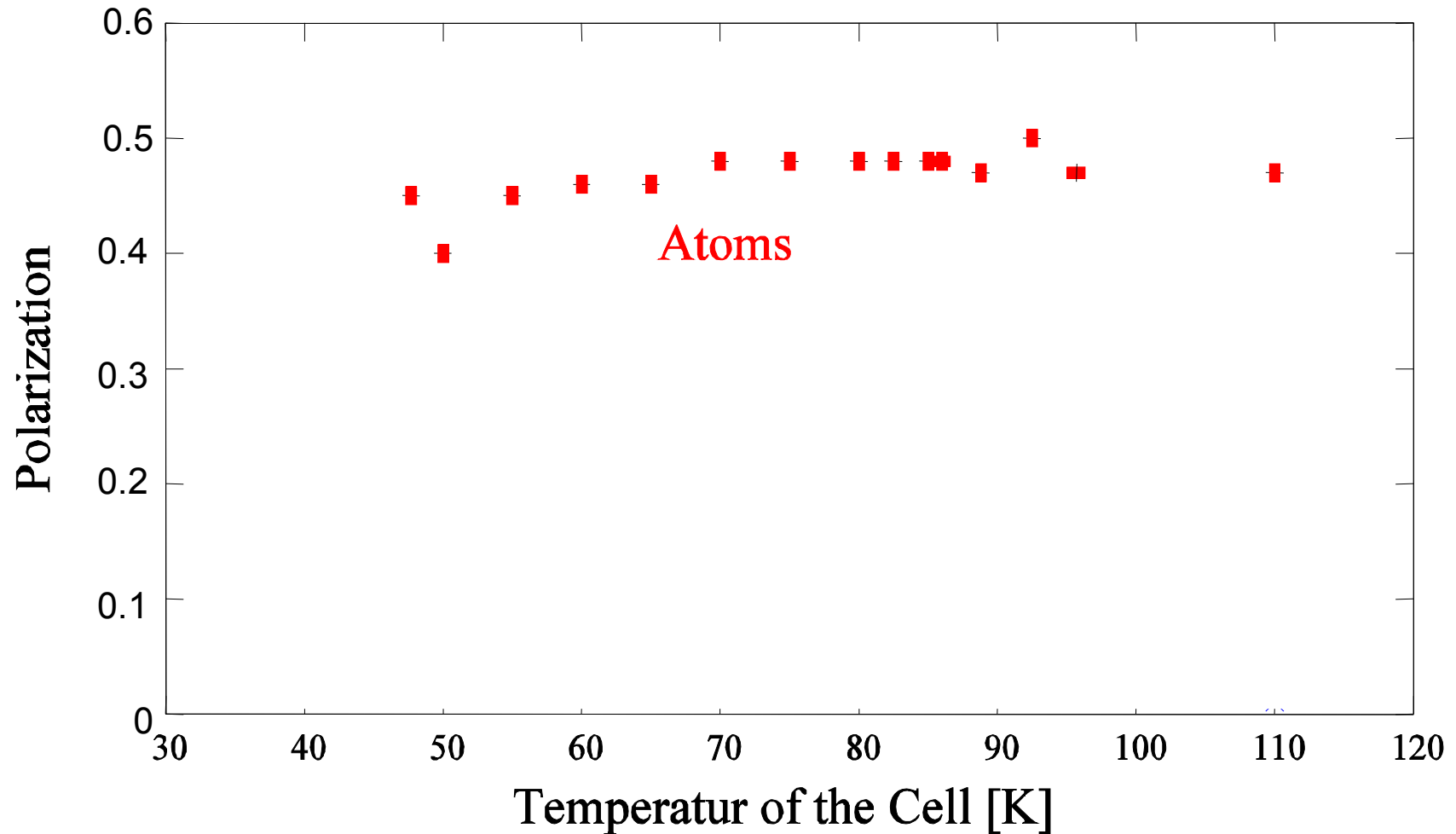
Polarized H₂ Molecules

Polarization of Hydrogen Atoms as Function of the Magnetic Field
(Surface: Gold, T = 47 K, HFS 1, Q = 3 keV)



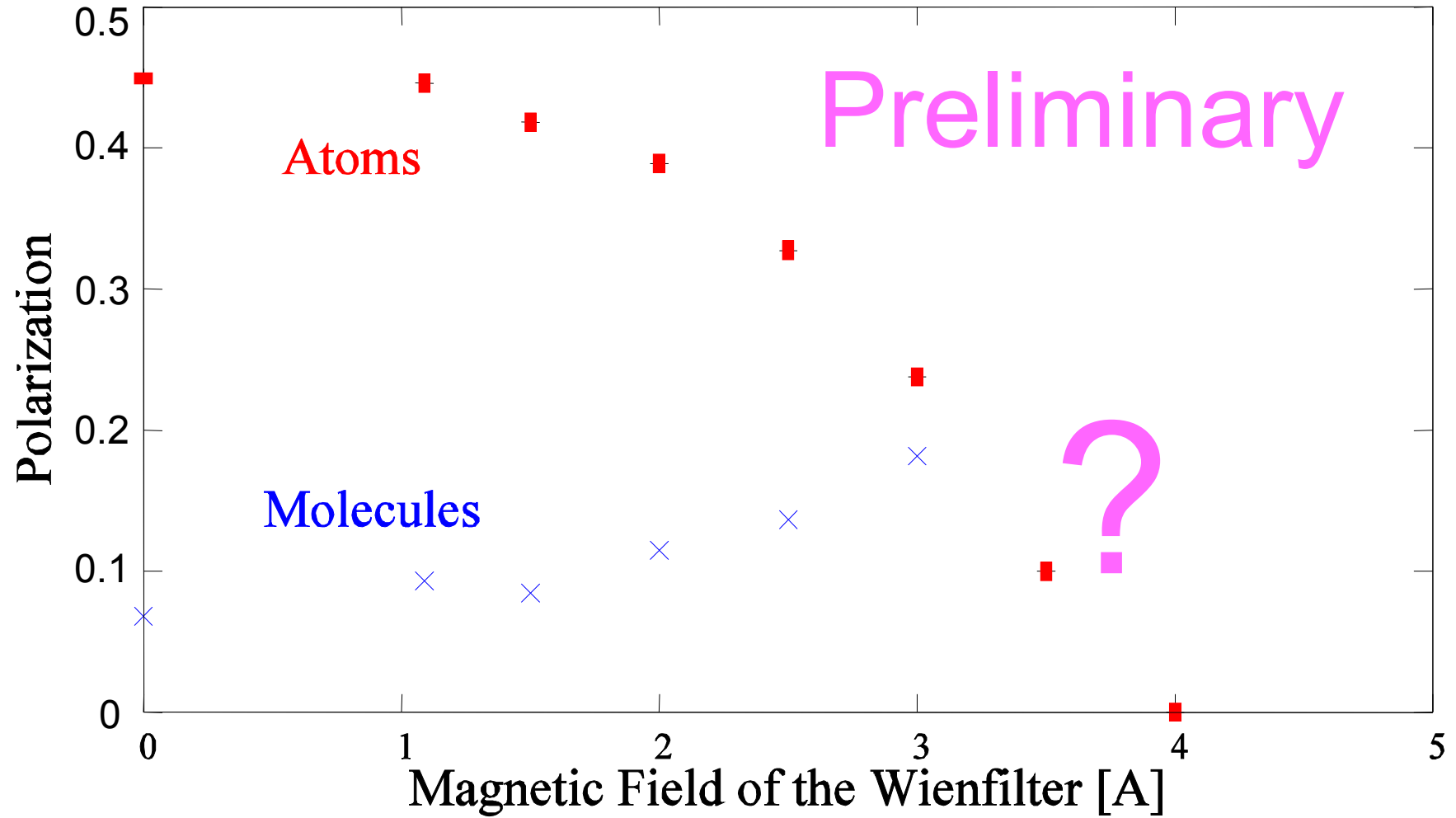
Polarized H₂ Molecules

Polarization of Hydrogen Molecules and Atoms
(Surface: Gold, HFS 1, B = 0.28 T, Q = 4 keV)

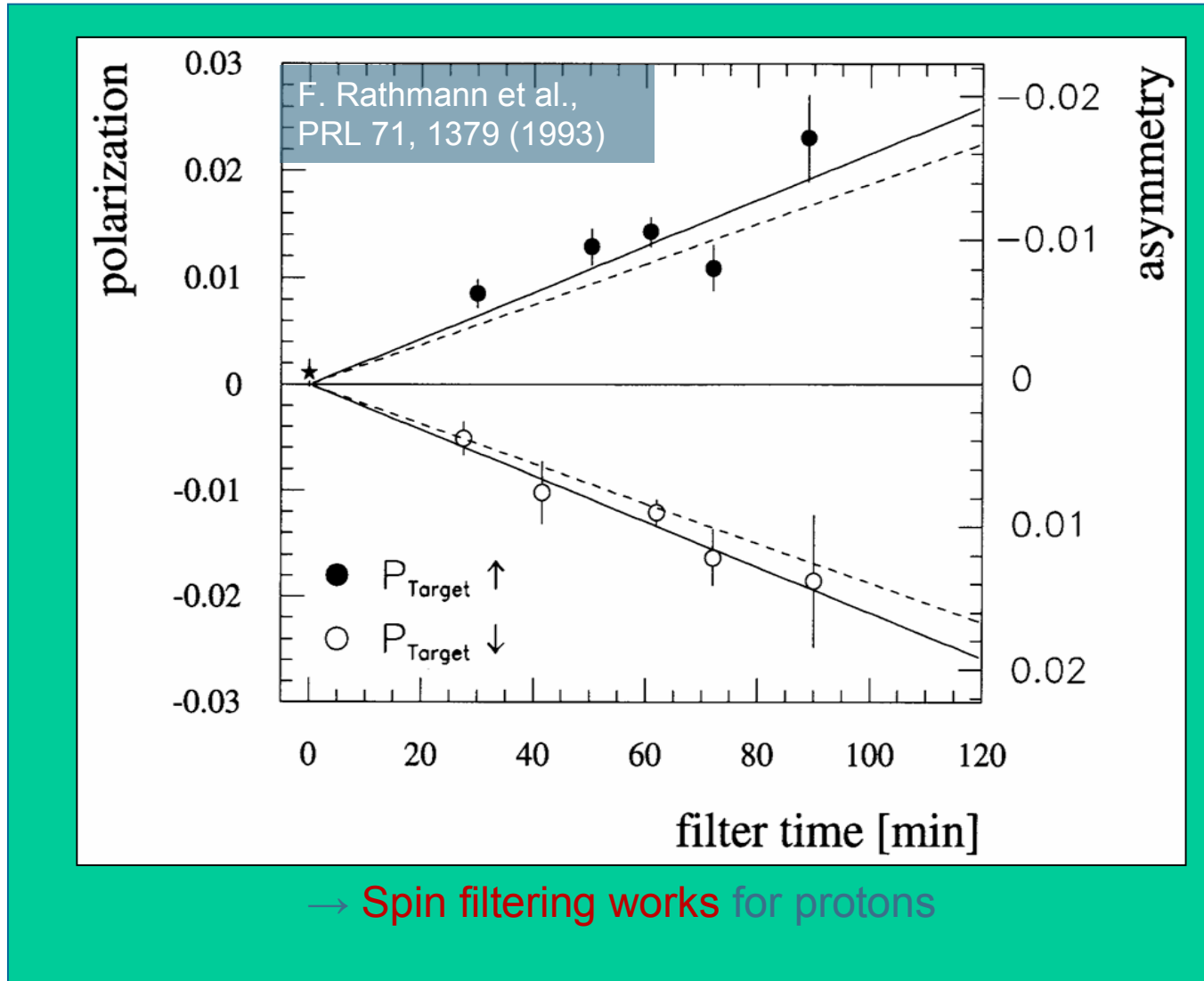


Polarized H₂ Molecules

Wienfilter Function of the Polarization

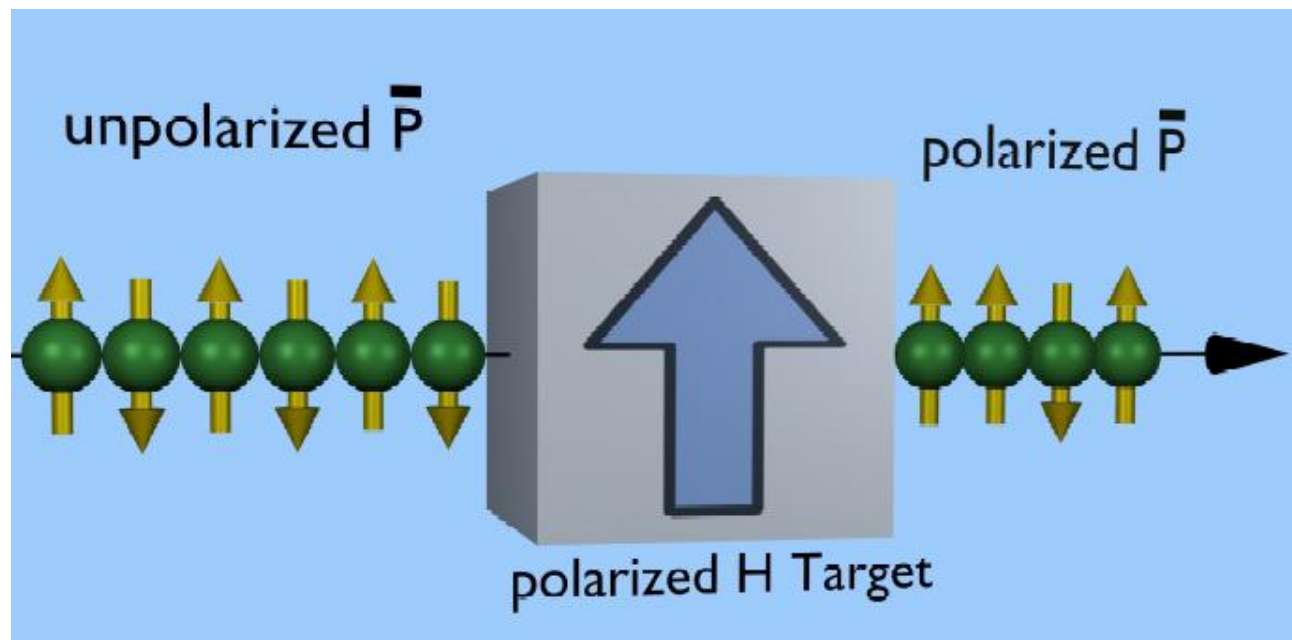


Spin-filtering at TSR: „FILTEX“ – proof-of-principle



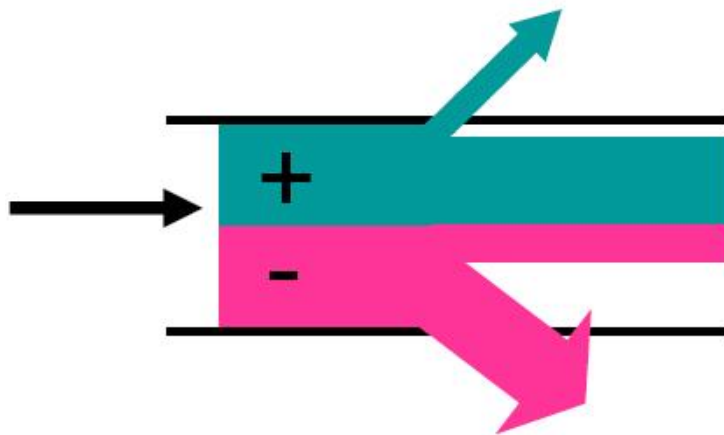
Spin-filtering

Polarization build-up of an initially unpolarized particle beam by repeated passage through a polarized hydrogen target in a storage ring:



PAX

Polarizing antiprotons: Two Methods: Loss



selective loss

discard (one) **substate**
(more than the other)

Eur. Phys. J. A **34**, 447–461 (2007)
DOI 10.1140/epja/i2007-10462-x

THE EUROPEAN
PHYSICAL JOURNAL A

Special Article – Tools for Experiment and Theory

A surprising method for polarising antiprotons

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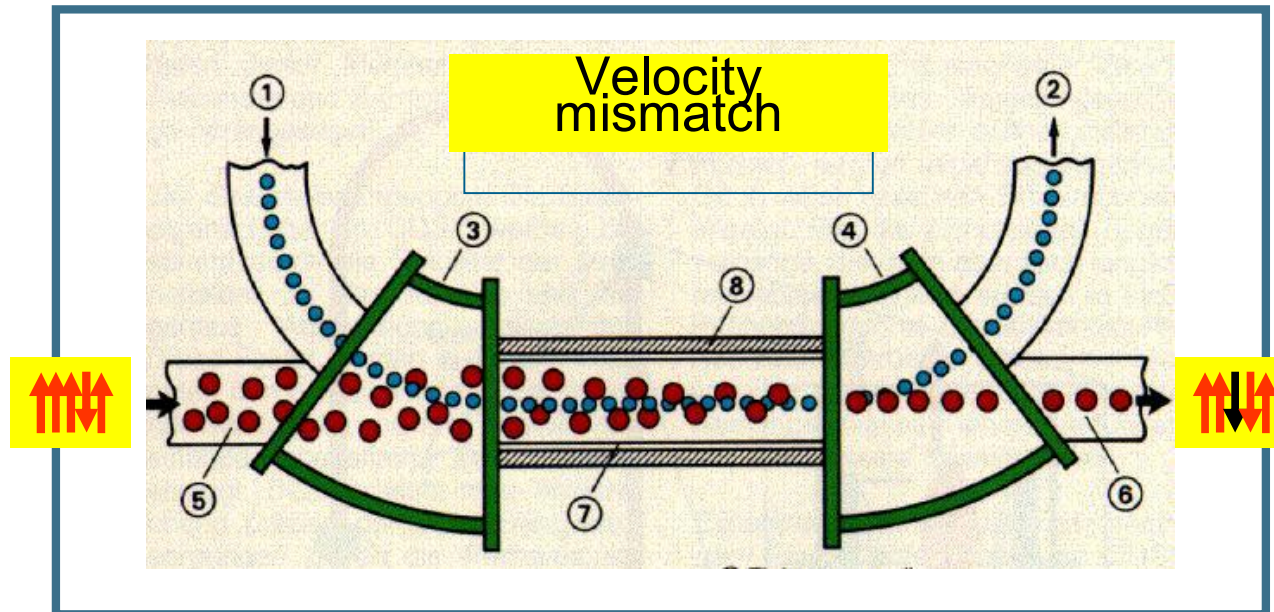
Communicated by E. De Sanctis

Abstract. We propose a method for polarising antiprotons in a storage ring by means of a polarised positron beam moving parallel to the antiprotons. If the relative velocity is adjusted to $v/c \approx 0.002$ the cross-section for spin-flip is as large as about $2 \cdot 10^{13}$ barn as shown by new QED calculations of the triple spin cross-

→ Need for an experimental test of this idea!

ep spin flip studies at COSY: Idea

- Use **proton** beam and co-moving **electrons**
- Turn experiment around: $p \vec{e} \rightarrow p \vec{e}$ into $p \vec{e} \rightarrow p \vec{e}$
i.e. observe **depolarization** of a polarized proton



Spin-filtering studies at COSY

Main purpose:

1. Repeat spin-filtering with protons. No surprises expected
2. Commissioning of the experimental setup for AD

Proposal to COSY PAC will be submitted in July 2009

Low- β magnet installation at COSY

