

Transverse Lambda Polarization at HERMES

S.Belostotski

Team:

S.Belostotski

Yu.Naryshkin

D.Veretennikov

O.Grebenjuk

Help and support

K.Rith

N.Makins

G.Schnell

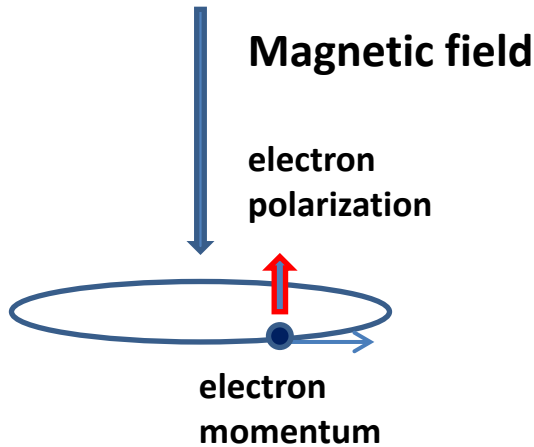
Published:

Phys.Rev.D 2014,

Phys.Rev.D 2007

Spontaneous Transverse Polarization

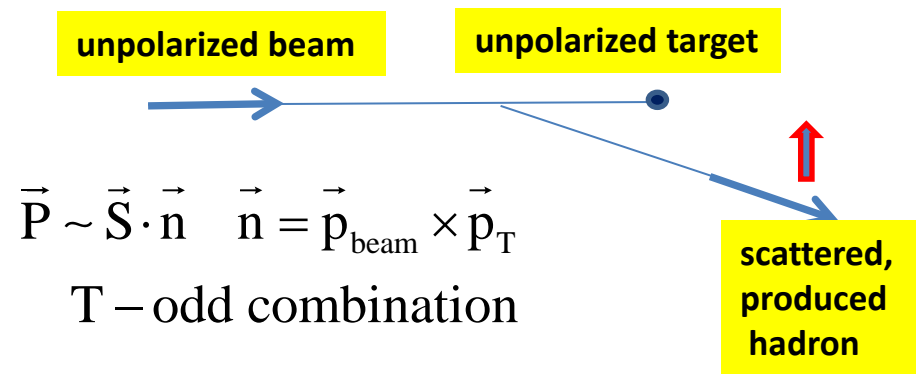
Sokolov-Ternov effect



Directed along magnetic field

Mechanism understood
in the frame of QED

Transverse polarization in hadron scattering



Polarization along normal \mathbf{n} to scattering
(production) plane

Well known phenomenon in pp , πp , pA
scattering in GeV energy domain.

In general Single Spin Asymmetry SSA

SSA at High Energies

Perturbative QCD in collinear factorization approach, e.g. contribution subprocess $qg \rightarrow qg$: quark of p_{beam} interacts with gluon of p_{target}

$$\sigma \sim q(x_b) \otimes g(x_t) \otimes \hat{\sigma}_{qg \rightarrow qg} \otimes D_{q \rightarrow \Lambda}(z) \quad (D_{q \rightarrow \Lambda}(z) \text{ in collinear approach})$$

$$P_{qg \rightarrow qg} \sim \alpha_s \frac{m_q}{\sqrt{S}} \text{ small} \Rightarrow 0 \text{ at } m_q \rightarrow 0$$

G.L.Kane, et al PRL 1978

At high energy substantial spin effects phenomena of produced hyperons are observed in disagreement with naïve perturbative QCD expectation

In order to see transverse spin transverse momentum to take into account:

$$D_{q \rightarrow \Lambda}(z) \Rightarrow D_{q \rightarrow \Lambda}(k_{\perp}, z) = D_{q \rightarrow \Lambda}^{\uparrow}(k_{\perp}, z) + D_{q \rightarrow \Lambda}^{\downarrow}(k_{\perp}, z)$$

$$P_{\Lambda} = \frac{\sigma_{\Lambda}^{\uparrow} - \sigma_{\Lambda}^{\downarrow}}{\sigma_{\Lambda}^{\uparrow} + \sigma_{\Lambda}^{\downarrow}} \sim \frac{D_{q \rightarrow \Lambda}^{\uparrow}(k_{\perp}, z) - D_{q \rightarrow \Lambda}^{\downarrow}(k_{\perp}, z)}{D_{q \rightarrow \Lambda}^{\uparrow}(k_{\perp}, z) + D_{q \rightarrow \Lambda}^{\downarrow}(k_{\perp}, z)}$$

***Collins Fragmentation Function,
P.J.Mulders and R.D.Tangerman,
Nucl.Phys 1996***

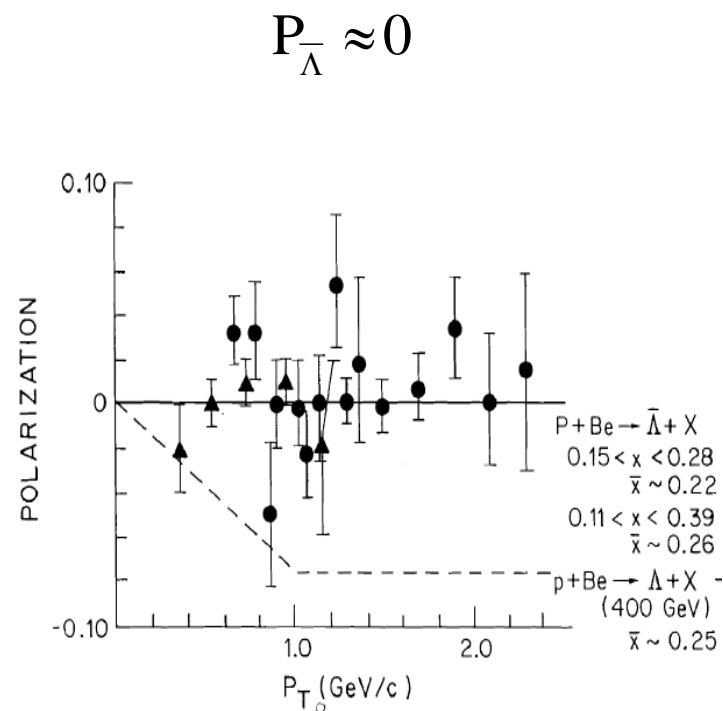
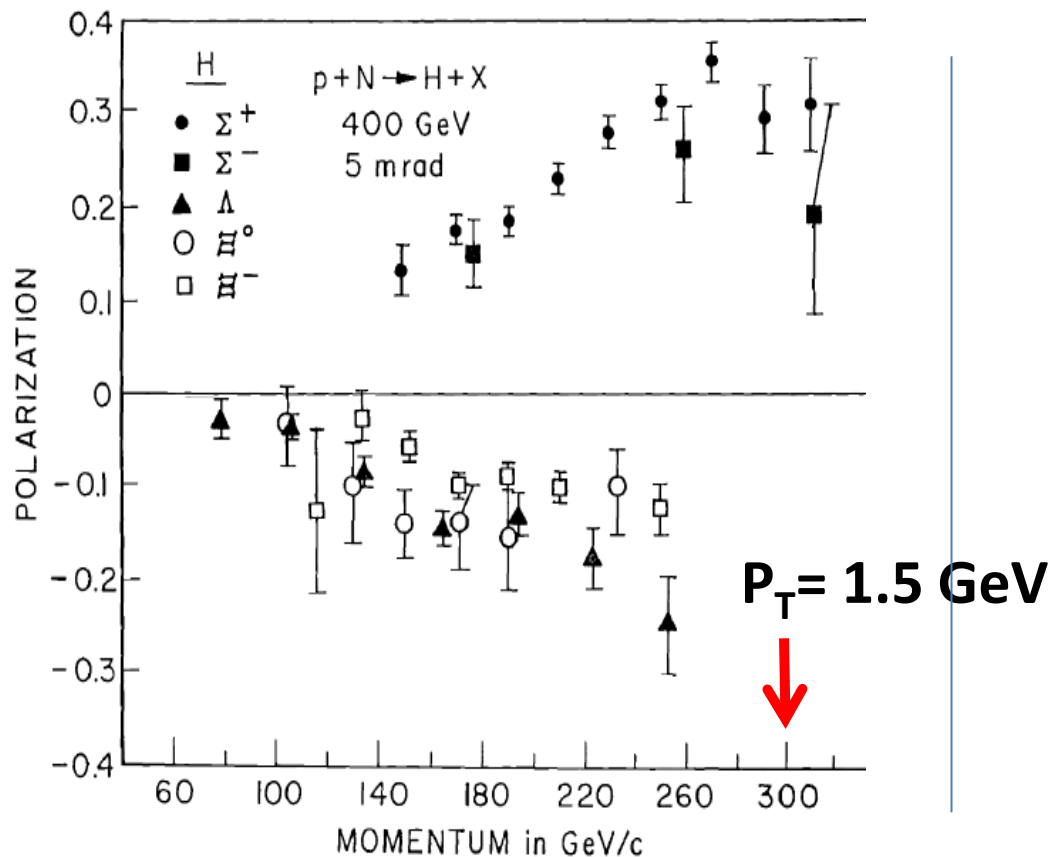
Transverse polarization is nonperturbative effect.

Unpolarized quark becomes polarized in hadronization process.

Recently: twist-3 FF (quark-gluon correlations) \rightarrow SSA

K.Kanazawa and Y.Koike 2013-2015 based on idea of Efremov and Teryaev

Hyperon Polarization . Famous Fermilab Data.



Interpretation: DeGrand and Mettinen PRD 1981,1985

Hyperon Octet Spin Structure $SU(3)_f$

	$\Delta(u+\bar{u})$	$\Delta(d+\bar{d})$	$\Delta(s+\bar{s})$
p(uud)	0.84	-0.43	-0.09
n(udd)	-0.43	0.84	-0.09
$\Lambda^0(uds)$	-0.16	-0.16	0.64
$\Sigma^+(uus)$	0.84	-0.09	-0.43
$\Sigma^0(uds)$	0.375	0.375	-0.43
$\Sigma^-(dds)$	-0.09	0.84	-0.43
$\Xi^0(uss)$	-0.43	-0.09	0.84
$\Xi^-(dss)$	-0.09	-0.43	0.84

$\Delta\Sigma=0.32$ HERMES/COMPASS

$$\Delta\Sigma = \Delta(u+\bar{u}) + \Delta(d+\bar{d}) + \Delta(s+\bar{s})$$

Assuming that

u(d) quark positively polarized

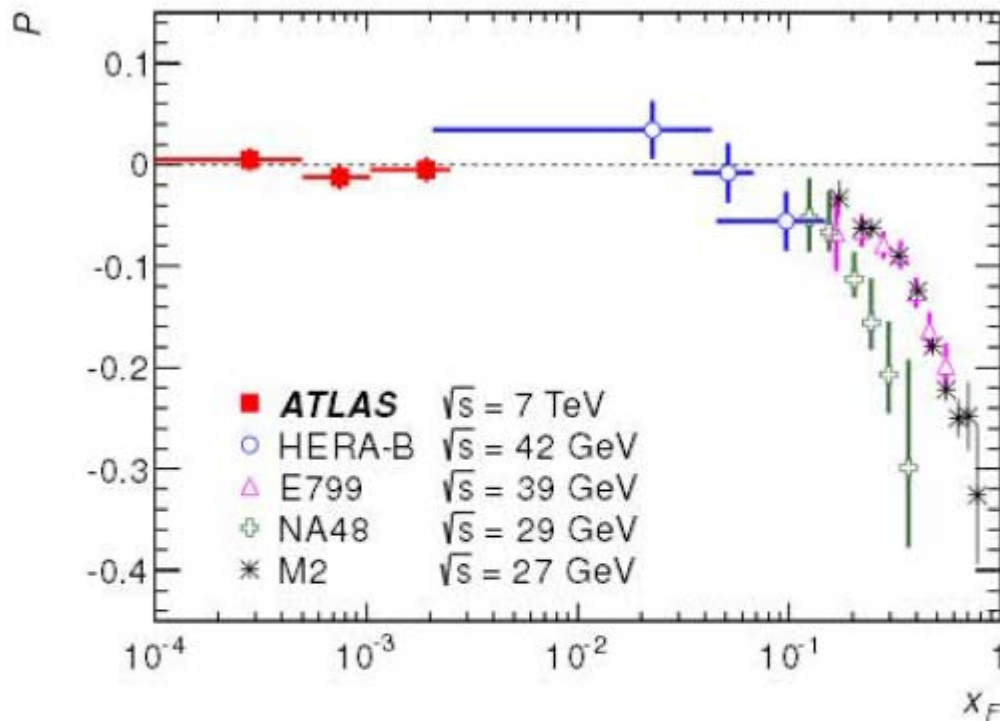
$\Lambda^0 \Xi^0 \Xi^-$ polarization negative

$\Sigma^+ \Sigma^-$ polarization positive

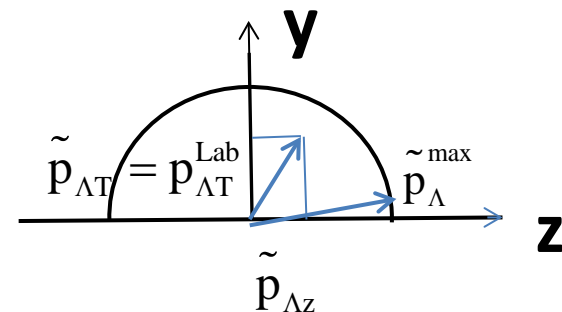
Lattice-QCD

$$\Delta u = \Delta d = -0.02 \quad \Delta s = 0.68 \quad (\pm 0.04)$$

More about Λ polarization. LHC/ATLAS Results.



x_F variable



$$x_F = \frac{\tilde{p}_{\Lambda z}}{\tilde{p}_{\Lambda}^{\max}} \quad -1 < x_F < +1$$

$\tilde{p}_{\Lambda z}, \tilde{p}_{\Lambda}^{\max}$ in beam-target c.m. frame,
z along beam particle momentum

FIG. 8. The Λ transverse polarization measured by ATLAS compared to measurements from lower center-of-mass energy experiments. HERA-B data are taken from Ref. [5], NA48 from Ref. [4], E799 data from Ref. [3], and M2 from Ref. [2]. The HERA-B results are transformed to positive values of x_F using Eq. (1).

Transverse Λ polarization summary

- P_Λ weakly depending of the beam momentum
- Its magnitude rises with p_T and x_F up to $P_\Lambda \approx 0,3 \div 0.4$ at $p_T = 1$ GeV
- At $1 < p_T < 3$ GeV P_Λ stays constant (at larger p_T no data available)
- In pp , pA inclusive production P_Λ negative
- Polarization of anti Λ compatible with zero
- For K^- , Σ^- beams P_Λ positive (valence s -quark)
- P_Λ shows a tendency to decrease with A (by 20% for Cu vs Be)
- At heavy ion collisions P_Λ similar to pp at $p_T > 2$ GeV
- P_Λ in electro/photo production practically unknown (2 tagged experiments CERN ,SLAC very poor statistics)

Transverse Λ polarization in photoproduction at HERMES

Deep Inelastic Scattering regime $\vec{e} + p, A \rightarrow e' + \underbrace{\bar{\Lambda} + X}_{\text{both detected, } Q^2 \text{ defined}}$ $Q^2 > 1 \text{ GeV}^2$ DLL

both detected,
 Q^2 defined

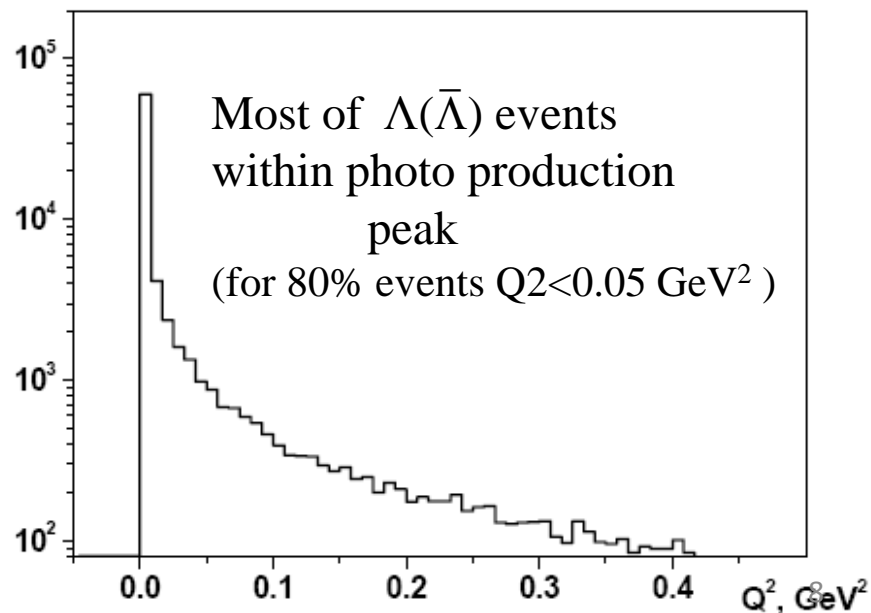
$$Q^2 = -4E_e E_{e'} \sin^2 \frac{\theta_{\text{Lab}}}{2} \quad Q^2 \rightarrow 0 \text{ at } \theta_{\text{Lab}} \rightarrow 0$$

Quasi-real photoproduction regime

$$e + p, A \rightarrow \bar{\Lambda} + X \quad Q^2 \approx 0$$

only detected
(inclusively)

Λ yield



Kinematics of Quasi-Real Photoproduction

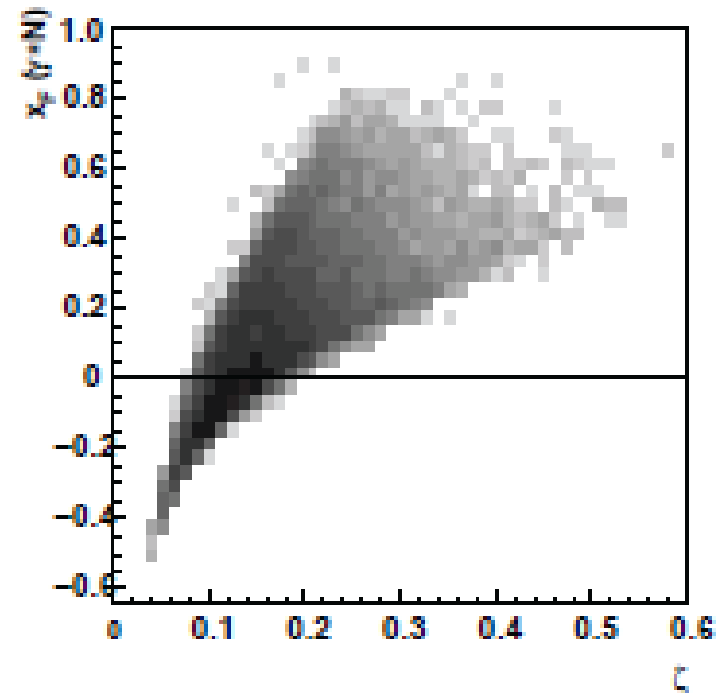
$$\gamma^* + p, A \rightarrow \Lambda + X \quad \Lambda \Rightarrow p\pi^- \quad \langle E_{\gamma^*} \rangle \approx 15 \text{ GeV} \text{ at HERMES}$$

$$\gamma^* + p, A \rightarrow \bar{\Lambda} + X \quad \bar{\Lambda} \Rightarrow \bar{p}\pi^+$$

As photons not tagged and only measured are Λ (anti Λ) momentum components $p_{\Lambda T}$ and $p_{\Lambda z}$, x_F cannot be built (!)

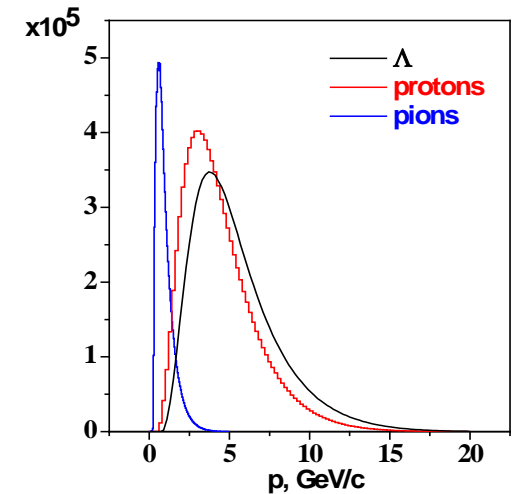
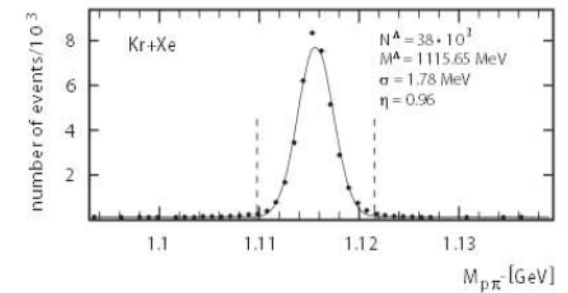
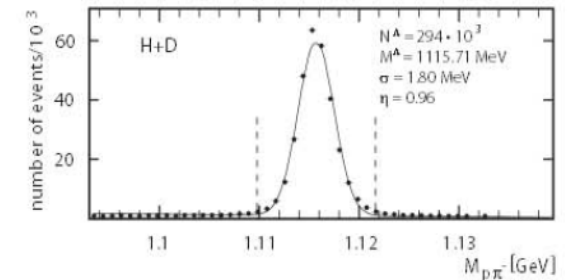
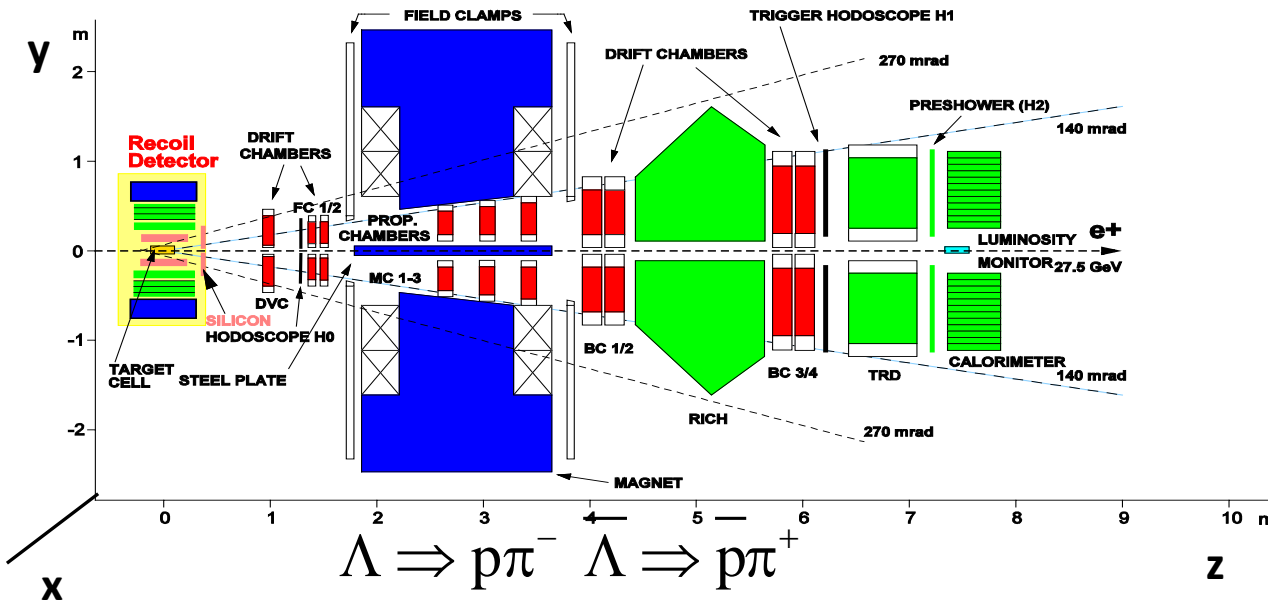
$$x_F \rightarrow \zeta = \frac{E_\Lambda + p_{z\Lambda}}{E_e + p_{z\Lambda}} \text{ light cone variable}$$

two variables define reaction kinematics
transverse momentum $p_{\Lambda T}$
and
light - cone variable ζ



x_F vs light - cone variable for HERMES
from PYTHIA MC

Λ detection with HERMES spectrometer



- Forward spectrometer with $\Delta p/p=0.01$ and $\Delta\theta=0.001$
- Polarized beam and H,D targets
- Gas targets H, D, He, Ne, Kr, Xe
- Up/down mirror symmetry (except solenoid fields)
- $0.6 < p_{\pi} < 2.5$ GeV cutoff low momentum pions
- Well-operational RICH, effective bgr suppression

Λ polarization measurement

Standard way to measure particle polarization is scattering off analyzing target:

$$\frac{d\sigma}{d\Omega_{\text{anal}}} = \frac{d\sigma}{d\Omega_0} (1 + PA \cos \phi) \quad P \text{ from azimuthal (left-right) asymmetry measurements.}$$

Difficulty: at high energy A small

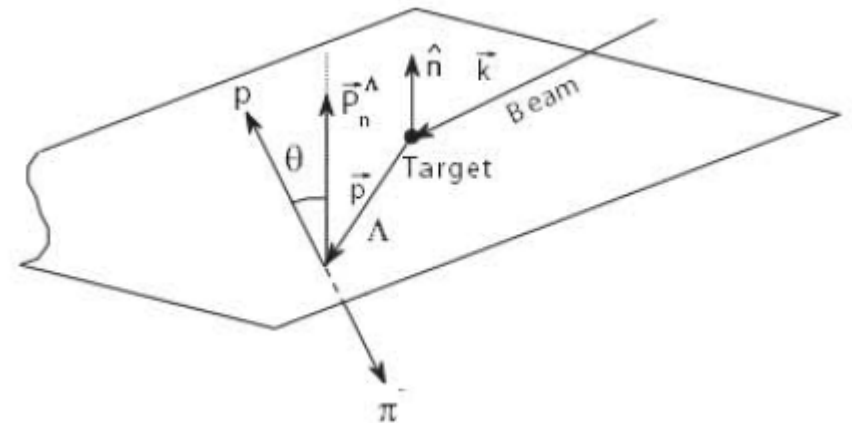
$\Lambda \rightarrow p\pi$ weak decay is self-analyzing:

$$\frac{d\sigma}{d\Omega_p} (\Lambda \rightarrow p\pi) = \frac{d\sigma}{d\Omega_{p0}} (1 + P_\Lambda \alpha \cos \theta_p) \quad \text{in } \Lambda \text{ rest frame}$$

$$\alpha = 0.642 \pm 0.13 \quad \text{for } \Lambda \rightarrow p\pi^-$$

$$\text{and } \alpha = -0.642 \quad \text{for } \bar{\Lambda} \rightarrow \bar{p}\pi^+$$

p_Λ from forward – backward decay asymmetry in Λ rest frame



Extraction of Λ polarization

$$\frac{dN}{d\Omega_p} = \frac{dN_0(\cos\theta_p)}{d\Omega_p} (1 + \alpha P_\Lambda \cos\theta_p)$$

$$\frac{dN_0(\cos\theta_p)}{d\Omega_p} \sim \varepsilon(\cos\theta_p)$$

Maximum likelihood (Moment method)

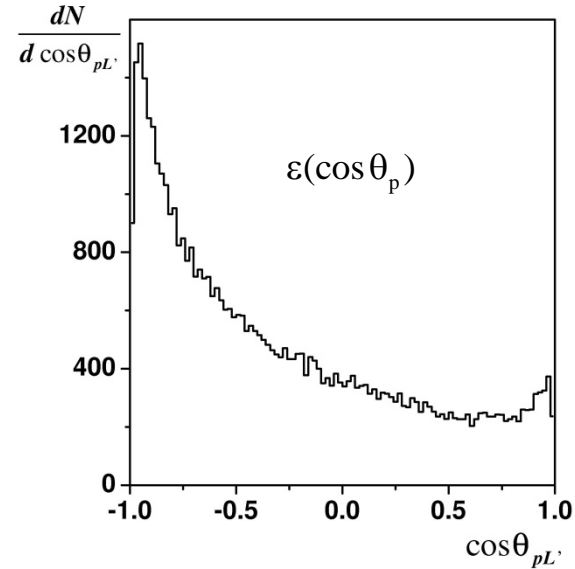
Normalized probability to detect Λ event

$$\omega_i = \frac{1 + \alpha P_\Lambda \cos\theta_{pi}}{1 + \alpha P_\Lambda \langle \cos\theta_p \rangle_0}$$

$$\langle \cos\theta_p \rangle_0 = \frac{\int \varepsilon(\cos\theta_p) \cos\theta_p d\Omega}{\int \varepsilon(\cos\theta_p) d\Omega} = 1$$

$$\cos\theta_p = \frac{p_{\pi x} p_{py} - p_{\pi y} p_{px}}{q \cdot p_{\Lambda T}}$$

$\cos\theta_p \rightarrow -\cos\theta_p$ at $y \rightarrow -y$



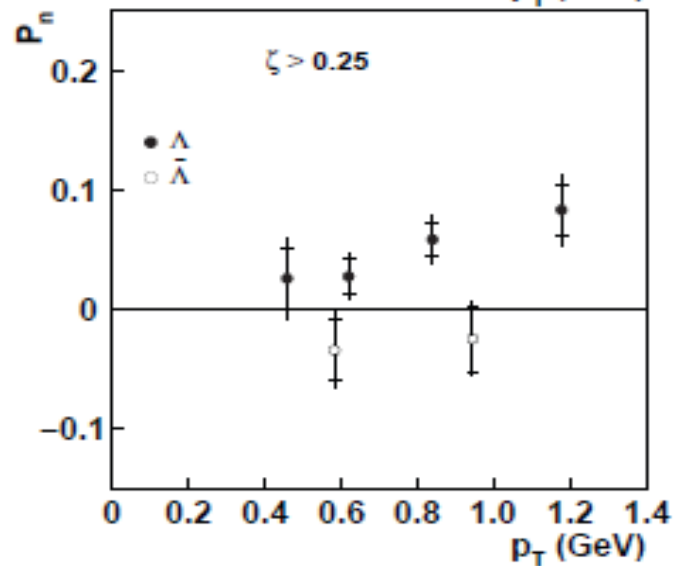
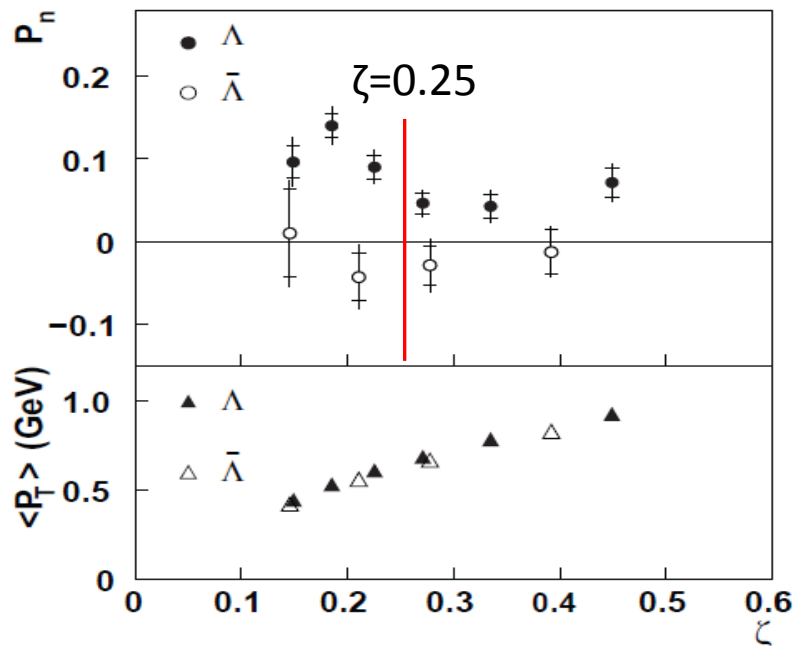
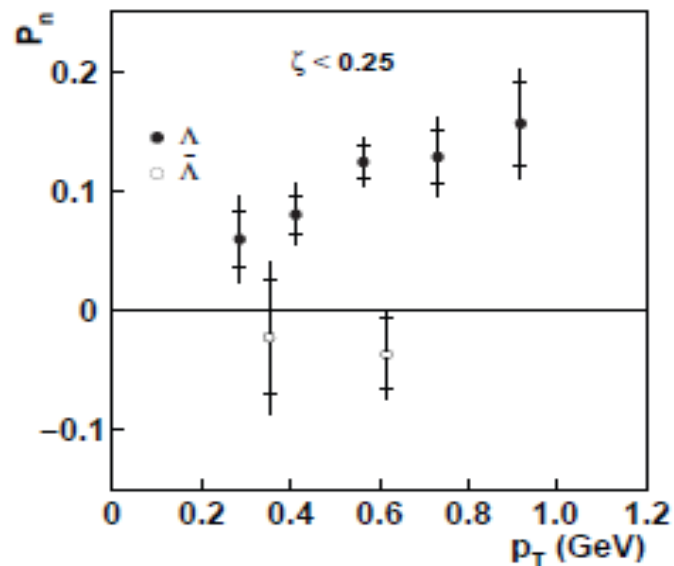
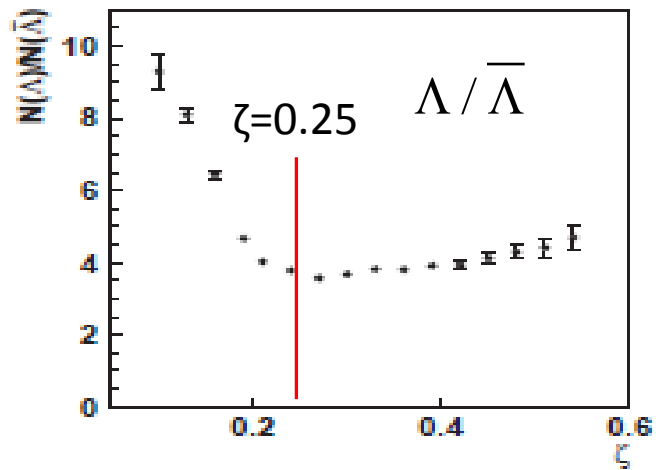
Up/Down mirror symmetry $\Rightarrow \varepsilon(-p_{\pi y}, -p_{py}) = \varepsilon(p_{\pi y}, p_{py})$ $\langle \cos\theta_p \rangle_0 = 0$

$$\omega_i = 1 + \alpha P_\Lambda \cos\theta_{pi} \quad L = \prod_{i=1}^{N_\Lambda} (1 + \alpha P_\Lambda \cos\theta_{pi}), \quad \frac{\partial \ln L}{\partial P_\Lambda} = \sum_{i=1}^{N_\Lambda} \frac{\alpha \cos\theta_{pi}}{1 + \alpha P_\Lambda \cos\theta_{pi}} = 0$$

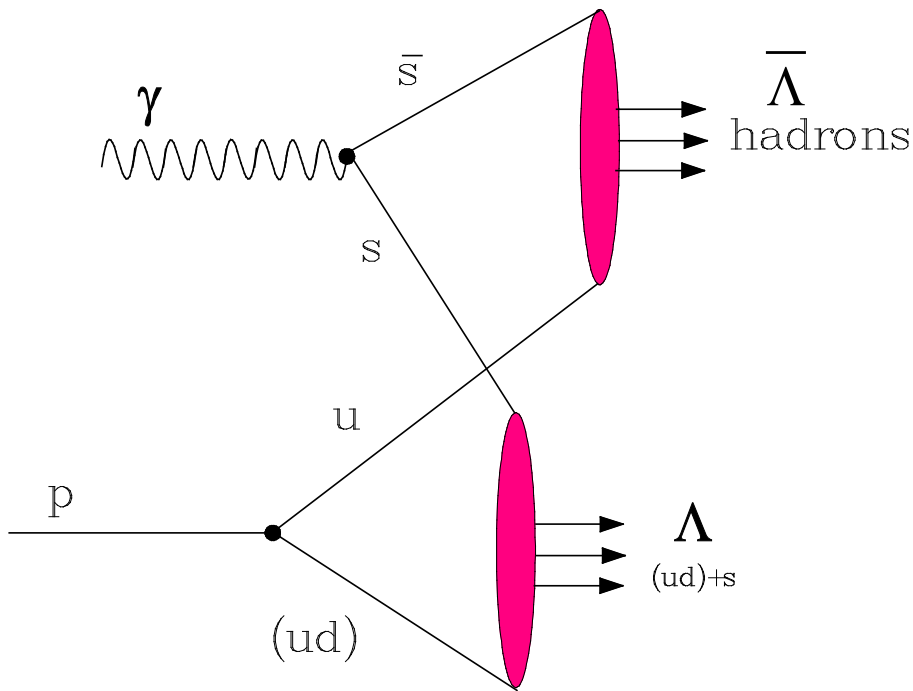
$$\sum_{i=1}^{N_\Lambda} \frac{\alpha \cos\theta_{pi}}{1 + \alpha P_\Lambda \cos\theta_{pi}} \approx \sum_{i=1}^{N_\Lambda} \alpha \cos\theta_{pi} (1 - \alpha P_\Lambda \cos\theta_{pi}) = 0 \quad \text{neglecting } (\alpha P_\Lambda \cos\theta_{pi})^n \quad n \geq 2$$

$$P_\Lambda = \frac{\sum_i^{N_\Lambda} \cos\theta_{pi}}{\alpha \sum_i^{N_\Lambda} \cos^2\theta_{pi}}$$

Results



Possible Interpretation



Current fragmentation

Λ production $s(\text{beam}) + (ud)_{0,1}(\text{string})$

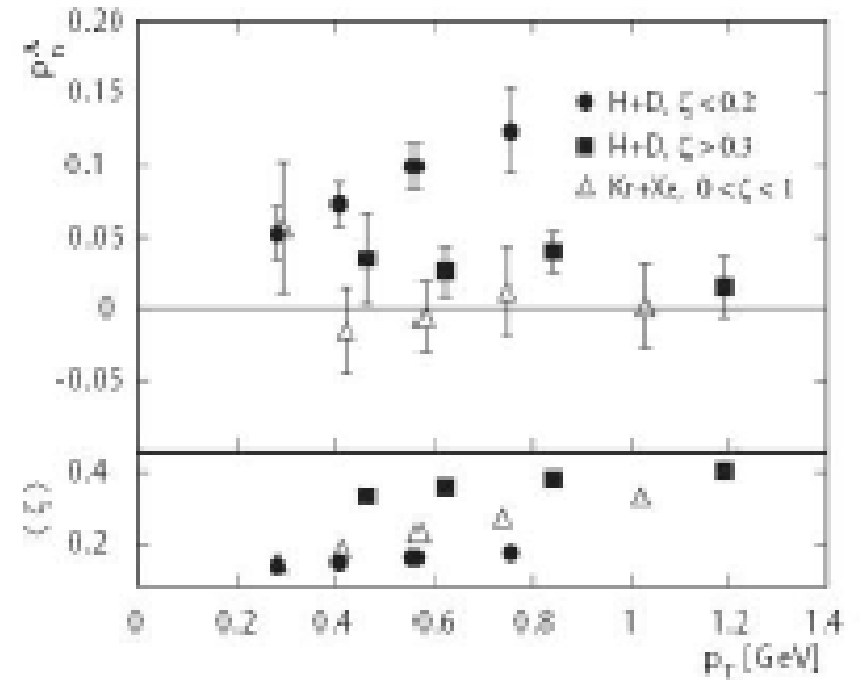
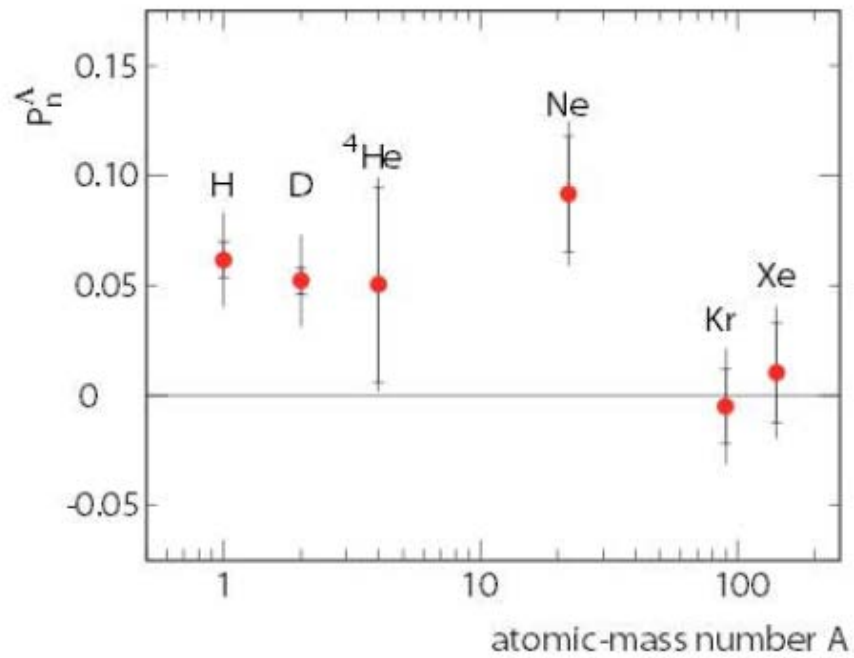
$\bar{\Lambda}$ production $\bar{s}(\text{beam}) + \overline{(ud)}_{0,1}(\text{string})$

Target fragmentation

Λ production $s(\text{beam}) + (ud)_{0,1}(\text{target})$

$\bar{\Lambda}$ production $\bar{s}(\text{beam}) + \overline{(ud)}_{0,1}(\text{target sea})$

A-dependence



Conclusion

- ❑ Polarization of Λ in inclusive photoproduction at $p_T < 1\text{GeV}$

$$P_\Lambda = 0.078 \pm 0.006_{\text{stat}} \pm 0.012_{\text{stat}}$$

positive as in the case of K^- or Σ^- beam

- ❑ Λ bar polarization compatible with zero:

$$P_{\Lambda\text{bar}} = -0.025 \pm 0.015_{\text{stat}} \pm 0.018_{\text{stat}}$$

- ❑ Λ polarization in target fragmentation ($\zeta < 0.25$) essentially larger than that in current fragment ($\zeta > 0.25$)

- ❑ Yield of Λ in target fragmentation surpasses substantially Λ bar

- ❑ A possible interpretation of production mechanism and (partly) observed polarizations relates photon dissociation to quark-antiquark pair

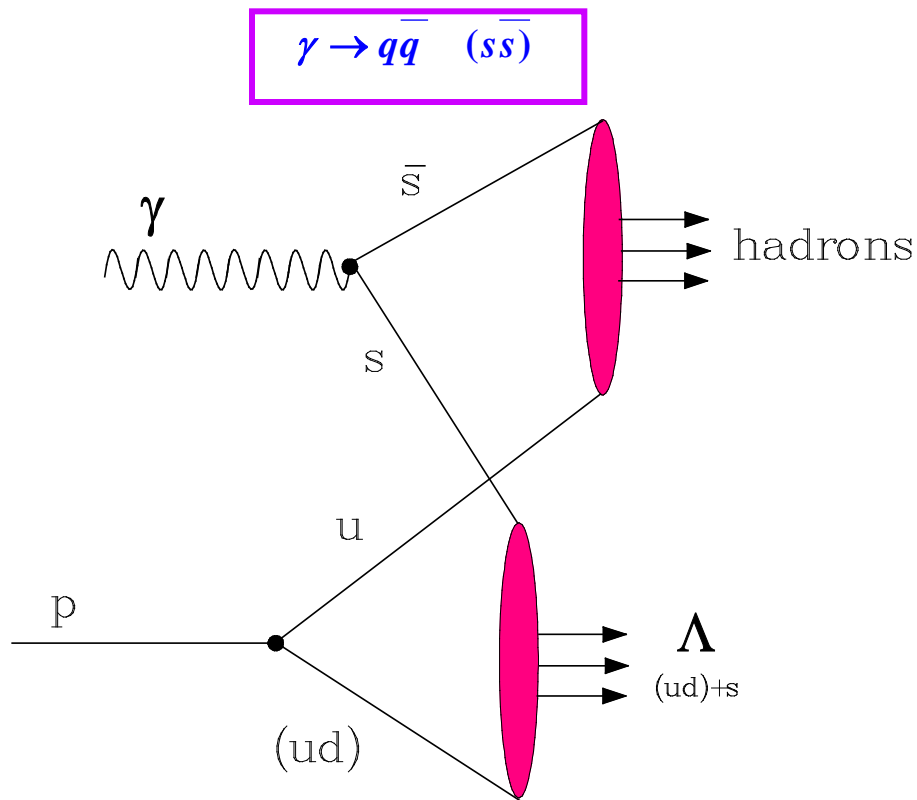
- ❑ A-dependence: polarization vanishes for $A \approx 100$

- ❑ Polarizations for H and D targets coincides within error bars

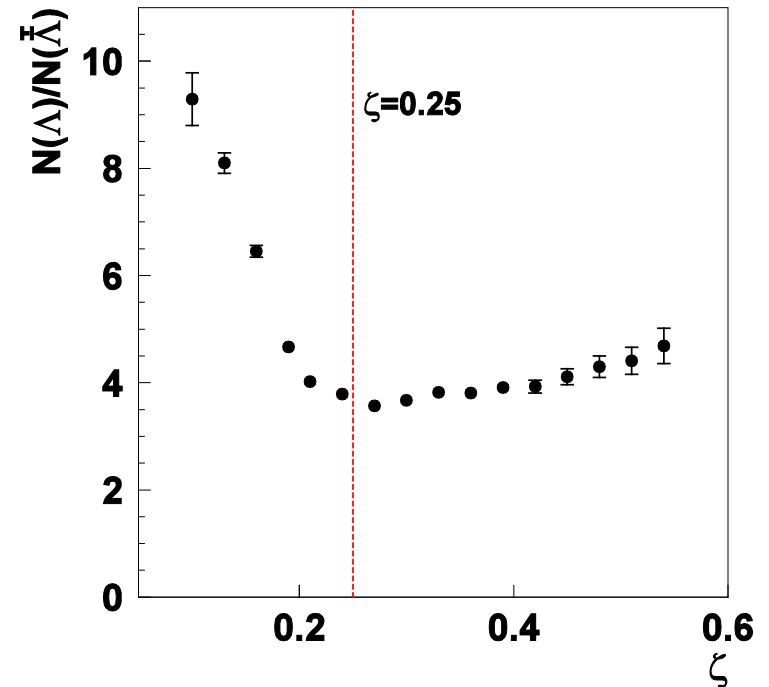
BACKUP

Λ photoproduction mechanism by PYTHIA

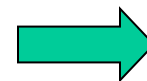
$$\langle E_\gamma \rangle = \langle E_e - E_{e'} \rangle \approx 15.6 \text{ GeV}$$



Λ to $\bar{\Lambda}$ yield ratio

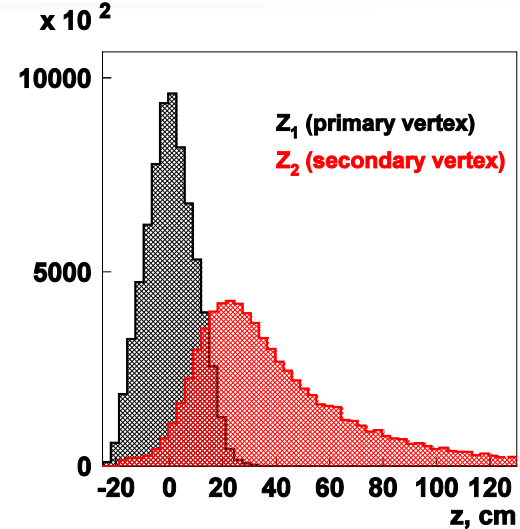
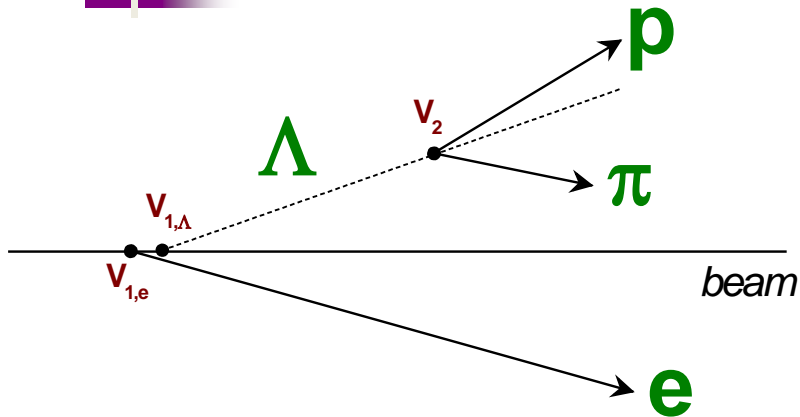


$$\zeta^\Lambda \approx \frac{E^\Lambda}{E_e} < 0.25 \text{ or } \sqrt{t} < 3.31 \text{ GeV}$$



*target (ud)
fragmentation
mechanism*

Λ and $\bar{\Lambda}$ events selection



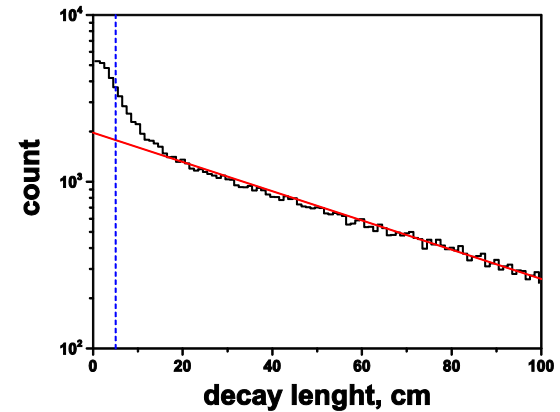
main bgr is $\pi^+ \pi^-$ and $K \pi$ pairs production;

bgr suppression cuts:

- Leading π rejection using threshold Cherenkov det. (1996-1997) or RICH (1998-2007)

- Vertex separation.

Distance between V_1 and V_2 vertices > 5 cm



Extraction of D_{Li} components from experimental data sample

$$\frac{dN}{d\Omega_p} = \frac{dN_0}{d\Omega_p} (1 + \alpha \vec{P}^\Lambda \cdot \hat{k}_p) = \frac{dN_0}{d\Omega_p} (1 + \alpha_\Lambda P_B \sum_{i=x,y,z} D_{Li}^\Lambda \cos \theta_i) \text{ in } \Lambda \text{ rest frame}$$

$$\alpha_{\Lambda \rightarrow p+\pi^-} = 0.642 \pm 0.013 \quad \alpha_{\Lambda \rightarrow \bar{p}+\pi^+} = -0.642 \pm 0.013$$

Spectrometer acceptance results in strong distortion of decay angular distribution,
intensive MC acceptance simulation (COMPSS)

For beam helicity balance case $[\mathbf{P}_B] = 0$

MC simulation of spectrometer acceptance is not needed, acceptance correction does not affect measured asymmetries. D_{Li} components are extracted using experimental data sample only !!

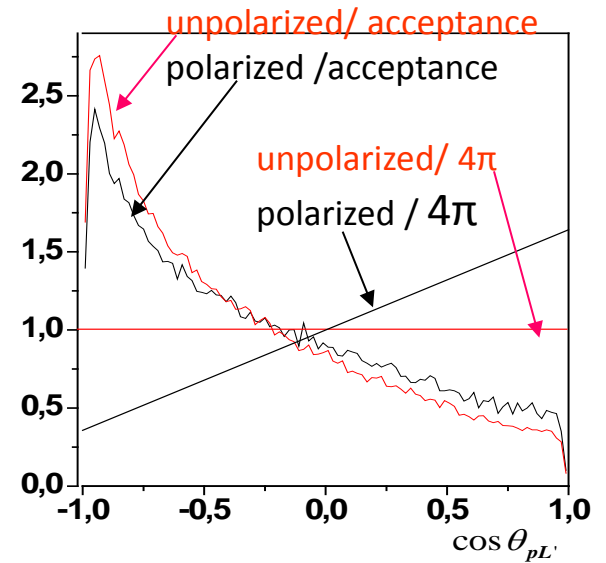
$$\sum_{k=x,y,z} D_{Lk} A_{ik} = \frac{1}{\alpha} \frac{B_i}{[\mathbf{P}_B^2]} \quad i = x, y, z$$

$$A_{ik} = \frac{1}{N^\Lambda} \sum_{v=1}^{N^\Lambda} (D^2(y) \cos \theta_i \cos \theta_k)_v$$

$$B_i = \frac{1}{N^\Lambda} \sum_{v=1}^{N^\Lambda} (P_B D(y) \cos \theta_i)_v$$

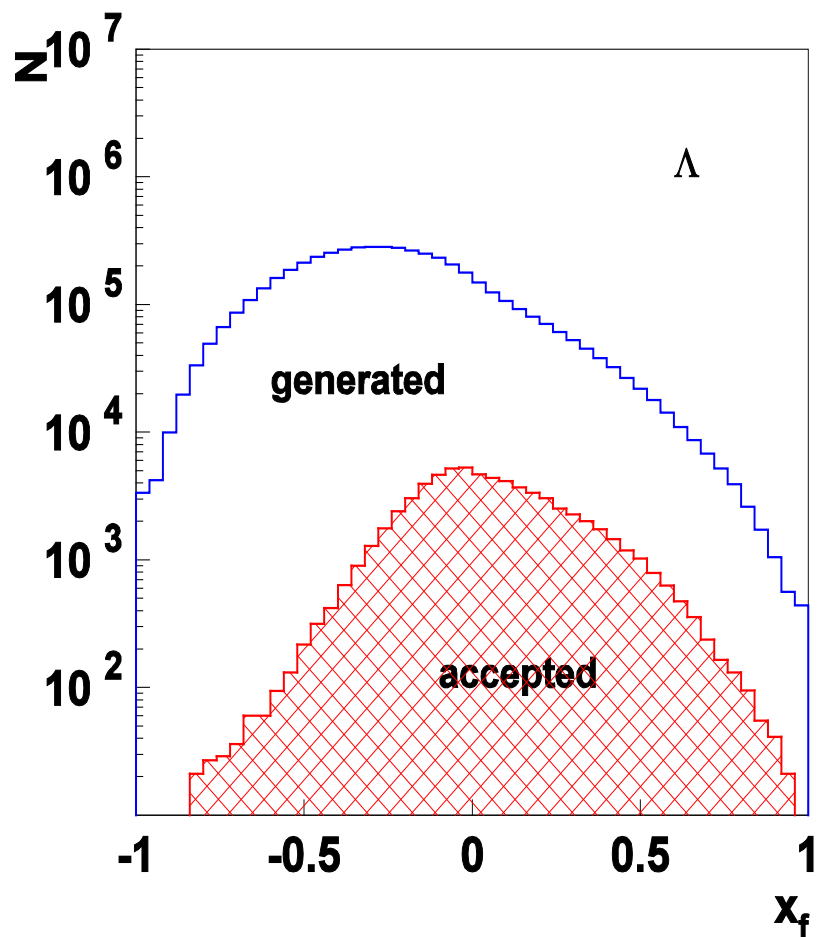
$$[\mathbf{P}_B^2] = \frac{\int P^2(t) L(t) dt}{\int L(t) dt}$$

average over experimental data sample



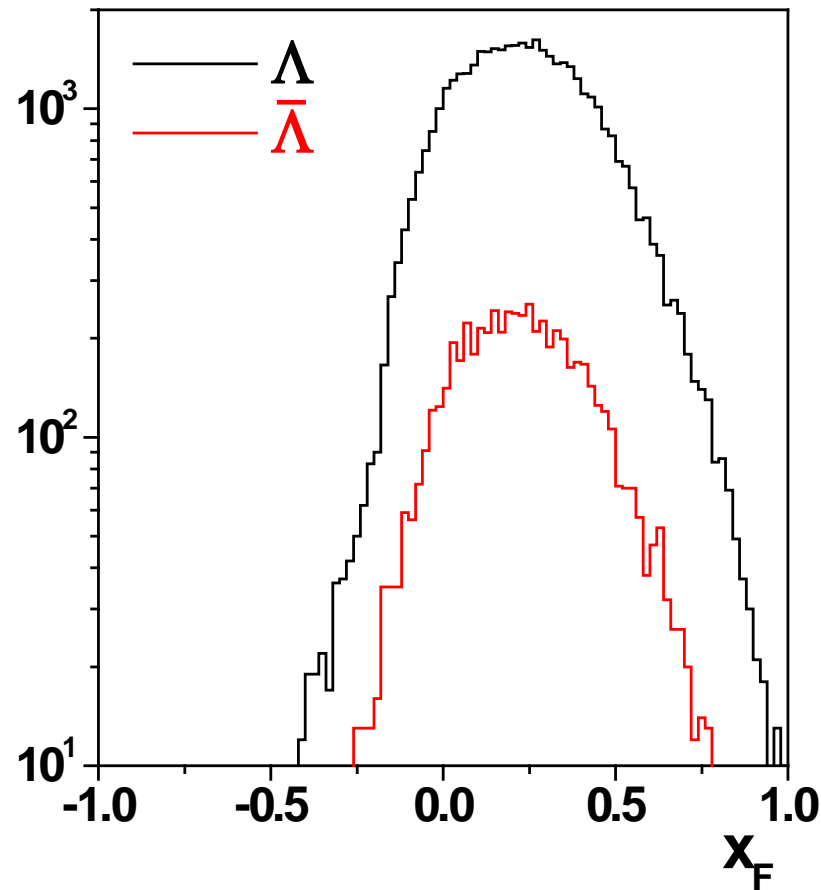
x_F distributions (semi-inclusive DIS)

Pythia MC for Λ



HERMES is a forward spectrometer

Experiment for Λ and $\bar{\Lambda}$



$\rightarrow p_{\Lambda}(\text{min}) \sim 1 \text{ GeV}$

Extraction of $D_{LL'}$

- Angular distribution of decay protons in Λ rest frame



$$\frac{dN}{d\Omega_p} = \frac{dN_0}{d\Omega_p} (1 + \alpha P_{\text{Beam}} \vec{D}_{LL'}^\Lambda \cdot \hat{k}_p)$$

$$\alpha = 0.642 \pm 0.013$$

$\frac{dN_0}{d\Omega_p} = \text{const}$ for 4π acceptance

for restricted acceptance

$\frac{dN_0}{d\Omega_p}$ depends on $\cos\theta_{pL'}$



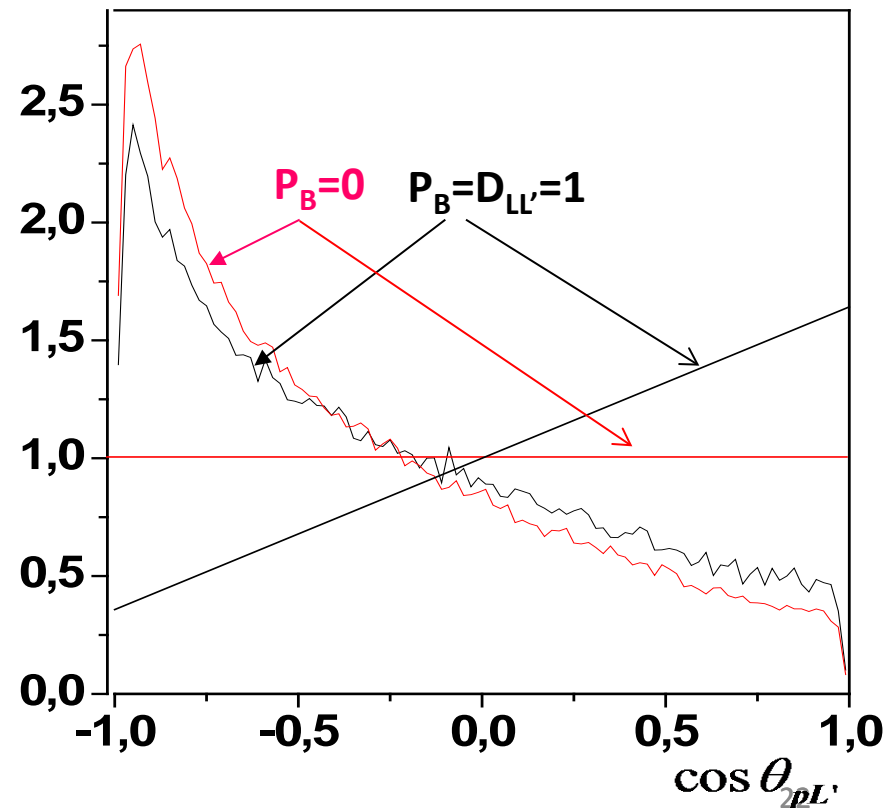
Distorted by spectrometer acceptance



May in principle be calculated using MC



difficulty to avoid false asymmetry induced by MC acceptance simulation



Effect of longitudinal magnetic field (solenoid)

$$\cos\theta_p = \frac{p_{\pi x} p_{py} - p_{\pi y} p_{px}}{p_{\Lambda T} q} \quad \cos\theta_p \rightarrow -\cos\theta_p \text{ at } y \rightarrow -y$$

$$\cos\theta_p \rightarrow -\cos\theta_p \text{ at } y \rightarrow -y$$

$\cos\theta_p$ in Λ rest frame,

$q = 101 \text{ Mev}$ decay momentum

$p_{\pi x,y}, p_{px,y}, p_{\Lambda T}$ in Lab frame (!)

Up/down mirror symmetry: acceptance function

$$\varepsilon^{\text{up}}(\cos\theta_p, \dots) \rightarrow \varepsilon^{\text{down}}(-\cos\theta_p, \dots) \text{ at } y \rightarrow -y$$

**Transverse magnetic field of the dipole magnet
(transverse pol target and spectrometer dipole)**

$$B_x^{\text{up}} = -B_x^{\text{down}} \quad B_y^{\text{up}} = B_y^{\text{down}} \quad B_z^{\text{up}} = -B_z^{\text{down}}$$

Lorentz force

$$F_x^{\text{up}} = F_x^{\text{down}} \quad F_y^{\text{up}} = -F_y^{\text{down}} \quad F_z^{\text{up}} = F_z^{\text{down}}$$

up / down mirror symmetry

**Longitudinal magnetic field of solenoid
(longitudinal pol target and RD solenoid)**

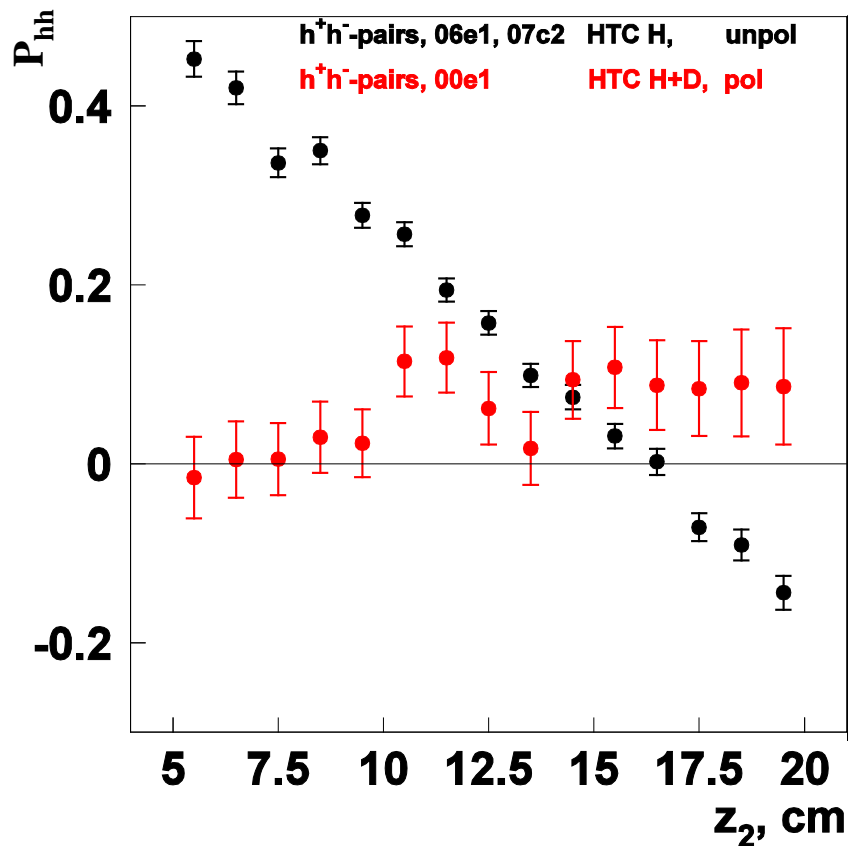
$$B_x^{\text{up}} = B_x^{\text{down}} \quad B_y^{\text{up}} = -B_y^{\text{down}} \quad B_z^{\text{up}} = B_z^{\text{down}}$$

Lorentz force

$$F_x^{\text{up}} = -F_x^{\text{down}} \quad F_y^{\text{up}} = F_y^{\text{down}} \quad F_z^{\text{up}} = -F_z^{\text{down}}$$

no up/down mirror symmetry

False polarization of hh pairs 00 and 06,07



*hh - pairs false polarization
1996-2005*

-20 < z₁ < 20 cm

pol + unpol 0.0088± 0.0034

pol 0.0232± 0.0054

unpol -0.0006± 0.0044

5 < z₁ < 20 cm (recoil target size)

pol + unpol 0.0455± 0.0059

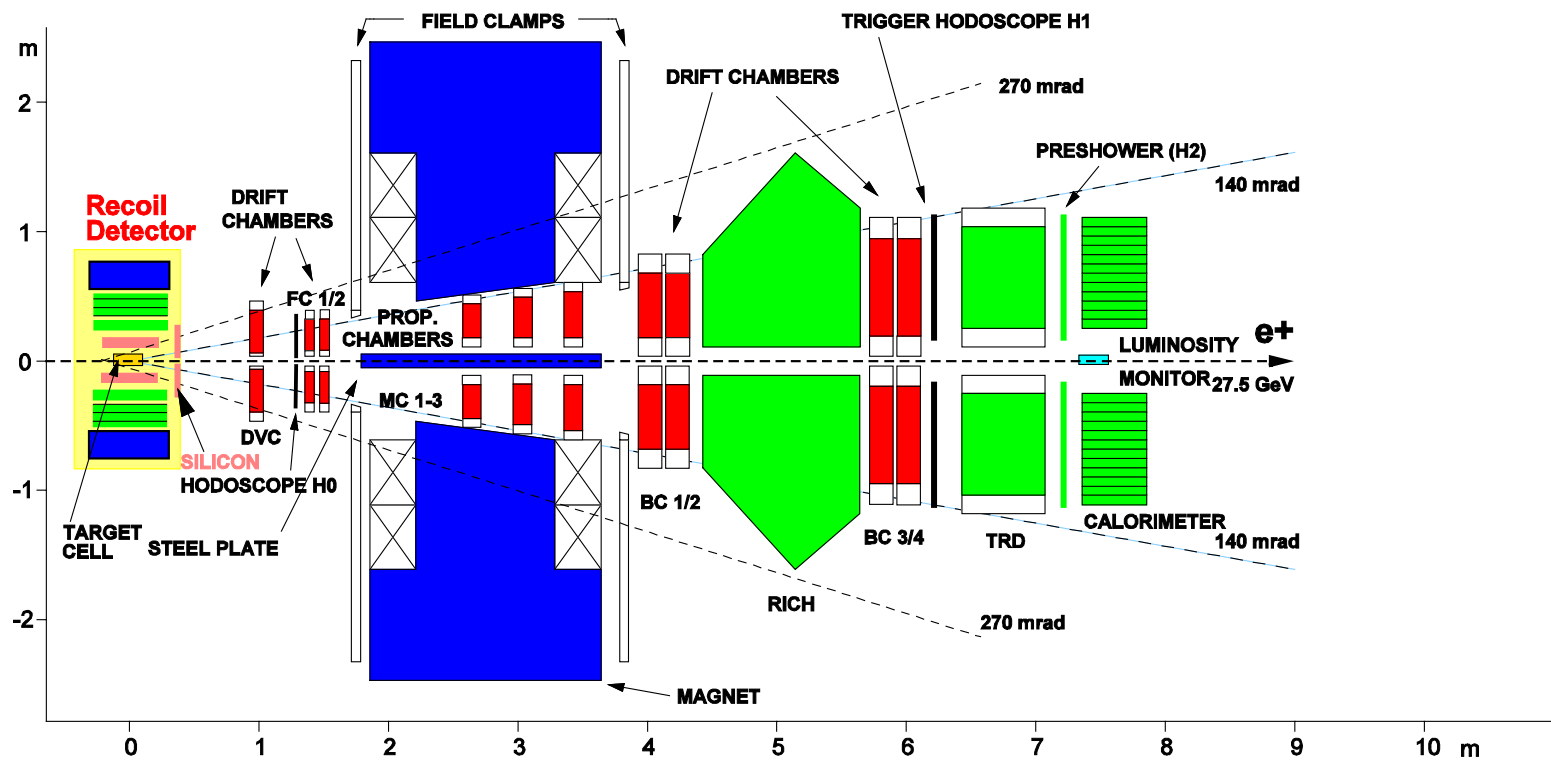
pol 0.0608± 0.0093

unpol 0.0355± 0.0076

Published false polarization of Ks is: $P_{Ks} = 0.012 \pm 0.004$

Published false polarization of hh pairs is: $P_{hh} = 0.012 \pm 0.002$

HERMES SPECTROMETER



HERMES dipole $BL=1.3 TM$ $\frac{\Delta p}{p} \approx 1\%$ $\Delta\theta_x, \Delta\theta_y \approx 1mrad$

 $-170 < \theta_x < +170mrad$ $-140 < \theta_y < -40mrad$ $140 > \theta_y > 40mrad$
 $40 < \theta < 220mrad$

Very good PID !!